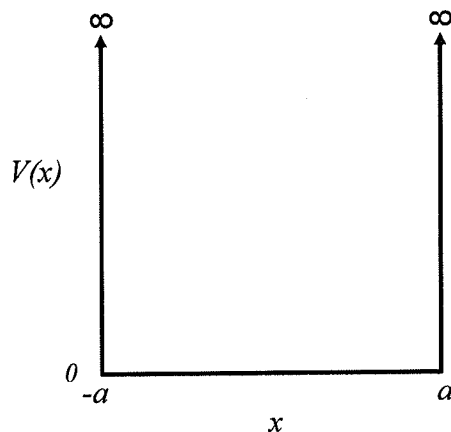


F-2007

PROBLEM 2



Consider the one-dimensional infinite-well potential shown above.

- [a] (4pts) Derive expressions for the energy eigenfunctions and energy eigenvalues for a particle in the one-dimensional infinite well shown above. Show your work.
- [b] (4pts) Now suppose a perturbation of the form

$$\Delta V(x) = V_o a \delta(x)$$

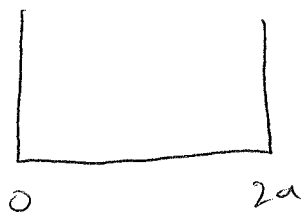
is added with $V_o \ll \frac{\hbar^2 \pi^2}{ma^2}$. Here $\delta(x)$ is the Dirac-delta function. According to first order perturbation theory, what are the eigenenergies of each state?

- [c] (2pts) According to first order perturbation theory, what is the wave function of the ground state? Write your answer in terms of a , V_o , fundamental constants, and the unperturbed wave functions $\phi_n(x)$. You do not have to normalize the wave function.

F-2007

①

② A Shift well



$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi$$

Assume $E > 0$

$$\rightarrow \frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi \rightarrow \frac{d^2\psi}{dx^2} = -k^2\psi$$

\uparrow
 k^2

$$\rightarrow \psi = A \sin kx + B \cos kx$$

ψ must be 0 at $x=0$

$\Rightarrow B \rightarrow 0$

$$\rightarrow \psi = A \sin kx \rightarrow \psi \text{ must be 0 at } 2a$$

$$k(2a) = n\pi$$

$$\rightarrow k = \frac{n\pi}{2a}$$

$$\rightarrow \psi = A \sin \frac{n\pi}{2a} x$$

Normalize $\rightarrow 1 = \int_0^{2a} A^2 \sin^2 \frac{n\pi}{2a} x dx = A^2 \left[\frac{x}{2} - \frac{\sin 2kx}{4k} \right]_0^{2a}$

$$= A^2 \left[\frac{x}{2} - \frac{\sin \frac{n\pi}{a} x}{\frac{4n\pi}{2a}} \right]_0^{2a}$$

$$= A^2 a = 1$$

$$\rightarrow A = \frac{1}{\sqrt{a}}$$

(A) Continued

(2)

$$\psi = \sqrt{\frac{1}{a}} \sin \frac{n\pi x}{2a}$$

$$\rightarrow \frac{\hbar^2 k^2}{4a^2} = \frac{2mE}{\hbar^2} \rightarrow \boxed{E_n = \frac{\hbar^2 n^2 \pi^2}{8ma^2}}$$

Shift well back $0 \rightarrow -a$

$$\psi = \sqrt{\frac{1}{a}} \sin \left[\frac{n\pi}{2a} (x-a) \right]$$

$$= \sin \frac{n\pi x}{2a} \cos \frac{n\pi a}{2a} + \cos \frac{n\pi x}{2a} \sin \frac{n\pi a}{2a}$$

$$(F) \quad n = 1, 3, 5 \rightarrow \frac{1}{\sqrt{a}} \cos \frac{n\pi x}{2a}$$

$$n = 2, 4, 6 \rightarrow \frac{1}{\sqrt{a}} \sin \frac{n\pi x}{2a}$$

$$(B) \quad \langle n | V | n \rangle = E_0^{(1)} = \int_{-a}^a \frac{1}{a} \cos^2 \frac{n\pi x}{2a} V_0 a \delta(x) dx$$

$$= \frac{V_0 a}{a} = \boxed{V_0 = E_0^{(1)} \quad (F \quad n = \text{odd})}$$

$$= \int_{-a}^a \frac{1}{a} \sin^2 \frac{n\pi x}{2a} V_0 a \delta(x) dx$$

$$= \boxed{0 \quad (F \quad n = \text{even})}$$

(3)

$$\textcircled{C} \quad \psi_0^{(1)} = \sum_{k \neq n} \frac{\langle k | V | n \rangle}{E_n - E_k} |k\rangle$$

$$\langle k | V | \phi \rangle = \frac{1}{a} \int_{-a}^a \cos \frac{kx}{2a} \sin \frac{n\pi x}{2a} dx \quad \text{or } a V_0 f(x) dx$$

$$= 0 \quad \text{IF } n \text{ IS EVEN}$$

$$= \frac{1}{a} \int_{-a}^a \cos \frac{kx}{2a} \cos \frac{n\pi x}{2a} dx \quad \text{or } V_0 dx$$

$$= V_0$$

$$\psi = \phi_1 + \sum_{\substack{k \neq 0 \\ \text{odd}}} \frac{(V_0 \cancel{m}^2)}{E_n^2 - E_k^2 (1-n^2)} \phi_k$$

$$E_0 - E_k = \frac{\hbar^2 k^2}{8m a^2} - \frac{n^2 \hbar^2 k^2}{8m a^2}$$

$$= \frac{\hbar^2 k^2}{8m a^2} (1-n^2)$$

PROBLEM 1: The Delta-Function Potential

Let us consider a single particle of mass m moving in one dimension with the Hamiltonian

$$H = T + V(x),$$

where the kinetic energy is

$$T = \frac{P^2}{2m} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2},$$

the potential energy is

$$V(x) = -V_0 \delta(x),$$

and $\delta(x)$ is the Dirac delta function.

- (a) [2 points] Find an expression for the discontinuity of the derivative of the wave function at $x = 0$.
- (b) [3 points] Find the ground state wave function.
- (c) [2 points] Find the ground state energy.
- (d) [3 points] Find the expectation value for the kinetic energy, $\langle T \rangle$.

5-2010

①

① $H = T + V(x)$

$V = -V_0 f(x)$

① $-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi$

(all $E < 0$ (bound state))

$\frac{d^2 \psi}{dx^2} = \frac{2mE}{\hbar^2} \psi$

\uparrow
 κ^2

$\rightarrow \psi = A e^{\kappa x} + B e^{-\kappa x}$

\uparrow \uparrow
blows up at ∞ blows up at $-\infty$

$\psi = \begin{cases} A e^{\kappa x} & x < 0 \\ B e^{-\kappa x} & x > 0 \end{cases}$

$\rightarrow \text{at } x=0 \rightarrow \boxed{A=B}$

\equiv

lim $\frac{-\hbar^2}{2m} \int_{-e}^e \frac{d^2 \psi}{dx^2} dx - V_0 \int_{-e}^e f(x) \psi(x) dx = \int_{-e}^e E \psi(x) dx$

$\epsilon \rightarrow 0$

$= \frac{-\hbar^2}{2m} \Delta \frac{d\psi}{dx} = V_0 \psi(0)$

$\rightarrow \boxed{\left. \frac{d\psi}{dx} \right|_{x=0} = -\frac{2mV_0}{\hbar^2} \psi(0)}$

(B)

$$\frac{d\psi}{dx} \Big|_{x>0} = -k_B$$

$$\frac{d\psi}{dx} \Big|_{x<0} = k_A = k_B$$

$$\Rightarrow \Delta \frac{d\psi}{dx} = \hbar^2 k_B = \frac{\hbar^2 m v_0}{t^2} \psi$$

$$\rightarrow k = \frac{m v_0}{t^2}$$

$$\Rightarrow \psi = B e^{-\frac{m v_0}{t^2} |x|}$$

even function because of

$$\rightarrow \text{Find } B \rightarrow 1 = B^2 \int_{-\infty}^{\infty} e^{-\frac{2m v_0}{t^2} |x|} dx \quad \text{absolute value}$$

$$= 2B^2 \int_0^{\infty} e^{-\frac{2m v_0}{t^2} x} dx$$

$$= \frac{-2t^2}{2m v_0} (0 - 1) = \frac{t^2}{m v_0} B^2 = 1$$

$$\rightarrow B = \frac{\sqrt{m v_0}}{t}$$

$$\rightarrow \psi = \frac{\sqrt{m v_0}}{t} e^{-\frac{m v_0}{t^2} |x|}$$

(3)

③ ~~$\frac{2mE}{\hbar^2} = \frac{mV_0}{\hbar^2}$~~

↑ dimensions

$$\frac{2mE}{\hbar^2} = \frac{m^2 V_0^2}{\hbar^4} \rightarrow \boxed{E = -\frac{mV_0^2}{2\hbar^2}}$$

$$\textcircled{D} \langle T \rangle = \int_{-\infty}^{\infty} \frac{-\hbar^2}{2m} \frac{mV_0}{\hbar^2} e^{-\frac{mV_0}{\hbar^2}|x|} \frac{d^2}{dx^2} e^{-\frac{mV_0}{\hbar^2}|x|} dx$$

or $\langle T \rangle = \langle E \rangle - \langle V \rangle$

$$\begin{aligned} \Rightarrow \langle V \rangle &= \int_{-\infty}^{\infty} -V_0 \delta(x) \frac{mV_0}{\hbar^2} e^{-\frac{mV_0}{\hbar^2}|x|} dx \\ &= -\frac{V_0^2 m}{\hbar^2} \end{aligned}$$

$\langle E \rangle$ is a constant

$$\begin{aligned} \rightarrow \langle T \rangle &= -\frac{mV_0^2}{2\hbar^2} + \frac{2V_0^2 \hbar}{2\hbar^2} \\ &= \frac{mV_0^2}{2\hbar^2} \end{aligned}$$

(4)

① Try other way

$$\langle T \rangle = -\frac{V_0}{2} \int_{-\infty}^{\infty} \frac{m^2 V_0^2}{\hbar^4} e^{-\frac{2mV_0}{\hbar^2} |x|} dx$$

↑
even function

$$= \frac{-V_0^3 m^2}{\hbar^4} \int_0^{\infty} e^{-\frac{2mV_0}{\hbar^2} |x|} dx$$

$$= \frac{+V_0^3 m^2}{\hbar^4} \frac{\hbar^2}{2mV_0} = \boxed{\frac{V_0^2 m}{2\hbar^2}}$$

S-2013

Problem 1: Bound States and Scattering for a Delta-Function Well

Consider a delta-function for a 1-D system,

$$V(x) = -g \delta(x) \quad (1)$$

where $g > 0$. We will consider the states of a particle of mass m interacting with this potential for both $E < 0$ and $E > 0$.

This potential has a single bound state $E_b < 0$.

- (a) [1 pt] Explain why the bound state wavefunction for the particle will have the form $\Psi(x) = ce^{-|x|/\lambda}$. (You don't need to solve for anything to answer this question.)
- (b) [2 pts] Derive the boundary conditions for $\Psi(x)$ and $\partial_x \Psi(x)$ at $x = 0$.
- (c) [1 pt] Using the boundary conditions at $x = 0$, determine the value of λ .
- (d) [1 pts] What is the energy of the bound state, E_b ? What is the normalization constant c ?
- (e) [2 pts] What is the uncertainty in position, Δx for the particle in this bound state?
- (f) [2 pts] Next consider a scattering state for this particle with energy $E > 0$

$$\begin{aligned} \Psi(x) &= e^{ikx} + ae^{-ikx}, \quad x < 0 \\ &= be^{ikx}, \quad x > 0 \end{aligned} \quad (2)$$

For this state, $E = \frac{\hbar^2 k^2}{2m}$

Using the boundary conditions you found in part (b), determine a and b , and the transmission and reflection coefficients for this scattering state.

S-2013

①

① $V(x) = -g\delta(x)$

② $E < 0$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi \rightarrow \frac{d^2\psi}{dx^2} = \frac{-2mE}{\hbar^2} \psi$$

$E < 0$

$$\hookrightarrow \frac{d^2\psi}{dx^2} = \frac{2mE}{\hbar^2} \psi$$

\uparrow
 κ^2

$$\psi(x) = A e^{\kappa x} + B e^{-\kappa x}$$
$$= C e^{-\kappa|x|}$$

Both terms blow up
at either end

③ $\psi(x) = \begin{cases} A e^{\kappa x} & x < 0 \\ B e^{-\kappa x} & x > 0 \end{cases}$

\rightarrow

$A = B$	First boundary Condition
---------	-----------------------------

Discontinuity in derivative at $x=0$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} - g\delta(x)\psi = E\psi$$

(2)

(B) Continuity

$$\int_{-\epsilon}^{+\epsilon} -\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} dx - \int_{-\epsilon}^{+\epsilon} g \delta(x) \psi dx = \int_{-\epsilon}^{+\epsilon} E \psi dx$$

$$\underbrace{-\frac{\hbar^2}{2m} \Delta \left(\frac{d\psi}{dx} \right)}_{\text{Left side}}$$

$$\underbrace{g \psi(0)}_{\text{Middle term}}$$

$$\underbrace{E \psi(\epsilon + \epsilon)}_{\text{Right side}}$$

$$\lim_{\epsilon \rightarrow 0} = 0$$

$$\rightarrow \Delta \left(\frac{d\psi}{dx} \right) = -\frac{2m}{\hbar^2} g \psi(0) \quad \text{Second boundary condition}$$

$$\rightarrow \left. \frac{d\psi}{dx} \right|_{x>0} = -Bk e^{-kx} \quad A=B$$

$$\rightarrow \left. \frac{d\psi}{dx} \right|_{x<0} = A k e^{kx}$$

(C)

$$\rightarrow -2Bk e^{-kx} = -\frac{2m}{\hbar^2} g \psi(0)$$

$$\rightarrow k = \frac{1}{\lambda} \rightarrow -\frac{2B}{\lambda} e^{-\frac{x}{\lambda}} = -\frac{2m}{\hbar^2} g \psi(0)$$

$$\rightarrow -\frac{2B}{\lambda} = -\frac{2m}{\hbar^2} g B$$

$$\rightarrow \frac{1}{\lambda} = \frac{m g}{\hbar^2} \rightarrow \boxed{\lambda = \frac{\hbar^2}{m g}}$$

(3)

(D)

$$\cancel{E = \frac{mg}{t^2}} = \frac{2mE}{t^2}$$

$$\frac{1}{\lambda} = \sqrt{\frac{2mE}{\hbar^2}} \rightarrow \frac{1}{\lambda^2} = \frac{2mE}{\hbar^2}$$

$$\rightarrow E = \frac{\hbar^2}{2m\lambda^2}$$

$$= \frac{\hbar^2 m^2 g^2}{2m\hbar^4} = \boxed{-\frac{mg^2}{2\hbar^2} = E}$$

$$1 = \int_{-\infty}^{\infty} |\psi(x)|^2 dx = 2 \int_0^{\infty} |\psi(x)|^2 dx$$

$$= 2B^2 \int_0^{\infty} e^{-2kx} dx$$

$$= 2B^2 \left[\frac{e^{-2kx}}{-2kx} \right]_0^{\infty}$$

$$= \frac{2B^2}{2k} \left(- (0 - 1) \right)$$

$$= \frac{2B^2}{2k} = \frac{B^2 \lambda}{\hbar} = 1$$

$$\hookrightarrow B^2 = \frac{1}{\lambda} = \frac{mg}{\hbar^2}$$

$$\rightarrow \boxed{B = \frac{\sqrt{mg}}{\hbar}}$$

$$\textcircled{E} \Delta x = \langle \Delta x^2 \rangle = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}$$

$$\rightarrow \langle x^2 \rangle = 2 \int_0^{\infty} x^2 \frac{m\gamma}{\hbar^2} e^{-2x/\lambda} dx$$

$$= \frac{2m\gamma}{\hbar^2} \int_0^{\infty} x^2 e^{-2x/\lambda} dx$$

$$= \frac{2m\gamma}{\hbar^2} \left[\frac{\Gamma(2+1)}{\left(\frac{2}{\lambda}\right)^3} \right]$$

$$= \frac{4m\gamma}{\hbar^2} \lambda^3 = \frac{\hbar^6}{m\gamma^3} \frac{m\gamma}{2\hbar^2} = \frac{1}{2} \frac{\hbar^4}{m\gamma^2}$$

$$\rightarrow \langle x \rangle = \frac{m\gamma}{\hbar^2} \int_{-\infty}^{\infty} x e^{2ix/\lambda} dx$$

[Handwritten scribbles and crossed-out work follow this equation]

$$e^{2ix/\lambda}$$

[Handwritten scribbles and crossed-out work follow this equation]

is an even function

Since absolute value $f(-x) = f(x)$

x is odd

$$\rightarrow \langle x \rangle = 0$$

5

(E) Continued

→

$$\Delta x = \frac{1}{\sqrt{2}} \frac{t^2}{mg}$$

(F)

$$1 + a = b$$

$$\left. \frac{d\psi}{dx} \right|_{x>b} = Lk e^{Lkx} \rightarrow Lk b$$

$$\left. \frac{d\psi}{dx} \right|_{x<b} = Lk e^{Lkx} - Lk a e^{-Lkx} \rightarrow Lk - Lk a$$

$$\rightarrow \Delta = Lk b - Lk + Lk a = \frac{-2m}{t^2} g(1+a)$$

$$= Lk(1+a) - Lk + Lk a = \frac{-2m}{t^2} g(1+a)$$

$$= 2Lk a = \frac{-2m}{t^2} g(1+a)$$

~~~~~

$$\frac{-2mg}{t^2} - \frac{2m}{t^2} a$$

$$\rightarrow a \left( 2Lk + \frac{2mg}{t^2} \right) = \frac{-2mg}{t^2}$$

$$\rightarrow a = \frac{-mg}{t^2} \left( \frac{1}{Lk + \frac{mg}{t^2}} \right) = \frac{-mg}{Lk t^2 + mg}$$

$$R = \frac{m^2 g^2}{(Lk t^2 + mg)^2}$$

$$a^2 = R = \frac{m^2 g^2}{(-Lk t^2 + mg)(Lk t^2 + mg)} = \frac{m^2 g^2}{k^2 t^4 + m^2 g^2} = R$$

Ⓟ Continued

⑥

$$I = 1 - \frac{m^2 g}{(Lk t^2 + m g)^2} = \frac{(Lk t^2 + m g)^2 - m^2 g^2}{(Lk t^2 + m g)^2}$$

$$= \frac{-k^2 t^4 + m^2 g^2 + 2Lk t^2 m g - m^2 g^2}{(Lk t^2 + m g)^2} = \frac{2Lk t^2 m g - k^2 t^4}{(Lk t^2 + m g)^2} = T$$

$$b = 1 + a \rightarrow \frac{Lk t^2 + m g - m g}{Lk t^2 + m g}$$

$$\rightarrow b = \frac{Lk t^2}{Lk t^2 + m g}$$

Probably better way to get T

$$15 \quad b^2 = \left( \frac{-Lk t^2}{-Lk t^2 + m g} \right) \left( \frac{Lk t^2}{Lk t^2 + m g} \right)$$

$$= \frac{k^2 t^4}{k^2 t^2 - Lk t^2 m g + Lk t^2 m g + m^2 g^2}$$

$$= \frac{k^2 t^2}{k^2 t^2 + m^2 g^2} = T$$

$$R + T = 1 \quad \checkmark$$