

### Problem 3: Identical particles (10 pts)

Two non-interacting particles of mass  $m$  are trapped in a 1-dimensional infinite box of length  $L$  situated between  $x = 0$  and  $x = L$ . (In the cases you are considering fermions, assume them to all be spin up.)

- (a) [1 points] Write down the single particle energy eigenvalues and wavefunctions.
- (b) [1 points] Write down the energy eigenvalues and wavefunctions for two distinguishable particles. Label the states by  $n_1$  for particle 1 and  $n_2$  for particle 2.
- (c) [2 points] An energy measurement of the *two identical particle* system yields  $E = \hbar^2\pi^2/mL^2$ . Write down the state vector/wave function of the system.
- (d) [2 points] Suppose instead the energy of the two identical particle system is measured to be  $E = 5\hbar^2\pi^2/mL^2$ . What is the wave function?  
*Hint: there are two possibilities.*
- (e) [2 points] Show that the fermion state you found in part (d) is an eigenfunction of the Hamiltonian, with the appropriate eigenvalue.
- (f) [1 points] Write down the wavefunction for two identical spin-up fermions in the  $n_1 = 2$  and  $n_2 = 2$  state.
- (g) [1 points] If instead you had three particles in the orthonormal states  $\Psi_1, \Psi_2$ , and  $\Psi_3$ , construct the three particle state for identical fermions.

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### Quantum #3

a) For an infinite well b/w 0 and L, our solution is:

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

for any single particle

b) Assuming the particles are distinguishable, we simply use the above equations

$$\psi_{n_1} = \sqrt{\frac{2}{L}} \sin\left(\frac{n_1 \pi x}{L}\right)$$

$$E_{n_1} = \frac{n_1^2 \pi^2 \hbar^2}{2mL^2}$$

$$\psi_{n_2} = \sqrt{\frac{2}{L}} \sin\left(\frac{n_2 \pi x}{L}\right)$$

$$E_{n_2} = \frac{n_2^2 \pi^2 \hbar^2}{2mL^2}$$

c) Considering two identical spin-up fermions, the exclusion principle prevents both particles from being in the same state. Since the only combination that yields  $E_{n_1, n_2} = \frac{(n_1^2 + n_2^2) \pi^2 \hbar^2}{2mL^2} = \frac{\pi^2 \hbar^2}{mL^2}$  is  $n_1 = n_2 = 1$ , this state is disallowed by the exclusion principle, therefore

$$\psi_{1,1} = 0$$

d) If  $E_{n_1, n_2} = \frac{5\pi^2 \hbar^2}{mL^2}$ , our possible configurations are  $n_1 = 1, n_2 = 3$ ;  $n_1 = 3, n_2 = 1$

Our general wavefunction for identical fermions is:

$$\psi_{n_1, n_2} = \frac{1}{\sqrt{2}} (\psi_{n_1}(x_1) \psi_{n_2}(x_2) - \psi_{n_1}(x_2) \psi_{n_2}(x_1))$$

This yields the following potential wave functions:

$$\begin{aligned} \psi_{13} &= \frac{1}{\sqrt{2}} (\psi_1(x_1) \psi_3(x_2) - \psi_1(x_2) \psi_3(x_1)) \\ &= \frac{\sqrt{2}}{L} \left( \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) - \sin\left(\frac{\pi x_2}{L}\right) \sin\left(\frac{3\pi x_1}{L}\right) \right) \end{aligned}$$

$$\psi_{31} = \frac{\sqrt{2}}{L} \left( \sin\left(\frac{3\pi x_1}{L}\right) \sin\left(\frac{\pi x_2}{L}\right) - \sin\left(\frac{3\pi x_2}{L}\right) \sin\left(\frac{\pi x_1}{L}\right) \right)$$

#3 (cont.)

e) We know that  $H \psi_n = E_n \psi_n$  where  $H = \frac{-\hbar^2}{2m} \left( \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} \right)$

$\Rightarrow$  For the  $\psi_{13}$  state:

$$\frac{\partial^2}{\partial x_1^2} \psi_{13} = \frac{\sqrt{2}}{L} \left( -\frac{\pi^2}{L^2} \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) + \left(\frac{3\pi}{L}\right)^2 \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) \right)$$

$$\frac{\partial^2}{\partial x_2^2} \psi_{13} = \frac{\sqrt{2}}{L} \left( -\left(\frac{3\pi}{L}\right)^2 \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) + \left(\frac{\pi}{L}\right)^2 \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) \right)$$

$$\begin{aligned} \hookrightarrow H \psi_{13} &= \frac{-\hbar^2}{2m} \left( \frac{\sqrt{2}}{L} \left[ -\frac{\pi^2}{L^2} \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) + \frac{9\pi^2}{L^2} \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) \right] \right. \\ &\quad \left. + \frac{\sqrt{2}}{L} \left[ -\frac{9\pi^2}{L^2} \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) + \frac{\pi^2}{L^2} \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) \right] \right) \\ &= \frac{\hbar^2 10\pi^2}{2m L^2} \left( \frac{\sqrt{2}}{L} \left[ \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) + \sin\left(\frac{\pi x_1}{L}\right) \sin\left(\frac{3\pi x_2}{L}\right) \right] \right) \\ &= \frac{10\hbar^2 \pi^2}{2m L^2} \psi_{13} \quad \checkmark \end{aligned}$$

\* A similar process will reach the same conclusion for the  $\psi_{31}$  state

f) The  $n_1 = n_2 = 2$  state is disallowed by the exclusion principle

$$\hookrightarrow \psi_{22} = 0$$

g) My guess is this follows something like a cyclic permutation

$$\begin{aligned} \Rightarrow \psi_{n_1 n_2 n_3} &= \frac{1}{\sqrt{6}} \left[ \psi_{n_1}(x_1) \psi_{n_2}(x_2) \psi_{n_3}(x_3) - \psi_{n_1}(x_2) \psi_{n_2}(x_3) \psi_{n_3}(x_1) \right. \\ &\quad - \psi_{n_1}(x_3) \psi_{n_2}(x_1) \psi_{n_3}(x_2) + \psi_{n_1}(x_3) \psi_{n_2}(x_2) \psi_{n_3}(x_1) \\ &\quad \left. + \psi_{n_1}(x_2) \psi_{n_2}(x_1) \psi_{n_3}(x_3) + \psi_{n_1}(x_1) \psi_{n_2}(x_3) \psi_{n_3}(x_2) \right] \end{aligned}$$