

Hartree Fock HW

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1. Approximate the many-body wavefunctions of different systems in which the single electrons are under the harmonic oscillator potential. Recall we do this via the Slater determinant:

$$\Psi(\{\gamma_i\}) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \psi_1(\gamma_1) & \psi_1(\gamma_2) & \dots & \psi_1(\gamma_N) \\ \psi_2(\gamma_1) & \psi_2(\gamma_2) & \dots & \psi_2(\gamma_N) \\ \dots & \dots & \dots & \dots \\ \psi_N(\gamma_1) & \psi_N(\gamma_2) & \dots & \psi_N(\gamma_N) \end{vmatrix} \quad (1)$$

We will assume the space and spin components of the single-body wavefunctions are separable:

$$\psi(\gamma) = \phi(\mathbf{r})\chi(\boldsymbol{\sigma}) \quad (2)$$

- (a) Electrons a and b are spin-up and spin-down respectively, both in ground state:

$$\begin{aligned} \Psi(\gamma_a, \gamma_b) &= \frac{1}{\sqrt{2}} \begin{vmatrix} \psi_a(\gamma_a) & \psi_a(\gamma_b) \\ \psi_b(\gamma_a) & \psi_b(\gamma_b) \end{vmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_0(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) & \phi_0(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) \\ \phi_0(\mathbf{r}_a)\chi_-(\boldsymbol{\sigma}_a) & \phi_0(\mathbf{r}_b)\chi_-(\boldsymbol{\sigma}_b) \end{vmatrix} \\ &= \frac{1}{\sqrt{2}} \phi_0(\mathbf{r}_a)\phi_0(\mathbf{r}_b) [\chi_+(\boldsymbol{\sigma}_a)\chi_-(\boldsymbol{\sigma}_b) - \chi_+(\boldsymbol{\sigma}_b)\chi_-(\boldsymbol{\sigma}_a)] \end{aligned} \quad (3)$$

- (b) Electrons a and b are in the ground state and first excited state respectively, both spin-up:

$$\begin{aligned} \Psi(\gamma_a, \gamma_b) &= \frac{1}{\sqrt{2}} \begin{vmatrix} \psi_a(\gamma_a) & \psi_a(\gamma_b) \\ \psi_b(\gamma_a) & \psi_b(\gamma_b) \end{vmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{vmatrix} \phi_0(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) & \phi_0(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) \\ \phi_1(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) & \phi_1(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) \end{vmatrix} \\ &= \frac{1}{\sqrt{2}} \chi_+(\boldsymbol{\sigma}_a)\chi_+(\boldsymbol{\sigma}_b) [\phi_0(\mathbf{r}_a)\phi_1(\mathbf{r}_b) - \phi_0(\mathbf{r}_b)\phi_1(\mathbf{r}_a)] \end{aligned} \quad (4)$$

- (c) Electron a is spin-up in the ground state. Electrons b and c are spin-up and spin-down respectively in the first excited state:

$$\begin{aligned} \Psi(\gamma_a, \gamma_b, \gamma_c) &= \frac{1}{\sqrt{6}} \begin{vmatrix} \psi_a(\gamma_a) & \psi_a(\gamma_b) & \psi_a(\gamma_c) \\ \psi_b(\gamma_a) & \psi_b(\gamma_b) & \psi_b(\gamma_c) \\ \psi_c(\gamma_a) & \psi_c(\gamma_b) & \psi_c(\gamma_c) \end{vmatrix} \\ &= \frac{1}{\sqrt{6}} \left[\psi_a(\gamma_a) \begin{vmatrix} \psi_b(\gamma_b) & \psi_b(\gamma_c) \\ \psi_c(\gamma_b) & \psi_c(\gamma_c) \end{vmatrix} - \psi_a(\gamma_b) \begin{vmatrix} \psi_b(\gamma_a) & \psi_b(\gamma_c) \\ \psi_c(\gamma_a) & \psi_c(\gamma_c) \end{vmatrix} + \psi_a(\gamma_c) \begin{vmatrix} \psi_b(\gamma_a) & \psi_b(\gamma_b) \\ \psi_c(\gamma_a) & \psi_c(\gamma_b) \end{vmatrix} \right] \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{\sqrt{6}} \left\{ \phi_0(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) \begin{vmatrix} \phi_1(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) & \phi_1(\mathbf{r}_c)\chi_+(\boldsymbol{\sigma}_c) \\ \phi_1(\mathbf{r}_b)\chi_-(\boldsymbol{\sigma}_b) & \phi_1(\mathbf{r}_c)\chi_-(\boldsymbol{\sigma}_c) \end{vmatrix} \right. \\
&\quad - \phi_0(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) \begin{vmatrix} \phi_1(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) & \phi_1(\mathbf{r}_c)\chi_+(\boldsymbol{\sigma}_c) \\ \phi_1(\mathbf{r}_a)\chi_-(\boldsymbol{\sigma}_a) & \phi_1(\mathbf{r}_c)\chi_-(\boldsymbol{\sigma}_c) \end{vmatrix} \\
&\quad \left. - \phi_0(\mathbf{r}_c)\chi_+(\boldsymbol{\sigma}_c) \begin{vmatrix} \phi_1(\mathbf{r}_a)\chi_+(\boldsymbol{\sigma}_a) & \phi_1(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_b) \\ \phi_1(\mathbf{r}_a)\chi_-(\boldsymbol{\sigma}_a) & \phi_1(\mathbf{r}_b)\chi_-(\boldsymbol{\sigma}_b) \end{vmatrix} \right\} \\
&= \frac{1}{\sqrt{6}} \{ \phi_0(\mathbf{r}_a)\phi_1(\mathbf{r}_b)\phi_1(\mathbf{r}_c)\chi_+(\boldsymbol{\sigma}_a) [\chi_+(\boldsymbol{\sigma}_b)\chi_-(\boldsymbol{\sigma}_c) - \chi_+(\boldsymbol{\sigma}_c)\chi_-(\boldsymbol{\sigma}_b)] \\
&\quad - \phi_0(\mathbf{r}_b)\phi_1(\mathbf{r}_a)\phi_1(\mathbf{r}_c)\chi_+(\boldsymbol{\sigma}_b) [\chi_+(\boldsymbol{\sigma}_a)\chi_-(\boldsymbol{\sigma}_c) - \chi_+(\boldsymbol{\sigma}_c)\chi_-(\boldsymbol{\sigma}_a)] \\
&\quad - \phi_0(\mathbf{r}_c)\phi_1(\mathbf{r}_a)\phi_1(\mathbf{r}_b)\chi_+(\boldsymbol{\sigma}_c) [\chi_+(\boldsymbol{\sigma}_a)\chi_-(\boldsymbol{\sigma}_b) - \chi_+(\boldsymbol{\sigma}_b)\chi_-(\boldsymbol{\sigma}_a)] \} \tag{5}
\end{aligned}$$

2. Recall:

$$\Psi(\{\gamma_i\}) = \frac{1}{\sqrt{N!}} \sum_{k \in \mathbb{K}} \epsilon_k \prod_{j=1}^N \psi_{k_j}(\gamma_j) \tag{6}$$

$$\begin{aligned}
\langle \Psi(\{\gamma_i\}) | \Psi(\{\gamma_i\}) \rangle &= \frac{1}{N!} \sum_{k \in \mathbb{K}} \sum_{k' \in \mathbb{K}} \int d\gamma_1 \dots d\gamma_N \epsilon_k \epsilon_{k'} \prod_{j=1}^N [\psi_{k_j}^*(\gamma_j)] \prod_{j=1}^N [\psi_{k'_j}(\gamma_j)] \\
&= \frac{1}{N!} \sum_{k \in \mathbb{K}} \sum_{k' \in \mathbb{K}} \prod_{i=1}^N \left[\int d\gamma_i \psi_{k_i}^*(\gamma_i) \psi_{k'_i}(\gamma_i) \right] \epsilon_k \epsilon_{k'} \\
&= \frac{1}{N!} \sum_{k \in \mathbb{K}} \sum_{k' \in \mathbb{K}} \prod_{i=1}^N [\delta_{k_i, k'_i}] \epsilon_k \epsilon_{k'} \\
&= \frac{1}{N!} \sum_{k \in \mathbb{K}} \epsilon_k \epsilon_k \\
&= \frac{1}{N!} \sum_{k \in \mathbb{K}} 1 \\
&= \frac{1}{N!} N! \\
&= 1 \tag{7}
\end{aligned}$$