

PROBLEM 2: Oscillator Model of Angular Momentum

Arbitrary angular momentum can be constructed from spin-1/2. The latter can be described in terms of the Pauli matrices

$$\mathbf{S} = \frac{\hbar}{2} \boldsymbol{\sigma}.$$

The construction of a general angular momentum can be done by introducing two sets of independent harmonic oscillators, in terms of creation (a_{ζ}^{\dagger}) and annihilation (a_{ζ}) operators,

$$[a_{+}, a_{-}] = 0, \quad [a_{+}^{\dagger}, a_{-}^{\dagger}] = 0, \quad [a_{\zeta}, a_{\zeta'}^{\dagger}] = \delta_{\zeta, \zeta'},$$

with $\zeta, \zeta' = \pm$ indexing oscillators of type \pm . Now define

$$\mathbf{J} = \frac{\hbar}{2} a^{\dagger} \boldsymbol{\sigma} a,$$

where a is a two component operator,

$$a = \begin{pmatrix} a_{+} \\ a_{-} \end{pmatrix}.$$

- a) Given the form of the Pauli matrices, give the explicit form for J_x , J_y , J_z in terms of a_{ζ}^{\dagger} and a_{ζ} operators (2 Points).
- b) Show that $J_{\pm} = J_x \pm iJ_y$ have particularly simple forms in terms of a_{ζ} and a_{ζ}^{\dagger} operators (1 Point).
- c) Compute the commutator $[J_x, J_y]$. How is this generalized for the other components? (2 Points)

- d) Show that

$$J^2 = J_z^2 + J_{+}J_{-} + i[J_x, J_y],$$

and then write this in terms of the number operators for the two harmonic oscillators,

$$n_{+} = a_{+}^{\dagger} a_{+}, \quad n_{-} = a_{-}^{\dagger} a_{-}.$$

Show that this implies that the eigenvalues of J^2 are $j(j+1)\hbar^2$, where j is an integer or an integer plus $\frac{1}{2}$ (Hint: apply the J^2 operator in the two harmonic oscillator state $|n_{+}, n_{-}\rangle$) (3 Points).

- e) Using the properties of the harmonic oscillators, show that the state in which J^2 has the eigenvalue $j(j+1)\hbar$ and $J_z = m\hbar$ can be constructed from the state in which both n_{+} and n_{-} have the value zero, $|0\rangle$, by

$$|jm\rangle = \frac{(a_{+}^{\dagger})^{j+m}}{\sqrt{(j+m)!}} \frac{(a_{-}^{\dagger})^{j-m}}{\sqrt{(j-m)!}} |0\rangle.$$

(2 Points)

(a)

In general, we have

$$\vec{J} = \frac{\hbar}{2} a^\dagger \vec{\sigma} a,$$

$$(\vec{J} = a^\dagger \vec{S} a)$$

where

$$a = \begin{pmatrix} a_+ \\ a_- \end{pmatrix}.$$

We know

$$\vec{J} = J_x + J_y + J_z$$

and

$$\vec{\sigma} = \sigma_x + \sigma_y + \sigma_z,$$

so we must have

$$\begin{aligned} J_x &= \frac{\hbar}{2} (a_+^\dagger \ a_-^\dagger) \sigma_x \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger \ a_-^\dagger) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger \ a_-^\dagger) \begin{pmatrix} a_- \\ a_+ \end{pmatrix} \end{aligned}$$

$$\boxed{J_x = \frac{\hbar}{2} (a_+^\dagger a_- + a_-^\dagger a_+)}$$

(a), cont'd...

$$\begin{aligned} J_y &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \sigma_y \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \begin{pmatrix} -ia_- \\ ia_+ \end{pmatrix} \end{aligned}$$

$$\boxed{J_y = -\frac{i\hbar}{2} (a_+^\dagger a_- - a_-^\dagger a_+)}$$

$$\begin{aligned} J_z &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \sigma_z \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a_+ \\ a_- \end{pmatrix} \\ &= \frac{\hbar}{2} (a_+^\dagger a_-^\dagger) \begin{pmatrix} a_+ \\ -a_- \end{pmatrix} \end{aligned}$$

$$\boxed{J_z = \frac{\hbar}{2} (a_+^\dagger a_+ - a_-^\dagger a_-)}$$

(b)

Now we want to show that $J_{\pm} = J_x \pm iJ_y$ can be written in simple form.

$$\begin{aligned} J_+ &= J_x + iJ_y \\ &= \frac{\hbar}{2} (a_+^\dagger a_- + a_-^\dagger a_+) + i \cdot \frac{-i\hbar}{2} (a_+^\dagger a_- - a_-^\dagger a_+) \\ &= \frac{\hbar}{2} [a_+^\dagger a_- + \cancel{a_-^\dagger a_+} + a_+^\dagger a_- - \cancel{a_-^\dagger a_+}] \\ &\boxed{J_+ = \hbar (a_+^\dagger a_-)} \end{aligned}$$

and

$$\begin{aligned} J_- &= J_x - iJ_y \\ &= \frac{\hbar}{2} (a_+^\dagger a_- + a_-^\dagger a_+) - i \cdot \frac{-i\hbar}{2} (a_+^\dagger a_- - a_-^\dagger a_+) \\ &= \frac{\hbar}{2} (\cancel{a_+^\dagger a_-} + a_-^\dagger a_+ - \cancel{a_+^\dagger a_-} + a_-^\dagger a_+) \\ &\boxed{J_- = \hbar (a_-^\dagger a_+)} \end{aligned}$$

(c)

Now we want to compute $[J_x, J_y]$. We have

$$\begin{aligned} [J_x, J_y] &= J_x J_y - J_y J_x \\ &= \left(\frac{\hbar}{2}\right)^2 [(a_+^\dagger a_- + a_-^\dagger a_+)(i a_-^\dagger a_+ - i a_+^\dagger a_-)] \\ &\quad - \left(\frac{\hbar}{2}\right)^2 [(i a_-^\dagger a_+ - i a_+^\dagger a_-)(a_+^\dagger a_- + a_-^\dagger a_+)] \\ &= \frac{i\hbar^2}{4} [a_+^\dagger a_- a_-^\dagger a_+ - a_+^\dagger a_- a_+^\dagger a_- + a_-^\dagger a_+ a_+^\dagger a_- - a_-^\dagger a_+ a_-^\dagger a_+] \\ &\quad - \frac{i\hbar^2}{4} [a_-^\dagger a_+ a_+^\dagger a_- + a_-^\dagger a_+ a_-^\dagger a_+ - a_+^\dagger a_- a_+^\dagger a_- - a_+^\dagger a_- a_-^\dagger a_+] \end{aligned}$$

We know $[a_+, a_+^\dagger] = 1$ and $[a_-, a_-^\dagger] = 1$. We also know $[a_+, a_-^\dagger] = 0$ and $[a_-, a_+^\dagger] = 0$. Lastly, we can say $[a_+^\dagger, a_-^\dagger] = 0$ and $[a_+, a_-] = 0$. This implies

$$a_+ a_+^\dagger - a_+^\dagger a_+ = 1 \quad (1)$$

$$a_- a_-^\dagger - a_-^\dagger a_- = 1 \quad (2)$$

$$a_+ a_-^\dagger = a_-^\dagger a_+ \quad (3)$$

$$a_- a_+^\dagger = a_+^\dagger a_- \quad (4)$$

$$a_+^\dagger a_-^\dagger = a_-^\dagger a_+^\dagger \quad (5)$$

$$a_+ a_- = a_- a_+ \quad (6)$$

Simplifying our original expression,

$$\begin{aligned} [J_x, J_y] &= \frac{i\hbar^2}{4} [2a_+^\dagger a_- a_-^\dagger a_+ + 0 - 2a_-^\dagger a_+ a_+^\dagger a_-] \\ &= \frac{i\hbar^2}{2} (a_+^\dagger a_- a_-^\dagger a_+ - a_-^\dagger a_+ a_+^\dagger a_-) \end{aligned}$$

(c), cont'd...

Using (1), (2), (5), and (6), we have

$$\begin{aligned} [J_x, J_y] &= \frac{i\hbar^2}{2} [a_+^\dagger a_+ + a_-^\dagger a_+^\dagger a_- a_+ - a_-^\dagger a_- - a_-^\dagger a_+^\dagger a_- a_+] \\ &= \frac{i\hbar^2}{2} [a_+^\dagger a_+ + a_-^\dagger a_+^\dagger a_- a_+ - a_-^\dagger a_- - a_-^\dagger a_+^\dagger a_- a_+] \\ &= \frac{i\hbar^2}{2} (a_+^\dagger a_+ - a_-^\dagger a_-) \end{aligned}$$

$$\boxed{[J_x, J_y] = i\hbar J_z}$$

The generalized expression is

$$[J_i, J_j] = i\hbar \epsilon_{ijk} J_k,$$

where

$$\epsilon_{ijk} = \begin{cases} 1, & \text{even permutations} \\ -1, & \text{odd permutations} \\ 0, & \text{otherwise} \end{cases}$$