

### Problem 5: Interaction Picture of Quantum Mechanics

The “Interaction Picture” of quantum mechanics is in some ways in-between the Schrödinger formulation and the Heisenberg formulation.

Consider a system with the Hamiltonian  $H = H_0 + V(t)$  where  $H_0$  is independent of time and  $V(t)$  may or may not be time dependent. The Interaction Picture is defined by the transformation of the Schrödinger states:

$$\begin{aligned} |\psi\rangle_I &= U_0^{-1} |\psi\rangle_S \\ U_0 &= e^{-\frac{i}{\hbar}(t-t_0)H_0}. \end{aligned} \quad (1)$$

The subscripts  $I$  and  $S$  refer to the Interaction Picture and Schrödinger Picture respectively.  $t_0$  is a time when the pictures coincide, and we will set  $t_0 = 0$  for this problem.

- (a) [1 pt] Show that  $U_0$  is a unitary operator. Why is it important for the transformation between pictures be unitary?
- (b) [3 pts] The transformation between  $|\psi\rangle_S$  and  $|\psi\rangle_I$  implies that there is also a transformation of the observables between the pictures. If  $A_S$  and  $A_I$  are operators for an observable in the Schrödinger and Interaction pictures respectively, derive the relation between  $A_S$  and  $A_I$ . Show that this implies that  $H_0$  is the same in the two pictures.
- (c) [3 pts] Derive the differential equation that determines the time dependence of the Interaction Picture states,  $|\psi(t)\rangle_I$ . Be sure to show and explain your work. Explain why the Interaction Picture may be particularly useful when  $V(t)$  is “small”.
- (d) [1 pt] Define the eigenstates of  $H_0$  to be

$$H_0|\lambda\rangle_S = E_\lambda|\lambda\rangle_S \quad (2)$$

Show that if  $V(t) = 0$ , the Interaction Picture energy eigenstates  $|\lambda\rangle_I$  are equal to  $|\lambda(t=0)\rangle_S$  and independent of time.

- (e) [2 pts] Consider a potential of the form

$$V(t) = 0, \quad t \leq 0 \quad V(t) \neq 0, \quad t > 0 \quad (3)$$

The system is in a state  $|\psi_0\rangle_I$  for  $t < 0$ . For  $t > 0$  the Interaction Picture state will depend on time. It can be expanded as:

$$|\psi(t)\rangle_I = \sum_{\lambda} c_{\lambda}(t) |\lambda(0)\rangle_I \quad (4)$$

In this expression,  $c_{\lambda}(t)$  are time-dependent expansion coefficients for the state and  $|\lambda(0)\rangle_I$  is the complete set of time-independent eigenstates of  $H_0$  in the interaction picture.

Use the time dependence found in part (c) to derive a set of coupled equations relating  $c_{\lambda}(t)$  and  $\partial_t c_{\lambda}(t)$ .

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## Quantum #5

a) An operator is unitary if:  $U^\dagger U = \mathbb{I}$

$$\hookrightarrow U = \exp\left[-\frac{i(t-t_0)H_0}{\hbar}\right]$$

$$U^\dagger = \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right]$$

$$U^\dagger U = \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] \exp\left[-\frac{i(t-t_0)H_0}{\hbar}\right]$$

$$= 1 \quad \checkmark$$

Our transformation must be a unitary operator b/c it preserves lengths of vectors, as well as the angles b/w them

b) This transformation is best seen in terms of an expectation value. In the Schrödinger picture:

$$\langle A \rangle = \langle \psi | A_S | \psi \rangle_S$$

since we can only multiply by 1 and we want  $|\psi\rangle_I = U_0^{-1} |\psi\rangle_S$

$$= \langle \psi | U_0 U_0^{-1} A_S U_0 U_0^{-1} | \psi \rangle_S$$

$$= \langle \psi | U_0^{-1} A_S U_0 | \psi \rangle_I$$

$$\hookrightarrow U_0^{-1} A_S U_0 = A_I$$

\*Not sure I've fully answered this part

$$c) |\psi\rangle_I = \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] |\psi\rangle_S$$

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle_I = i\hbar \frac{\partial}{\partial t} \left( \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] |\psi\rangle_S \right)$$

$$= i\hbar \left( \frac{H_0}{i\hbar} \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] |\psi\rangle_S + \left[ \frac{i(t-t_0)H_0}{\hbar} \right] \frac{\partial}{\partial t} |\psi\rangle_S \right)$$

$$= -H_0 \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] |\psi\rangle_S + (H_0 + V) |\psi\rangle_S \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right]$$

$$= V \exp\left[\frac{i(t-t_0)H_0}{\hbar}\right] |\psi\rangle_S$$

$$= V |\psi\rangle_I$$

#5 (cont.)

$$c) i\hbar \frac{\partial}{\partial t} \langle n | \psi \rangle_I = \langle n | V | \psi \rangle_I$$

$$i\hbar \frac{\partial}{\partial t} c_n(t) = \sum_m \langle n | V | m \rangle \langle m | \psi \rangle_I$$

$$= \sum_m c_m V_{nm}$$

expanding  $c_n(t)$  as  $c_n^{(0)}(t) + \lambda c_n^{(1)}(t) + \dots$

$$i\hbar \frac{\partial}{\partial t} (c_n^{(0)}(t) + \lambda c_n^{(1)}(t) + \dots) = \sum_m (c_m^{(0)}(t) + \lambda c_m^{(1)}(t))$$