

PROBLEM 3

A particle of mass m has a potential energy represented by two one-dimensional harmonic oscillator potentials centered at $\pm a$:

$$V(x) = \frac{1}{2}K(x-a)^2 + \frac{1}{2}K(x+a)^2$$

- ~~[a]~~ (3 pts) What are the eigenvalues of the particle given this potential V ? You may derive this result from first principles or deduce the result from the well known eigenvalues of a particle moving in a single harmonic-oscillator potential.
- ~~[b]~~ (3 pts) The normalized ground-state eigenfunction of the particle is given by

$$\phi(x) = \frac{1}{\pi^{1/4}\Delta^{1/2}} \exp\left(-\frac{x^2}{2\Delta^2}\right)$$

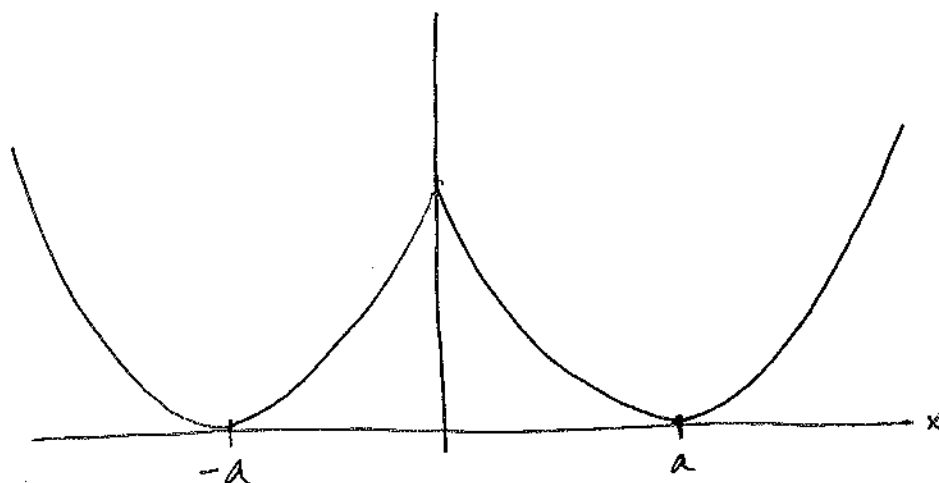
Use Schrodinger's equation to determine the constant Δ in terms of K , m , and fundamental constants.

- [c] (4 pt) The potential well at $x = -a$ suddenly disappears, leaving the particle in a new potential

$$U(x) = \frac{1}{2}K(x-a)^2$$

Suppose that *before the sudden change*, the particle was in the ground state of the double-well potential $V(x)$. Derive an expression for the probability that after the sudden change the particle will be in the ground state of the single well potential $U(x)$. Express your answer in terms of a and Δ .

(a)



We have

$$\begin{aligned} V(x) &= \frac{1}{2} K (x-a)^2 + \frac{1}{2} K (x+a)^2 \\ &= \frac{1}{2} K [(x-a)^2 + (x+a)^2] \\ &= \frac{1}{2} K (x^2 + 2a^2) \\ &= m\omega^2 (x^2 + a^2). \end{aligned}$$

So our Hamiltonian is

$$H = \frac{p^2}{2m} + m\omega^2 (x^2 + a^2).$$

We know

$$\begin{aligned} a^+ &= \sqrt{\frac{m\omega}{2\hbar}} x - \frac{i}{\sqrt{2\hbar m\omega}} p \\ a &= \sqrt{\frac{m\omega}{2\hbar}} x + \frac{i}{\sqrt{2\hbar m\omega}} p, \end{aligned}$$

so

$$\begin{aligned} p &= -i\sqrt{\frac{\hbar m\omega}{2}} (a - a^+) \\ x &= \sqrt{\frac{\hbar}{2m\omega}} (a + a^+) \end{aligned}$$

(a), cont'd...

To avoid confusion with the operators, let $x = a = \alpha$.
Plugging in p and x , we have

$$\begin{aligned} H &= \frac{1}{2m} \left(-i\sqrt{\frac{m\omega}{2}} (a - a^\dagger) \right)^2 + m\omega^2 \left[\left(\sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger) \right)^2 + \alpha^2 \right] \\ &= -\frac{\hbar\omega}{4} (a - a^\dagger)^2 + \frac{\hbar\omega}{2} (a + a^\dagger)^2 + m\omega^2 \alpha^2 \\ &= -\frac{\hbar\omega}{4} (aa - aa^\dagger - a^\dagger a + a^\dagger a^\dagger) + \frac{\hbar\omega}{2} (aa + aa^\dagger + a^\dagger a + a^\dagger a^\dagger) + m\omega^2 \alpha^2 \\ &= \frac{\hbar\omega}{4} (aa + a^\dagger a^\dagger) + \frac{3\hbar\omega}{4} (aa^\dagger + a^\dagger a) + m\omega^2 \alpha^2. \end{aligned}$$

But

$$\begin{aligned} [a, a^\dagger] &= 1 \\ aa^\dagger - a^\dagger a &= 1 \\ aa^\dagger &= a^\dagger a + 1 \end{aligned}$$

So

$$\begin{aligned} H &= \frac{\hbar\omega}{4} (aa + a^\dagger a^\dagger) + \frac{3\hbar\omega}{4} (2a^\dagger a + 1) + m\omega^2 \alpha^2 \\ H &= \frac{\hbar\omega}{4} (aa + a^\dagger a^\dagger) + \frac{3\hbar\omega}{2} (a^\dagger a + \frac{1}{2}) + m\omega^2 \alpha^2. \end{aligned}$$

Then

$$\begin{aligned} H|n\rangle &= \frac{\hbar\omega}{4} [\sqrt{n(n-1)} |n-2\rangle + \sqrt{(n+1)(n+2)} |n+2\rangle] \\ &\quad + \frac{3\hbar\omega}{2} [n |n\rangle + \frac{1}{2} |n\rangle] + m\omega^2 \alpha^2 |n\rangle \end{aligned}$$

and

$$\begin{aligned} \langle n|H|n\rangle &= \frac{\hbar\omega}{4} [\sqrt{n(n-1)} \langle n|n-2\rangle + \sqrt{(n+1)(n+2)} \langle n|n+2\rangle] \\ &\quad + \frac{3\hbar\omega}{2} [n \langle n|n\rangle + \frac{1}{2} \langle n|n\rangle] + m\omega^2 \alpha^2 \langle n|n\rangle \end{aligned}$$

(a), cont'd...

$$\langle n|H|n \rangle = \frac{3\hbar\omega}{2} \left(n + \frac{1}{2} \right) + m\omega^2 a^2$$

Thus, our energy eigenvalues are

$$E = \frac{3\hbar\omega}{2} \left(n + \frac{1}{2} \right) + m\omega^2 a^2$$

(b)

We know

$$a|0\rangle = 0.$$

So

$$\left(\sqrt{\frac{m\omega}{2\hbar}} x + \frac{i}{\sqrt{2\hbar m\omega}} p \right) \psi_0 = 0$$

$$\left(\sqrt{\frac{m\omega}{2\hbar}} x + \frac{i}{\sqrt{2\hbar m\omega}} \left(-i\hbar \frac{d}{dx} \right) \right) \psi_0 = 0$$

$$\sqrt{\frac{\hbar}{2m\omega}} \frac{d\psi_0}{dx} = - \sqrt{\frac{m\omega}{2\hbar}} x \psi_0$$

$$\frac{d\psi_0}{dx} = - \sqrt{\frac{m\omega}{2\hbar}} \sqrt{\frac{2m\omega}{\hbar}} x \psi_0$$

$$\frac{d\psi_0}{dx} = - \frac{m\omega}{\hbar} x \psi_0$$

and

$$\int \frac{1}{\psi_0} d\psi_0 = \int -\frac{m\omega}{\hbar} x dx$$

$$\ln \psi_0 = -\frac{m\omega}{2\hbar} x^2 + C_1$$

$$\psi_0 = C e^{-\frac{m\omega}{2\hbar} x^2}$$

Normalizing...

$$\int_{-\infty}^{\infty} |\psi_0|^2 dx = 1$$

$$C^2 \int_{-\infty}^{\infty} e^{-\frac{m\omega}{\hbar} x^2} dx = 1$$

$$C^2 \left(\sqrt{\pi \frac{\hbar}{m\omega}} \right) = 1$$

$$C = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4}$$

(b), cont'd...

Then

$$\psi_0(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$$

We are given

$$\phi(x) = \frac{1}{\pi^{1/4} \Delta^{1/2}} \exp \left(-\frac{x^2}{2\Delta^2} \right).$$

So

$$\left(\frac{m\omega}{\pi\hbar} \right)^{1/4} = \frac{1}{\pi^{1/4} \Delta^{1/2}}$$

$$\Delta^{1/2} = \frac{1}{\pi^{1/4}} \left(\frac{\pi\hbar}{m\omega} \right)^{1/4}$$

$$= \left(\frac{\hbar}{m\omega} \right)^{1/4}$$

$$\Delta = \left(\frac{\hbar}{m\omega} \right)^{1/2}.$$

This would give us

$$\phi(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} \exp \left(-\frac{m\omega}{2\hbar} x^2 \right),$$

which is what we expect. Thus,

$$\Delta = \left(\frac{\hbar}{m\omega} \right)^{1/2}$$

but $K = \omega^2 m$, so

$$\frac{\hbar}{m\omega} = \frac{\hbar}{m} \cdot \frac{1}{\sqrt{K/m}} = \hbar / \sqrt{mK}$$

and

$$\Delta = \left(\frac{\hbar^2}{mK} \right)^{1/4}$$

(c)

Let ψ_o^i be the wavefunction before the change, in the ground state, and let ψ_o^f be the ground state wave function after the change.

Then, in general, our probability that the particle will be in the ground state after the change is

$$p = \left| \int_{-\infty}^{\infty} \psi_o^{*f} \cdot \psi_o^i dx \right|^2.$$