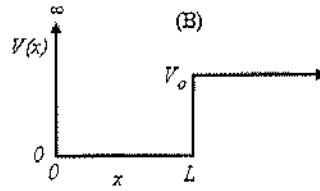
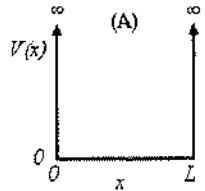


PROBLEM 2



- a) [2 pts] Calculate the energy eigenvalues for a particle of mass m in the one-dimensional infinite well shown in Figure A.
- b) [4 pts] For the time-independent Schrödinger Equation corresponding to potential (B), find a transcendental equation in E giving the eigenenergies in terms of V_0 , L , m , and \hbar .
- c) [4 pts] For the time-independent Schrödinger Equation corresponding to potential (B), what is the smallest value of V_0 that gives one bound state? What is the smallest value of V_0 that gives two bound states?

(a)



The Schrödinger equation tells us

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi$$

$$\frac{d^2 \psi}{dx^2} = -k^2 \psi,$$

where

$$k = \frac{\sqrt{2mE}}{\hbar},$$

for this potential. The general solution is

$$\psi(x) = A \sin(kx) + B \cos(kx).$$

Boundary conditions require that

$$\psi(0) = \psi(L) = 0,$$

so

$$\psi(0) = A \sin(0) + B \cos(0) = 0$$

implies $B=0$. Then

$$\psi(L) = A \sin(kL) = 0$$

implies

$$kL = n\pi$$

$$k = \frac{n\pi}{L}. \quad (n=1,2,3,\dots)$$

(a), cont'd...

But we said

$$k = \frac{\sqrt{2mE}}{\hbar},$$

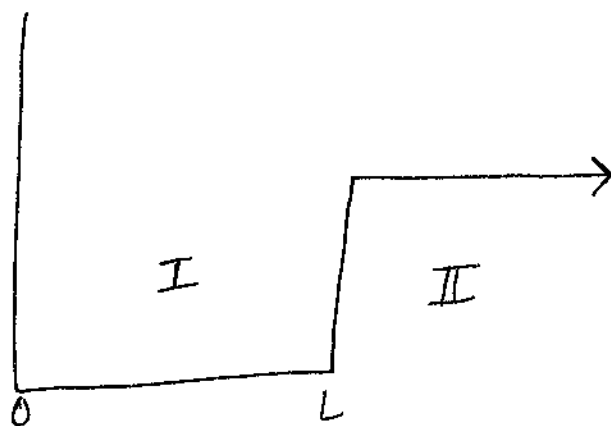
so our energy eigenvalues are

$$\frac{\sqrt{2mE}}{\hbar} = \frac{n\pi}{L}$$

$$2mE = \frac{n^2 \pi^2 \hbar^2}{L^2}$$

$$E = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

(b)



In region I, we have

$$\psi_I(x) = A \sin(kx) + B \cos(kx)$$

and in region II, we have

$$\psi_{II}(x) = D e^{\lambda x} + F e^{-\lambda x}$$

Our boundary conditions in region I still require that we have $B=0$, so

$$\psi_I = A \sin(kx) = \sqrt{\frac{2}{L}} \sin(kx)$$

As $x \rightarrow \infty$, we must have $D=0$ because the solution blows up. So

$$\psi_{II}(x) = F e^{-\lambda x}$$

Since the Schrödinger equation tells us

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + (V_0 - E) \psi = 0$$

$$\frac{d^2 \psi}{dx^2} = \frac{2m(V_0 - E)}{\hbar^2} \psi,$$

we know

$$\lambda = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$$

(b), cont'd...

Continuity requires that

$$\psi_{\text{I}}(L) = \psi_{\text{II}}(L)$$

and

$$\left. \frac{d\psi_{\text{I}}}{dx} \right|_L = \left. \frac{d\psi_{\text{II}}}{dx} \right|_L.$$

So

$$\begin{aligned} \sqrt{\frac{2}{L}} \sin(kL) &= F e^{-\lambda L} \\ F &= \sqrt{\frac{2}{L}} \sin(kL) e^{\lambda L} \end{aligned}$$

and

$$\begin{aligned} k \sqrt{\frac{2}{L}} \cos(kL) &= -\lambda F e^{-\lambda L} \\ k \sqrt{\frac{2}{L}} \cos(kL) &= -\lambda \sqrt{\frac{2}{L}} \sin(kL) \\ -\frac{k}{\lambda} &= \tan(kL). \end{aligned}$$

But

$$k = \frac{\sqrt{2mE}}{\hbar}$$

and

$$\lambda = \frac{\sqrt{2m(V_0 - E)}}{\hbar},$$

so

$$\begin{aligned} \sqrt{\frac{2mE}{2m(V_0 - E)}} &= -\tan\left(\frac{\sqrt{2mE}}{\hbar} L\right) \\ \sqrt{\frac{V_0 - E}{E}} &= -\cot\left(\frac{\sqrt{2mE}}{\hbar} L\right) \end{aligned}$$

(b), cont'd...

We could have left our transcendental equation in the form

$$k \cot(kL) = -\lambda$$

$$\frac{\sqrt{2mE}}{\hbar} \cot\left(\frac{\sqrt{2mE}}{\hbar} L\right) = -\frac{\sqrt{2m(V_0 - E)}}{\hbar},$$

as well.