

PROBLEM 1: Stationary States

For a quantum system with a time independent Hamiltonian (\mathbf{H}), the wave function ($\Psi(x, t)$) is a linear combination of stationary state solutions ($\Psi_n(x, t)$) to the Schrödinger equation:

$$\Psi_n(x, t) = u_n(x) \exp(-iE_n t/\hbar)$$

where $u_n(x)$ are eigenfunctions of the Hamiltonian

$$\mathbf{H}u_n(x) = E_n u_n(x)$$

and they form a complete orthonormal basis.

- (a) Evaluate the uncertainty in the energy for a system in a stationary state with the wave function $\Psi(x, t) = \Psi_n(x, t)$. [Show all work.] (2 Points)
- (b) Derive the time evolution operator $\mathbf{U}(t, t_0)$ in terms of the Hamiltonian (\mathbf{H}), and apply it to a stationary state $\Psi_n(x, t_0 = 0)$. Describe the change in the stationary state. (2 Points)

Now consider a particle that starts out in a normalized wave function

$$\Psi(x, 0) = c_1 u_1(x) + c_2 u_2(x)$$

where the $u_n(x)$ are real eigenfunctions of the Hamiltonian and c_n are real.

- (c) Determine an expression for the wave function $\Psi(x, t)$ at subsequent times. (2 Points)
- (d) Evaluate the probability density and describe its motion in time. (3 Points)
- (e) Determine the uncertainty in the energy ΔE with $\Delta t = \tau$ that is the period of oscillation in (d). (1 Points)

S-2012

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①

$$\textcircled{1} \textcircled{A} \psi_n(x,t) = U_n(x) \exp\left(-i \frac{E_n t}{\hbar}\right)$$

$$\Delta E = \sqrt{\langle E^2 \rangle - \langle E \rangle^2}$$

$$\rightarrow \langle E^2 \rangle = \int_{-\infty}^{\infty} \psi_n^* E^2 \psi \, dx$$

$$= \int_{-\infty}^{\infty} U_n(x)^2 \underset{\substack{\uparrow \\ \text{constant}}}{E^2} dx = E^2 \underbrace{\int_{-\infty}^{\infty} U_n^2 dx}_{\text{normalized} = 1}$$

$$\rightarrow \langle E^2 \rangle = E^2$$

$$\rightarrow \langle E \rangle = \int_{-\infty}^{\infty} \psi_n^* E \psi_n \, dx = E \underbrace{\int_{-\infty}^{\infty} \psi_n^* \psi_n \, dx}_{=1}$$

$$= E$$

$$\rightarrow \Delta E = \sqrt{E^2 - E^2} = \boxed{0 = \Delta E}$$

(2)

(B) $\psi_n(x, t_0 = 0) = \psi_n(x)$

~~$i\hbar \frac{\partial \psi_n}{\partial t} = H \psi_n$~~ \rightarrow ~~$\psi_n = e^{-iE_n t/\hbar}$~~

time evolution operator \rightarrow ~~$U(t)$~~

$$U(t) \psi_n(x) = \psi_n(x, t)$$

$$\rightarrow \text{or} \quad i\hbar \frac{\partial \psi}{\partial t} = H \psi$$

$$\rightarrow i\hbar \frac{\partial}{\partial t} [U(t) \psi(x)] = H U(t) \psi(x)$$

$\rightarrow \psi(x)$ does not depend on time

$$\rightarrow i\hbar \frac{\partial U(t)}{\partial t} = H U(t)$$

$$\rightarrow \frac{\partial U(t)}{\partial t} = -\frac{iH}{\hbar} U(t)$$

$$\rightarrow U(t) = e^{-iHt/\hbar}$$

$$\Rightarrow U(t) \psi_n(x) = e^{-iE_n t/\hbar} \psi_n(x)$$

\rightarrow Standing wave that oscillates with angular frequency of $\frac{E_n}{\hbar} (\omega t)$

(3)

$$(C) \quad \psi(x, 0) = C_1 u_1(x) + C_2 u_2(x)$$

$$\psi(x, t) = U(t) \psi(x, 0) = C_1 e^{-iE_1 t/\hbar} u_1(x) + C_2 e^{-iE_2 t/\hbar} u_2(x)$$

$$(D) \quad \text{Probability density} = \langle \psi(x, t) | \psi(x, t) \rangle$$

$$= \int C_1 e^{iE_1 t/\hbar} u_1(x) + C_2 e^{iE_2 t/\hbar} u_2(x) \left[\right]$$

$$\times \int C_1 e^{-iE_1 t/\hbar} u_1(x) + C_2 e^{-iE_2 t/\hbar} u_2(x) \left[\right]$$

$$= C_1^2 u_1^2(x) + C_1 C_2 e^{i\frac{t}{\hbar}(E_1 - E_2)} u_1 u_2 \\ + C_2^2 u_2^2(x) + C_2 C_1 e^{i\frac{t}{\hbar}(E_2 - E_1)} u_2 u_1$$

$$= \left[C_1^2 u_1^2(x) + C_2^2 u_2^2(x) + 2 C_1 C_2 u_1 u_2 \cos\left(\frac{(E_1 - E_2)t}{\hbar}\right) \right]$$

The probability of finding a particle in $x+dx$ of
space oscillates with time

$$\text{if } E_1 = E_2 \quad |C_1 u_1 + C_2 u_2|^2 = \text{Probability}$$

density,

The probability
density is

stationary
with time

(E) $\Delta E = \sqrt{\langle E^2 \rangle - \langle E \rangle^2}$

(4)

At

~~$\Delta E = \sqrt{\langle E^2 \rangle - \langle E \rangle^2}$~~

This is the proper way to calculate $\langle E \rangle$ math should write

$\langle \psi | H | \psi \rangle = \langle \psi | i\hbar \frac{\partial}{\partial t} | \psi \rangle \stackrel{\downarrow}{=} \langle \psi | E | \psi \rangle$

$\langle \psi | H | \psi \rangle = \langle \psi | -\hbar^2 \frac{\partial^2}{\partial x^2} | \psi \rangle = \langle E^2 \rangle$

$\rightarrow \langle E^2 \rangle = -\hbar^2 \int_{-\infty}^{\infty} \psi^* \frac{\partial^2 \psi}{\partial x^2} dx$

$\rightarrow \frac{\partial^2}{\partial x^2} \psi = \left(\left(-\frac{E_1}{\hbar^2} \right) \left(-\frac{E_1}{\hbar^2} \right) \right) e^{-E_1 x / \hbar} \psi_1(x)$

$+ \left(-\frac{E_2}{\hbar^2} \right) \left(-\frac{E_2}{\hbar^2} \right) e^{-E_2 x / \hbar} \psi_2(x)$

$= -\frac{E_1^2}{\hbar^2} \psi_1 - \frac{E_2^2}{\hbar^2} \psi_2$

And you'll get cross terms that are orthogonal $\int \psi_1 \psi_2 = 0$

$= E_1^2 + E_2^2$

$\langle E \rangle = i\hbar \int_{-\infty}^{\infty} \psi^* \frac{\partial \psi}{\partial t} dx$

$\frac{d}{dt} \psi = -\frac{E_1}{\hbar} \psi_1 + -\frac{E_2}{\hbar} \psi_2$

$E_1 + E_2$

$\Delta E = \sqrt{E_1^2 + E_2^2 - E_1 - E_2}$

$= E_1(E_1 - 1) + E_2(E_2 - 1)$

$-E_1 - E_2 = -2E_1E_2$

Math

$(E_2 - E_1)^2$

$= E_2^2 + E_1^2 - 2E_1E_2$

F-2014

PROBLEM 1: Stationary and Non-Stationary States

Consider a quantum system whose particles are in the following state:

$$\Psi(x, t) = \frac{1}{\sqrt{8}}\psi_1(x)e^{-iE_1t/\hbar} - i\sqrt{\frac{3}{8}}\psi_3(x)e^{-iE_3t/\hbar} + \frac{1}{\sqrt{2}}\psi_5(x)e^{-iE_5t/\hbar}, \quad (1)$$

where $\psi_n(x)$, $n = 1, 2, 3 \dots$ are stationary states of the Hamiltonian governing the system,

$$H\psi_n(x) = E_n\psi_n(x).$$

Answer the following questions:

- a) Do you expect $\langle x \rangle$, $\langle x^2 \rangle$ and $\langle E \rangle$ to be time dependent or time independent? Discuss briefly, but do not calculate. (2 Points)
- b) Is the uncertainty ΔE positive, negative or zero? Is ΔE time dependent or time independent? Again, discuss briefly but do not calculate. (2 Points)
- c) Is $\Psi(t)$ above a solution of the time dependent Schrodinger equation? Demonstrate. (2 Points)
- d) If the stationary states $\psi_1(x)$, $\psi_3(x)$ and $\psi_5(x)$ are eigenstates of the harmonic oscillator, will any of your answers to part a) change? Justify. (2 Points)
- e) Now assume the particles are in the state

$$\Psi(x, t) = \psi_3(x)e^{-iE_3t/\hbar}.$$

Answer parts a) and b) for this state. (2 Points)

$$\textcircled{1} \textcircled{A} \frac{d\langle A \rangle}{dt} = \langle \frac{1}{i\hbar} [H, A] \rangle$$

$|F\rangle$ in an eigenbasis of
The Hamiltonian.

$$[X, H] \neq 0$$

$$[X^2, H] \neq 0$$

$$[E, H] \text{ doesn't make sense but } [H, H] = 0$$

No time dependence for E

but time dependence for x & x^2

$$\textcircled{B} \Delta E > 0 \quad \text{Proportional to } E_2 - E_1$$

↑

has to be real

$$\textcircled{C} i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

$$\rightarrow i\hbar \frac{\partial \psi}{\partial t} = \frac{1}{\sqrt{2}} \psi_1 \left(-\frac{LE_1}{\hbar} \right) e^{-LE_1 t/\hbar} - L \sqrt{\frac{3}{8}} \psi_3 \left(-\frac{LE_3}{\hbar} \right) e^{-LE_3 t/\hbar}$$

$$+ \frac{1}{\sqrt{2}} \psi_5 e^{-LE_5 t/\hbar} \left(-\frac{LE_5}{\hbar} \right) + i\hbar$$

$$= E_1 \dots$$

ψ_2, ψ_4, ψ_6 is.

(2)

$$\textcircled{D} \langle x \rangle = \alpha \langle a^\dagger \rangle + \langle a \rangle$$

$$= \cancel{\psi_1} \cancel{\psi_2^0} + \cancel{\psi_2} \cancel{\psi_1^0} + \cancel{\psi_5} \psi_6 + \dots$$

Nothing with change

$\langle x^2 \rangle$ Just do the calculation (don't ~~use~~ using
about QFT)

~~QFT~~ $\textcircled{E} \langle x \rangle$ will be zero

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 \psi^2 dx$$

Time independent

$\langle E \rangle$ Time independent

$$\Delta E = 0$$

~~QFT~~

definite way