

2. (a) {2 pts} In a homogeneous, linear and isotropic conducting material whose electromagnetic properties (at low frequencies) are described by constant values of the permittivity, permeability, and conductivity respectively ϵ , μ , and σ , show that Maxwell's equations require that the electric field satisfy

$$\nabla^2 \mathbf{E} - \epsilon\mu \frac{\partial^2 \mathbf{E}}{\partial t^2} - \sigma\mu \frac{\partial \mathbf{E}}{\partial t} = 0, \quad (SI)$$

$$\nabla^2 \mathbf{E} - \frac{\epsilon\mu}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} - \frac{4\pi\sigma\mu}{c^2} \frac{\partial \mathbf{E}}{\partial t} = 0. \quad (Gaussian)$$

- (b) {2 pts} Given a plane polarized plane wave of angular frequency ω whose electric field is of the form

$$\mathbf{E}(z, t) = \text{Real} \{ \hat{i} E_0 e^{i(kz - \omega t)} \},$$

evaluate k^2 as a function of ϵ , μ , σ , and ω .

- (c) {2 pts} Find the real and imaginary parts of k assuming $\sigma \gg \omega\epsilon$.
 (d) {2 pts} Using your results from (c) find the skin depth δ of the conductor. The skin depth is defined by the depth at which the wave's amplitude decreases by e^{-1} , i.e.,

$$\frac{|\mathbf{E}(z + \delta, t)|}{|\mathbf{E}(z, t)|} = \frac{1}{e}$$

- (e) {2 pts} Using Maxwell's equations, find the magnetic field $\mathbf{H}(x, t)$ associated with $\mathbf{E}(z, t)$ given in (b) and discuss their phase difference when $\sigma \gg \omega\epsilon$.

Part (a) (Gaussian)

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From the law of induction,

$$\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}.$$

The Ampère-Maxwell law tells us

$$\vec{\nabla} \times \vec{H} = \frac{1}{c} (4\pi \vec{J}_f + \frac{\partial \vec{D}}{\partial t})$$

$$\vec{\nabla} \times \frac{1}{\mu} \vec{B} = \frac{4\pi}{c} \vec{J}_f + \frac{\epsilon}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times \frac{1}{\mu} \vec{B} = \frac{4\pi}{c} (\sigma \vec{E}) + \frac{\epsilon}{c} \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \frac{4\pi\mu\sigma}{c} \vec{E} + \frac{\epsilon\mu}{c} \frac{\partial \vec{E}}{\partial t}.$$

Taking the curl of Faraday's law,

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = -\frac{1}{c} \frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B})$$

$$\vec{\nabla} (\vec{\nabla} \cdot \vec{E}) - \vec{\nabla}^2 \vec{E} = -\frac{1}{c} \frac{\partial}{\partial t} \left(\frac{4\pi\mu\sigma}{c} \vec{E} + \frac{\epsilon\mu}{c} \frac{\partial \vec{E}}{\partial t} \right).$$

but

$$\vec{\nabla} \cdot \vec{E} = 0$$

since this is a conductor, so

$$-\vec{\nabla}^2 \vec{E} = -\frac{\epsilon\mu}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} - \frac{4\pi\sigma\mu}{c^2} \frac{\partial \vec{E}}{\partial t}$$

or

$$\boxed{\vec{\nabla}^2 \vec{E} - \frac{\epsilon\mu}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} - \frac{4\pi\sigma\mu}{c^2} \frac{\partial \vec{E}}{\partial t} = 0,}$$

as expected.

Part (b)

Plugging in

$$E = E_0 e^{i(kz - \omega t)}$$

into our expression from part (a)...

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) E - \frac{\epsilon \mu}{c^2} \frac{\partial^2}{\partial t^2} E - \frac{4\pi\sigma\mu}{c^2} \frac{\partial E}{\partial t} = 0$$

$$\frac{\partial^2}{\partial z^2} E - \frac{\epsilon \mu}{c^2} \frac{\partial^2}{\partial t^2} E - \frac{4\pi\sigma\mu}{c^2} \frac{\partial E}{\partial t} = 0$$

$$-k^2 E_0 e^{i(kz - \omega t)} + \frac{\epsilon \mu}{c^2} \omega^2 E_0 e^{i(kz - \omega t)} + \frac{4\pi\sigma\mu\omega}{c^2} i E_0 e^{i(kz - \omega t)} = 0$$

$$-k^2 + \frac{\epsilon \mu}{c^2} \omega^2 + i \frac{4\pi\sigma\mu\omega}{c^2} = 0$$

and

$$\boxed{k^2 = \frac{1}{c^2} (\epsilon \mu \omega^2 + i 4\pi \sigma \mu \omega)} .$$

Part (c)

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If $\sigma \gg \omega \epsilon$, then $\epsilon \ll \frac{\sigma}{\omega}$. So

$$\begin{aligned} k^2 &= \frac{1}{c^2} (\epsilon \mu \omega^2 + i 4\pi \sigma \mu \omega) \\ &= \frac{\omega^2}{c^2} \left(\epsilon \mu + i \frac{4\pi \sigma \mu}{\omega} \right). \end{aligned}$$

but $\epsilon \ll \sigma/\omega$, so

$$\begin{aligned} k^2 &\approx \frac{\omega^2}{c^2} \left(i \frac{4\pi \sigma \mu}{\omega} \right) \\ &= i \frac{4\pi \sigma \mu \omega}{c^2} \end{aligned}$$

and

$$\begin{aligned} k &= \sqrt{i \frac{4\pi \sigma \mu \omega}{c^2}} \\ &= \sqrt{i} \sqrt{\frac{4\pi \sigma \mu \omega}{c^2}}. \end{aligned}$$

Since $i = e^{\frac{i\pi}{2}}$, then $\sqrt{i} = e^{\frac{i\pi}{4}}$. We have

$$\begin{aligned} e^{\frac{i\pi}{4}} &= i \sin\left(\frac{\pi}{4}\right) + \cos\left(\frac{\pi}{4}\right) \\ &= i \frac{\sqrt{2}}{2} + \frac{\sqrt{2}}{2}, \end{aligned}$$

so

$$\begin{aligned} k &= \left(i \frac{\sqrt{2}}{2} + \frac{\sqrt{2}}{2} \right) \sqrt{\frac{4\pi \sigma \mu \omega}{c^2}} \\ &= \sqrt{\frac{2\pi \sigma \mu \omega}{c^2}} (i + 1). \end{aligned}$$

So

$$\begin{aligned} \operatorname{Re}(k) &= \sqrt{\frac{2\pi \sigma \mu \omega}{c^2}} \\ \operatorname{Im}(k) &= i \sqrt{\frac{2\pi \sigma \mu \omega}{c^2}} \end{aligned}$$

Part (d)

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We can write the electric field as

$$\begin{aligned}\vec{E} &= E_0 e^{i(\operatorname{Im}(k)z + \operatorname{Re}(k)z - \omega t)} \hat{z} \\ &= E_0 e^{i\left(\sqrt{\frac{2\pi\sigma\mu\omega}{c^2}} - \omega t\right)z - \sqrt{\frac{2\pi\sigma\mu\omega}{c^2}}z} \hat{z}\end{aligned}$$

In order for the amplitude to decrease by e^{-1} , we must have

$$z = \sqrt{\frac{c^2}{2\pi\sigma\mu\omega}},$$

which implies the skin depth δ must be

$$\boxed{\delta = \sqrt{\frac{c^2}{2\pi\sigma\mu\omega}}}.$$

Part (e)

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From the law of induction,

$$\vec{\nabla} \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}.$$

So ...

$$\vec{\nabla} \times \vec{E} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E & 0 & 0 \end{vmatrix} = \hat{i}(0) - \hat{j}\left(-\frac{\partial E}{\partial z}\right) + \hat{k}\left(-\frac{\partial E}{\partial y}\right)$$

$$\frac{\partial \vec{E}}{\partial z} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$ik E_0 e^{i(kz - \omega t)} \hat{j} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t}$$

$$\frac{\partial \vec{B}}{\partial t} = -ick E_0 e^{i(kz - \omega t)} \hat{j}$$

$$\vec{B} = -ick E_0 \int e^{i(kz - \omega t)} dt \hat{j}$$

$$= \frac{ck}{\omega} E_0 e^{i(kz - \omega t)} \hat{j}$$

but $\vec{H} = \frac{1}{\mu} \vec{B}$, so

$$\boxed{\vec{H} = \frac{kc}{\mu\omega} E_0 e^{i(kz - \omega t)} \hat{j}}$$