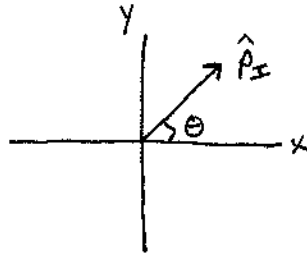


3. Light with wavenumber k and frequency ω is incident from vacuum normally along the z axis on a crystalline-plane parallel plate of thickness d . The crystal is cut so that a plane wave polarized along the x axis is refracted with index of refraction n_0 and one polarized along the y axis is refracted with index of refraction n_E .
- ~~(a)~~ [2 pts.] If the incident light is polarized at angle θ with respect to the x axis, write the form of the incident wave (using ω , θ , and space and time coordinates), assuming an incident amplitude of E_I .
- ~~(b)~~ [2 pts.] Now write the form of the refracted wave inside the crystal (using ω , θ , n_0 , n_E , and space and time coordinates) assuming that the refracted wave amplitude is E_T .
- ~~(c)~~ [6 pts.] Find the conditions of θ and d such that the refracted wave at d (the opposite side of the plate) is circularly polarized (either right or left).

Part (a)

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The incident light is polarized at an angle θ with respect to \hat{x} . So we have a setup like



where \hat{p}_\perp is the direction of polarization of our incident wave. We can write this as

$$\hat{p}_\perp = \cos \theta \hat{x} + \sin \theta \hat{y}.$$

Then

$$\vec{E}_i = E_\perp e^{i(kz - \omega t)} \hat{p}_\perp$$

$$\vec{E}_i = E_\perp e^{i(kz - \omega t)} (\cos \theta \hat{x} + \sin \theta \hat{y}).$$

Part (b)

We know that along the x -axis within the material, we have $k = k_o$. Along the y -axis, we have $k = k_E$. If $k = \frac{n\omega}{c}$, then

$$k_o = \frac{n_o \omega}{c}$$

$$k_E = \frac{n_E \omega}{c}.$$

The refracted wave is then given by

$$\vec{E}_t = E_T \left[e^{i(k_o z - \omega t)} \cos \theta \hat{x} + e^{i(k_E z - \omega t)} \sin \theta \hat{y} \right],$$

where E_T is the refracted wave amplitude.

Part (c)

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The wave is circularly polarized when the amplitudes in the x - and y -directions are equal, but offset by a phase difference of $\frac{\pi}{2}$.

Setting the amplitudes equal, we can put a restraint on θ .

$$E_{Tx} \cos \theta = E_{Ty} \sin \theta$$

$$\tan \theta = \frac{E_{Tx}}{E_{Ty}}.$$

To determine E_{Tx} and E_{Ty} , we want to apply the boundary conditions at the interface. Specifically, we want to apply the tangential boundary conditions

$$\vec{E}_{1,\parallel} - \vec{E}_{2,\parallel} = 0$$

$$\vec{H}_{1,\parallel} - \vec{H}_{2,\parallel} = 0,$$

which take into account the fact that there is no free current or free surface charge density. Then

$$\vec{E}_{1,\parallel} - \vec{E}_{2,\parallel} = 0$$

$$(\vec{E}_i + \vec{E}_r - \vec{E}_t) \times \hat{n} = 0,$$

which ends up giving us

$$E_i + E_r - E_t = 0.$$

Applying the second boundary condition...

$$\vec{H}_{1,\parallel} - \vec{H}_{2,\parallel} = 0$$

$$\frac{1}{\mu_0} \vec{B}_{1,\parallel} - \frac{1}{\mu_0} \vec{B}_{2,\parallel} = 0$$

$$\vec{B}_{1,\parallel} - \vec{B}_{2,\parallel} = 0$$

$$(\vec{B}_i + \vec{B}_r - \vec{B}_t) \times \hat{n} = 0$$

$$B_i - B_r - B_t = 0$$

$$B_i - B_r = B_t$$

Part (c), cont'd

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but

$$\vec{B} = n \hat{k} \times \vec{E},$$

so since $n = \sqrt{\epsilon\mu}$, and $\epsilon = \epsilon_0 = 1$ and $\mu = \mu_0 = 1$ in vacuum, we have

$$E_I - E_R = n E_T.$$

plugging in E_R from our first condition...

$$E_I - (E_T - E_I) = n E_T$$

$$2E_I = (n+1)E_T$$

$$E_T = \frac{2}{n+1} E_I.$$

Then

$$E_{Tx} = \frac{2}{n_0+1} E_I$$

$$E_{Ty} = \frac{2}{n_E+1} E_I$$

and

$$\begin{aligned} \tan \theta &= \frac{E_{Tx}}{E_{Ty}} \\ &= \frac{\frac{2}{n_0+1} E_I}{\frac{2}{n_E+1} E_I} \\ &= \frac{n_E+1}{n_0+1} \end{aligned}$$

$$\theta = \arctan \left(\frac{n_E+1}{n_0+1} \right).$$

Part (c), cont'd

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We want to determine the condition on d . At the end of the plate, $(z+d)$, we know that the phase difference between \vec{E}_i and \vec{E}_t is $\pi/2$. So...

$$[k_o(z+d) - \omega t] - [k_E(z+d) - \omega t] = \frac{\pi}{2}$$

$$(k_o - k_E)z + (k_o - k_E)d = \frac{\pi}{2}$$

$$(k_o - k_E)d = \frac{\pi}{2} - (k_o - k_E)z$$

$$d = \frac{\frac{\pi}{2} - (k_o - k_E)z}{k_o - k_E}$$

$$d = \frac{\pi - 2(k_o - k_E)z}{2(k_o - k_E)}$$