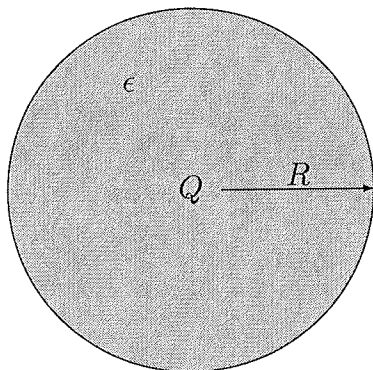


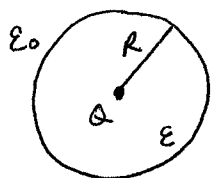
2. Consider a linear, homogeneous, isotropic, and non-dissipative dielectric (i.e., a dielectric where $\mathbf{D} = \epsilon \mathbf{E}$ and ϵ is a constant) in the shape of a sphere of radius R with a point charge Q embedded at its center.
- (a) {2 pts} Find the electric displacement vector \mathbf{D} , the electric field \mathbf{E} , and the polarization density \mathbf{P} inside the dielectric.
 - (b) {2 pts} Find the bound charge volume density ρ_D inside the dielectric.
 - (c) {1 pts} Find the total bound charge Q_D on the $r = R$ boundary of the dielectric.
 - (d) {2 pts} Find the net charge (free plus bound) at the center of the dielectric.
 - (e) {1 pts} Find the electric displacement vector \mathbf{D} , the electric field \mathbf{E} , and the polarization density \mathbf{P} , outside the dielectric sphere.
 - (f) {2 pts} Are \mathbf{D} and \mathbf{E} continuous at $r = R$? If not explain why.

(If you use Gaussian units you can put $\epsilon_0 = 1$.)

ϵ_0



Part (a)



We can use Gauss's law to determine the electric field inside the dielectric. So

$$\oint \vec{E} \cdot d\vec{a} = \frac{q_{enc}}{\epsilon}$$

$$E \oint da = \frac{Q}{\epsilon}$$

$$E(4\pi r^2) = \frac{Q}{\epsilon}$$

and so

$$\boxed{\vec{E} = \frac{1}{4\pi\epsilon} \frac{Q}{r^2} \hat{r}},$$

where r is the radius of our Gaussian sphere, and $r < R$. The displacement vector is

$$\begin{aligned} \vec{D} &= \epsilon \vec{E} \\ &= \epsilon \left(\frac{1}{4\pi\epsilon} \frac{Q}{r^2} \hat{r} \right) \end{aligned}$$

$$\boxed{\vec{D} = \frac{1}{4\pi} \frac{Q}{r^2} \hat{r}}.$$

We know

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P},$$

so

$$\begin{aligned} \vec{P} &= \vec{D} - \epsilon_0 \vec{E} \\ &= \frac{1}{4\pi} \frac{Q}{r^2} \hat{r} - \epsilon_0 \frac{1}{4\pi\epsilon} \frac{Q}{r^2} \hat{r} \end{aligned}$$

$$\boxed{\vec{P} = \frac{Q}{4\pi r^2} \left(1 - \frac{\epsilon_0}{\epsilon} \right) \hat{r}}$$

Part (b)

We want to find the bound charge volume density within the dielectric.
We know

$$\rho_b = - \vec{\nabla} \cdot \vec{P},$$

so

$$\begin{aligned} \rho_b &= - \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 P) \\ &= - \frac{1}{r^2} \frac{\partial}{\partial r} \left(\frac{Q}{4\pi} \left(1 - \frac{\epsilon_0}{\epsilon} \right) \right) \end{aligned}$$

and

$$\boxed{\rho_b = 0}.$$

Part (c)

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To find the bound charge on the surface, $r=R$, we first need to find the surface bound charge density, σ_b . We have

$$\begin{aligned}\sigma_b &= \vec{P} \cdot \hat{n} \big|_{\text{surface}} \\ &= \vec{P} \cdot \hat{r} \big|_{r=R} \\ &= \frac{Q}{4\pi R^2} \left(1 - \frac{\epsilon_0}{\epsilon}\right) \hat{r} \cdot \hat{r} \big|_{r=R} \\ &= \frac{Q}{4\pi R^2} \left(1 - \frac{\epsilon_0}{\epsilon}\right).\end{aligned}$$

Then the bound charge on the surface is

$$\begin{aligned}Q_b &= \int \sigma_b da \\ &= \int \frac{Q}{4\pi R^2} \left(1 - \frac{\epsilon_0}{\epsilon}\right) da \\ &= \int_0^{2\pi} \int_0^\pi \frac{Q}{4\pi R^2} \left(1 - \frac{\epsilon_0}{\epsilon}\right) R^2 \sin\theta d\theta d\phi\end{aligned}$$

and

$$\boxed{Q_b = Q \left(1 - \frac{\epsilon_0}{\epsilon}\right)}.$$

Part (d)

Now we want to find the net charge at the center of the sphere.
If we have

$$Q_b = Q \left(1 - \frac{\epsilon_0}{\epsilon}\right)$$

on the surface, then at the center we have

$$\begin{aligned} Q_b &= -Q_b \\ &= -Q \left(1 - \frac{\epsilon_0}{\epsilon}\right). \end{aligned}$$

Then since we have free charge

$$Q_f = Q$$

at the center, the total charge at the center is

$$\begin{aligned} Q_{tot} &= Q_b + Q_f \\ &= Q - Q \left(1 - \frac{\epsilon_0}{\epsilon}\right) \end{aligned}$$

$$\boxed{Q_{tot} = \frac{Q\epsilon_0}{\epsilon}}.$$

Part (c)

Outside of the dielectric,

$$\oint \vec{E} \cdot d\vec{a} = \frac{q_{enc}}{\epsilon_0}$$

$$E \oint da = \frac{q_{enc}}{\epsilon_0}$$

$$E(4\pi r^2) = \frac{q_{enc}}{\epsilon_0}$$

$$\boxed{\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}}$$

where $r \rightarrow R$. Then

$$\vec{D} = \epsilon_0 \vec{E}$$

$$\boxed{\vec{D} = \frac{Q}{4\pi r^2} \hat{r}}$$

Since we are outside of the material, we know

$$\boxed{\vec{P} = 0}$$

Part (f)

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At the surface, we have no free surface charge. Since the boundary condition on \vec{D} at the surface is

$$D_{\perp}^{\text{above}} - D_{\perp}^{\text{below}} = \sigma_{\text{free}},$$

then

$$D_{\perp}^{\text{above}} - D_{\perp}^{\text{below}} = 0$$

$$D_{\perp}^{\text{above}} = D_{\perp}^{\text{below}}$$

and the normal component of \vec{D} is continuous across the surface.

However,

$$E_{\perp}^{\text{above}} - E_{\perp}^{\text{below}} = \frac{\sigma_{\text{tot}}}{\epsilon_0}$$

$$E_{\perp}^{\text{above}} - E_{\perp}^{\text{below}} = \frac{\sigma_b}{\epsilon_0},$$

so the electric field's normal component is discontinuous by σ_b/ϵ_0 across the surface boundary.