

PROBLEM 3: Harmonic Oscillator in 3 Dimensions

Consider a particle subject to a 3-dimensional harmonic oscillator potential:

$$\begin{aligned} H &= H_x + H_y + H_z \\ &= \frac{P_x^2}{2m} + \frac{1}{2}m\omega_x^2 X^2 + \frac{P_y^2}{2m} + \frac{1}{2}m\omega_y^2 Y^2 + \frac{P_z^2}{2m} + \frac{1}{2}m\omega_z^2 Z^2 \end{aligned}$$

where $\omega_x = \omega_y = \omega$ and $\omega_z = 2\omega$. The wave function is given by:

$$\Psi_{n_x, n_y, n_z}(x, y, z) = \Phi_{n_x}(x)\Phi_{n_y}(y)\Phi_{n_z}(z)$$

- (a). For the ground state wave function, determine, ΔX , ΔP_x and their product $\Delta X \Delta P_x$ using Dirac notation and raising and lowering operators. (3 points)
- (b). For the ground state, $\Psi_{0,0,0}$, do you expect
- (i) ΔZ to be larger or smaller than ΔX ?
 - (ii) $\Delta Z \Delta P_z$ to be larger or smaller than $\Delta X \Delta P_x$?
- Explain your reasoning in each case. (2 points)
- (c). Assume that at $t = 0$ the particle is in the state:

$$|\Psi(t=0)\rangle = \frac{1}{\sqrt{6}}|\Psi_{1,0,0}\rangle + \frac{1}{\sqrt{3}}|\Psi_{2,1,0}\rangle + \frac{1}{2}|\Psi_{0,1,0}\rangle + \frac{1}{2}|\Psi_{1,0,1}\rangle$$

If one measures the total energy, E , what is the probability of obtaining $5\hbar\omega$? (2 points)

- (d). Immediately after the measurement performed in part (c), what harmonic oscillator state (or superposition of states) is the system in? (1 points)
- (e). Assume that at time $t = 0$ the measurement described in part (c) is performed and that the energy is found to be $E = 5\hbar\omega$. If the observable, X , is measured at a time $t > 0$, will its probability distribution be time dependent or time independent? Explain your reasoning. (2 points)

(a)

Let's represent our state in Dirac notation,

$$|n\rangle = |n_x n_y n_z\rangle.$$

We know, in general, that

$$X = \sqrt{\frac{\hbar}{2m\omega_x}} (a_x + a_x^\dagger)$$

and

$$P_x = -i\sqrt{\frac{m\hbar\omega_x}{2}} (a_x - a_x^\dagger).$$

So

$$\begin{aligned}\langle X \rangle_n &= \langle n | X | n \rangle \\ &= \langle n | \sqrt{\frac{\hbar}{2m\omega_x}} (a_x + a_x^\dagger) | n \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega_x}} [\langle n | a_x | n \rangle + \langle n | a_x^\dagger | n \rangle]\end{aligned}$$

In the ground state, we have

$$|n\rangle = |0\rangle = |0\ 0\ 0\rangle.$$

So

$$\begin{aligned}\langle X \rangle_0 &= \sqrt{\frac{\hbar}{2m\omega_x}} [0 + \langle 0 | 100 \rangle] \\ &= 0.\end{aligned}$$

(a), cont'd...

We also have

$$\begin{aligned}\langle X^2 \rangle_0 &= \frac{\hbar}{2m\omega_x} \left[\langle 000 | a_x a_x + a_x a_x^\dagger + a_x^\dagger a_x + a_x^\dagger a_x^\dagger | 000 \rangle \right] \\ &= \frac{\hbar}{2m\omega_x} \left[0 + 1 \langle 000 | 000 \rangle + 0 + \sqrt{2} \langle 000 | 200 \rangle \right] \\ &= \frac{\hbar}{2m\omega}\end{aligned}$$

Thus,

$$\begin{aligned}\Delta X_0 &= \sqrt{\langle X^2 \rangle_0 - \langle X \rangle_0^2} \\ \Delta X_0 &= \left(\frac{\hbar}{2m\omega} \right)^{1/2}\end{aligned}$$

We have

$$\begin{aligned}\langle P_x \rangle_0 &= \langle 0 | P_x | 0 \rangle \\ &= \langle 0 | -i \sqrt{\frac{m\hbar\omega_x}{2}} (a_x - a_x^\dagger) | 0 \rangle \\ &= 0\end{aligned}$$

and

$$\begin{aligned}\langle P_x^2 \rangle_0 &= \langle 0 | P_x^2 | 0 \rangle \\ &= \langle 0 | -\frac{m\hbar\omega_x}{2} (a_x - a_x^\dagger)^2 | 0 \rangle \\ &= -\frac{m\hbar\omega_x}{2} \langle 0 | a_x a_x - a_x a_x^\dagger - a_x^\dagger a_x + a_x^\dagger a_x^\dagger | 0 \rangle \\ &= -\frac{m\hbar\omega_x}{2} \left[0 - 1 \langle 0 | 0 \rangle - 0 + \sqrt{2} \langle 0 | 2 \rangle \right]\end{aligned}$$

(a), cont'd...

$$\langle p_x^2 \rangle_0 = \frac{m\hbar\omega}{2}$$

Then

$$\Delta p_x = \sqrt{\langle p_x^2 \rangle - \langle p_x \rangle^2}$$

$$\Delta p_x = \left(\frac{m\hbar\omega}{2} \right)^{1/2}$$

So

$$\Delta x \Delta p_x = \left(\frac{\hbar}{2m\omega} \right)^{1/2} \left(\frac{m\hbar\omega}{2} \right)^{1/2}$$

$$\Delta x \Delta p_x = \frac{\hbar}{2}$$

(b)

(i) I would expect Δz to be smaller than Δx since

$$\Delta z \propto \left(\frac{1}{2\omega} \right)^{1/2},$$

whereas

$$\Delta x \propto \left(\frac{1}{\omega} \right)^{1/2}.$$

(ii) I would expect $\Delta z \Delta p_z = \Delta x \Delta p_x$ because

$$\Delta z \propto \left(\frac{1}{2\omega} \right)^{1/2}$$

$$\Delta p_z \propto (2\omega)^{1/2},$$

so the result is independent of ω_z .

(c)

If we have

$$E = 5\hbar\omega,$$

we must have

or
since

$$\begin{aligned} n_x &= 2, n_y = 1, n_z = 0 \\ n_x &= 1, n_y = 0, n_z = 1 \end{aligned}$$

$$E = (n_x + n_y + 2n_z + 2)\hbar\omega.$$

Then

$$\begin{aligned} P(E=5\hbar\omega) &= \left| \frac{1}{\sqrt{3}} \right|^2 + \left| \frac{1}{2} \right|^2 \\ &= \frac{1}{3} + \frac{1}{4} \\ &= \frac{4}{12} + \frac{3}{12} \end{aligned}$$

$$P(E=5\hbar\omega) = \frac{7}{12}.$$

(d) Immediately after the measurement in part (c), our state is

$$|\psi'\rangle = \frac{1}{\sqrt{3}}|\psi_{210}\rangle + \frac{1}{2}|\psi_{101}\rangle.$$

Normalizing...

$$A^2 \left(\frac{1}{3} + \frac{1}{4} \right) = 1$$

$$A^2 \left(\frac{7}{12} \right) = 1$$

$$A = \sqrt{\frac{12}{7}}$$

So

$$|\psi'\rangle = \sqrt{\frac{4}{7}}|\psi_{210}\rangle + \sqrt{\frac{3}{7}}|\psi_{101}\rangle.$$

(e)

Since we collapsed our wavefunction into a superposition of states with the same eigenenergy, our wavefunction after the measurement represents a stationary state. Thus, I would expect the probability distribution of the observable X to be time-independent.

Mathematically, we have

$$\begin{aligned}\psi(x,t) &= U(t=0) \psi(t=0) \\ &= \sqrt{\frac{4}{7}} |\psi_{210}\rangle e^{-i(5\hbar\omega)t/\hbar} + \sqrt{\frac{3}{7}} |\psi_{101}\rangle e^{-i(5\hbar\omega)t/\hbar} \\ &= \left[\sqrt{\frac{4}{7}} |\psi_{210}\rangle + \sqrt{\frac{3}{7}} |\psi_{101}\rangle \right] e^{-i(5\hbar\omega)t/\hbar}\end{aligned}$$

So

$$\begin{aligned}\langle X \rangle &= \langle \psi(x,t) | X | \psi(x,t) \rangle \\ &= e^{i(5\hbar\omega)t/\hbar} \left[\sqrt{\frac{4}{7}} \langle \psi_{210} | + \sqrt{\frac{3}{7}} \langle \psi_{101} | \right] X \left[\sqrt{\frac{4}{7}} |\psi_{210}\rangle + \sqrt{\frac{3}{7}} |\psi_{101}\rangle \right] e^{-i(5\hbar\omega)t/\hbar} \\ &= \left(\sqrt{\frac{4}{7}} \langle \psi_{210} | + \sqrt{\frac{3}{7}} \langle \psi_{101} | \right) X \left(\sqrt{\frac{4}{7}} |\psi_{210}\rangle + \sqrt{\frac{3}{7}} |\psi_{101}\rangle \right)\end{aligned}$$

which is obviously time-independent.