

## Problem 2: WKB approximation (10 Points):

The one-dimensional Schrodinger equation,

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi = E\psi$$

can be rewritten as

$$\frac{d^2\psi}{dx^2} = -\frac{p^2}{\hbar^2}\psi,$$

where

$$p(x) \equiv \sqrt{2m[E - V(x)]}.$$

The wave function  $\psi(x)$  is often expressed as  $\psi(x) = A(x)e^{i\phi(x)}$  where  $A(x)$  is the amplitude and  $\phi(x)$  is the phase. Both  $A(x)$  and  $\phi(x)$  can be real.

- (a) Show that the amplitude is  $A = \frac{C}{\sqrt{\phi'}}$  where  $C$  is a constant and prime is the derivative with respect to  $x$ . (2 points)
- (b) (3 points) Let us assume that  $A''/A \ll (\phi')^2$  and  $A''/A \ll p^2/\hbar^2$ . Show that the wave function in the WKB approximation is

$$\psi(x) \simeq \frac{C}{\sqrt{p(x)}} e^{\pm \frac{i}{\hbar} \int p(x) dx}.$$

In parts (c)–(e), the potential energy of the one-dimensional harmonic oscillator is

$$V(x) = \frac{1}{2}m\omega^2 x^2.$$

- (c) Find the classical turning points  $x_1 < x_2$  for an energy  $E$ . (1 points)
- (d) Evaluate the phase  $\phi$  in terms of  $E$  and  $\omega$  with the WKB method. (3 points)
- (e) Apply the eigenvalue condition  $\phi = (n + \frac{1}{2})\pi\hbar$  and find energy eigenvalues  $E_n$ . (1 points)

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## Quantum #2

a) To show  $A = \frac{C}{\sqrt{\varphi'}}$ , assuming  $\psi(x) = A(x)e^{i\varphi(x)}$ , we substitute  $\psi$  into the rewritten Schrödinger eqn:  $\frac{d^2\psi}{dx^2} = -\frac{p^2}{\hbar^2} \psi$ ,  $p = \sqrt{2m(E-V(x))}$

$$\hookrightarrow \frac{d\psi}{dx} = A'e^{i\varphi(x)} + Ai\varphi'e^{i\varphi(x)}$$

$$\frac{d^2\psi}{dx^2} = A''e^{i\varphi(x)} + iA'\varphi'e^{i\varphi(x)} + iA'\varphi'e^{i\varphi(x)} + iA\varphi''e^{i\varphi(x)} - A(\varphi')^2e^{i\varphi(x)}$$

$$e^{i\varphi(x)} \cdot [A'' - A(\varphi')^2 + i[2A'\varphi' + A\varphi'']] = -\frac{p^2}{\hbar^2} A(x)e^{i\varphi(x)}$$

$$\text{Real: } -\frac{p^2}{\hbar^2} A = A'' - A(\varphi')^2$$

$$\text{Imaginary: } 0 = 2A'\varphi' + A\varphi''$$

$$0 = \frac{d}{dx}(A^2\varphi')$$

$$\hookrightarrow C^2 = A^2\varphi'$$

$$\hookrightarrow A = \frac{C}{\sqrt{\varphi'}} \checkmark$$

$$\frac{d}{dx}(A^2\varphi') = 2AA'\varphi' + A^2\varphi'' = 0$$

\* divide by A

$$2A'\varphi' + A\varphi'' = 0$$

b) Assuming  $\frac{A''}{A} \ll (\varphi')^2$  and  $\frac{A''}{A} \ll \frac{p^2}{\hbar^2}$ , we can now use the real equation from part A to show  $\psi(x) \approx \frac{C}{\sqrt{p(x)}} e^{\pm \frac{i}{\hbar} \int p(x) dx}$

$$-\frac{p^2}{\hbar^2} A = A'' - A(\varphi')^2$$

$$-\frac{p^2}{\hbar^2} \frac{C}{\sqrt{\varphi'}} = \frac{C}{\sqrt{\varphi'}} (\varphi')^2$$

$$\varphi' = \pm \frac{p}{\hbar} \Rightarrow \varphi = \pm \frac{1}{\hbar} \int_0^x p(x) dx$$

$$\hookrightarrow \psi(x) = A e^{i\varphi(x)}$$

$$= \frac{C}{\sqrt{\varphi'}} \exp\left[\pm \frac{i}{\hbar} \int p(x) dx\right]$$

## #2 (cont.)

c) The classical turning points occur when  $E = V(x)$

$$\hookrightarrow E = \frac{1}{2} m \omega^2 x^2$$

$$\hookrightarrow x = \pm \sqrt{\frac{2E}{m}} \omega$$

d) Note: In the region where  $E < V$ ,  $p(x)$  is imaginary

$E > V$ ,  $p(x)$  is real

$E = V$   $p(x)$  is 0

\* If  $p(x) = 0$ ,  $\psi(x) = 0$

\* If  $p(x)$  is real ( $E > V$ )

$$\begin{aligned}\psi(x) &= \int_{x_1}^{x_2} p(x) dx \\&= \int_{x_1}^{x_2} \left[ 2m \left( E - \frac{1}{2} m \omega^2 x^2 \right) \right]^{1/2} dx \\&= \int_{x_1}^{x_2} \left( 2mE - m^2 \omega^2 x^2 \right)^{1/2} dx \\&\quad * \text{let } a = \sqrt{2mE}, \quad u = m\omega x \\&= \int_{x_1}^0 \sqrt{a^2 - u^2} \frac{du}{m\omega} + \int_0^{x_2} \sqrt{a^2 - u^2} \frac{du}{m\omega} \\&= \int_0^{x_2} \frac{1}{m\omega} \sqrt{a^2 - u^2} du - \int_0^{x_1} \frac{1}{m\omega} \sqrt{a^2 - u^2} du \\&= \frac{1}{m\omega} \left[ \frac{\pi x_2^2}{4} - \frac{\pi x_1^2}{4} \right] \\&= \frac{\pi}{4m\omega} [x_2^2 - x_1^2] \\&= \frac{\pi}{4m\omega^2}\end{aligned}$$

See Griffiths QM  
ex. 8.4