

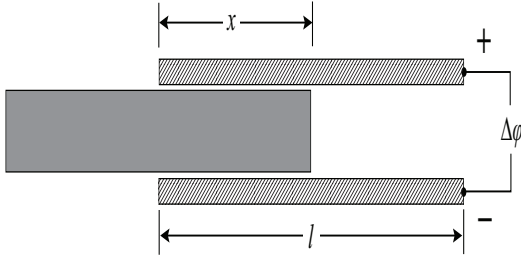
## **E & M Qualifier**

January 13, 2011

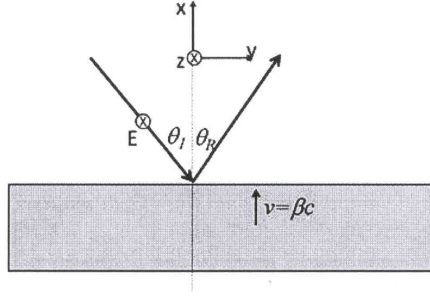
**To insure that the your work is graded correctly you MUST:**

1. use only the blank answer paper provided,
2. write only on one side of the page,
3. put your alias on every page,
4. put the problem # on every page,
5. start each problem by stating your units e.g., SI or Gaussian,
6. number every page starting with 1 for each problem,
7. put the total # of pages you use for that problem on every page,
8. staple your exam when done.

**Use only the reference material supplied (Schaum's Guides).**



1. A parallel plate capacitor has the region between its plates filled with a dielectric slab of dielectric constant  $K = \epsilon/\epsilon_0$  and mass  $m$ . The plate dimensions are: width  $w$ , length  $\ell$ , and plate separation  $d$ . The capacitor plates are connected to a battery of constant voltage  $V$  ( $\Delta\phi = V$  in the figure). Neglect the fringe field and friction, and assume the slab is constrained to move in the plane parallel to the capacitor plates.
  - (a) {2 pts} Compute the capacitance  $C \equiv q/V$  of this capacitor as a function of  $x$ .
  - (b) {2 pts} If the slab is withdrawn half way (to  $x = \ell/2$ ) and held in place, what is the magnitude and direction of the force on the slab caused by the electric field?
  - (c) {2 pts} At  $x = \ell/2$  the slab is released and given a velocity  $v_0$  to the right. Find the current supplied by the battery at the instant it is released.
  - (d) {2 pts} At  $x = \ell/2$  the slab is again released but with zero velocity. Describe the motion of the slab (in words). What is the maximum velocity achieved by the slab?
  - (e) {2 pts} Sketch the displacement of the slab versus time.



2. This problem investigates the shifting frequency of electromagnetic radiation that is reflected off a moving target. Incident and reflected frequencies and angles are not the same if the target is moving.

Assume that in the lab frame of reference, the target is a flat mirror traveling upward in the positive  $x$ -direction parallel to the mirror's normal with velocity  $\mathbf{v} = \beta c \hat{\mathbf{x}}$  (see the figure). Also assume the wave is a linearly polarized plane wave traveling in vacuum towards the moving mirror at angle  $\theta_I$  (relative to the mirror's normal). If the polarization is in the  $\hat{\mathbf{z}}$  direction, the incident electric field is given by

$$\mathbf{E}_I = E_0 \hat{\mathbf{z}} e^{i(\mathbf{k}_I \cdot \mathbf{r} - \omega_I t)},$$

with

$$\mathbf{k}_I = \frac{\omega_I}{c} (-\cos \theta_I \hat{\mathbf{x}} + \sin \theta_I \hat{\mathbf{y}}).$$

- (a) {2 pts} Write the Lorentz boost  $A$  as a function of  $\beta$  and  $\gamma \equiv \sqrt{1 - \beta^2}$  that transforms the Lab coordinates  $\mathbf{r}$  and  $ct$  to coordinates  $\mathbf{r}'$  and  $ct'$  co-moving with the mirror. Also give the inverse  $A^{-1}$  of the Lorentz boost  $A$  that transforms the moving coordinates  $\mathbf{r}'$  and  $ct'$  into Lab coordinates  $\mathbf{r}$  and  $ct$ .
- (b) {3 pts} By rewriting the above wave's phase in both reference frames, i.e.,

$$\mathbf{k}_I \cdot \mathbf{r} - \omega_I t = \mathbf{k}'_I \cdot \mathbf{r}' - \omega'_I t'$$

as a function of the co-moving mirror coordinates  $\mathbf{r}'$  and  $ct'$  (i.e., use  $A^{-1}$ ) find  $\mathbf{k}'_I$  and  $\omega'_I$  as observed in the co-moving frame. These will be functions of  $\beta, \gamma$ , and  $\theta_I$  as well as  $\omega_I$ .

- (c) {2 pts} By writing the incident wave vector just obtained in the moving frame in the form

$$\mathbf{k}'_I = \frac{\omega'_I}{c} (-\cos \theta'_I \hat{\mathbf{x}} + \sin \theta'_I \hat{\mathbf{y}}),$$

determine the incident angle  $\theta'_I$  as seen by observers moving with the mirror (e.g., give  $\cos \theta'_I$  as a function of  $\theta_I, \omega_I$  and the Lorentz parameters  $\beta, \gamma$ ).

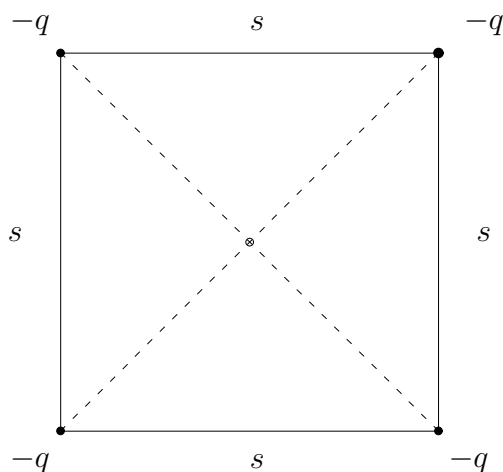
- (d) {3 pts} If, as seen by observers moving with the mirror, the reflected wave has the same frequency as the incident wave  $\omega'_R = \omega'_I$  and a reflection angle that is the same as the incidence angle  $\theta'_R = \theta'_I$ , i.e.,

$$\mathbf{k}'_R = \frac{\omega'_I}{c}(\cos \theta'_I \hat{\mathbf{x}} + \sin \theta'_I \hat{\mathbf{y}}),$$

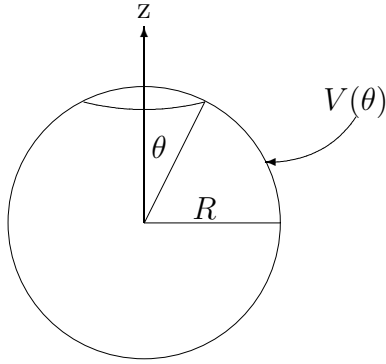
what is the frequency  $\omega_R$  of the reflected light as measured in the laboratory frame? Hint: again use

$$\mathbf{k}_R \cdot \mathbf{r} - \omega_R t = \mathbf{k}'_R \cdot \mathbf{r}' - \omega'_R t',$$

and the Lorentz boosts  $A$ .



3. Consider a square with sides of length  $s$  and charges  $-q$  at the corners as shown:
- (a) {2 pts} What is the potential at the center of the square if the potential is zero at  $\infty$ ?
  - (b) {2 pts} How much work does it take to bring in another charge  $-q$  from  $\infty$  to the center of the square?
  - (c) {3 pts} How much work does it take to assemble the original configuration of 4 negative charges (no charge at center)?
  - (d) {3 pts} Now suppose that instead of the 4 charges being located at the corners of a square, a net charge of  $-4q$  is distributed uniformly on the surface of a sphere of radius  $s$ . How much work does it take to bring in another charge  $q$  from  $\infty$  to the center of the sphere?



4. Consider an isolated spherical surface of radius  $R$  centered on the origin, that is kept at a known potential  $V(\theta)$ , i.e.,

$$\Phi(r = R, \theta) = V(\theta)$$

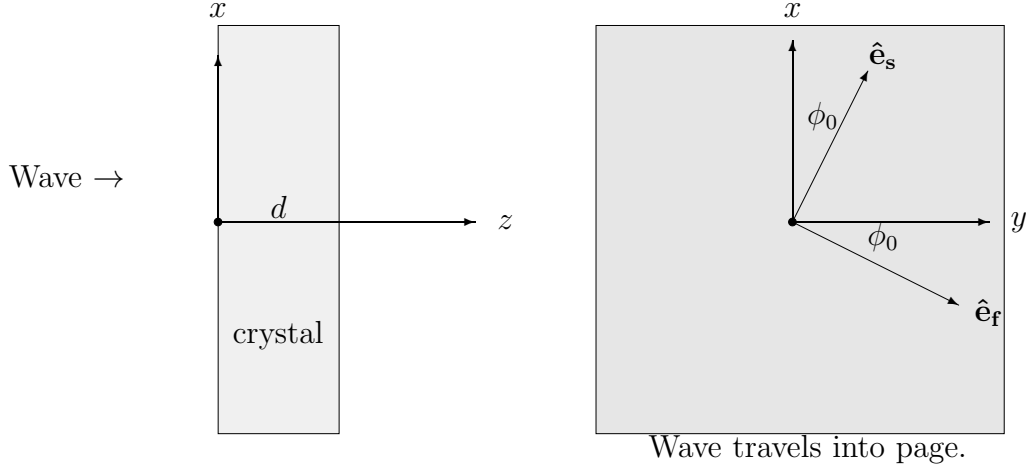
where  $(r, \theta, \phi)$  are the usual spherical polar coordinates, i.e.,  $\theta$  is measured with respect to a  $z$ -axis passing through the center of the sphere and  $\phi$  is the azimuthal angle about the  $z$ -axis measured from the  $x$  axis.

- (a) {2 pts} Write down expressions for the general solution to  $\nabla^2 \Phi(r, \theta) = 0$  for the electrostatic potential as a linear combination of Legendre polynomials in the respective regions  $0 \leq r < R$  and  $r > R$ . Assume that the potential vanishes at  $r \rightarrow \infty$  and has azimuthal symmetry i.e., no dependence on the angle  $\phi$ . Do not include terms that must vanish. Do not attempt to evaluate the constants that appear in the linear combination but do give the correct  $r$  dependence of each term.
- (b) {2 pts} What boundary conditions must your two expressions satisfy at the junction  $r = R$  to have a unique solution to Maxwell's equations?
- (c) {2 pts} If the particular surface potential imposed is

$$\Phi(r = R, \theta) = V_0 \cos \theta$$

where  $V_0$  is a constant, what is the explicit form of your potential for both regions  $r \leq R$  and  $r > R$ ?

- (d) {2 pts} Determine the resulting electric field on both sides of the  $r=R$  surface.
- (e) {2 pts} What is the surface charge density  $\sigma(\theta)$  on the spherical shell at  $r=R$ .



5. A plane polarized monochromatic light wave traveling in the  $+z$  direction enters a large flat slab of transparent crystal of thickness  $d$ , located between  $z = 0$  and  $z = d$ . This crystal has the property that the index of refraction depends on the direction of polarization as follows: Plane waves traveling in the  $z$  direction but polarized in the direction

$$\hat{\mathbf{e}}_s = \cos \phi_0 \hat{\mathbf{x}} + \sin \phi_0 \hat{\mathbf{y}},$$

travel with speed  $v_s = c/n_s < c$  but those polarized in the orthogonal direction

$$\hat{\mathbf{e}}_f = -\sin \phi_0 \hat{\mathbf{x}} + \cos \phi_0 \hat{\mathbf{y}},$$

travel with the faster speed  $v_f = c/n_f < c$  where  $n_s = n_f + \delta n$ .

Assume the wave, just after entering the crystal (i.e., for very small  $z \ll \lambda < d$ ), is polarized in the  $y$  direction and hence has the form

$$\mathbf{E}(z \approx 0, t) = E_0 \hat{\mathbf{y}} e^{-i\omega t}.$$

- (a) {4 pts} Prove that in general the initial plane wave becomes elliptically polarized when it reaches  $z = d$  by deriving the following expression

$$\mathbf{E}(z = d, t) = [E_x \hat{\mathbf{x}} + E_y \hat{\mathbf{y}}] e^{i(\bar{k}d - \omega t)},$$

where

$$\bar{k} \equiv \frac{\omega}{c} \left( \frac{n_s + n_f}{2} \right),$$

and

$$E_x = iE_0 \sin 2\phi_0 \sin \delta,$$

$$E_y = E_0 (\cos \delta - i \cos 2\phi_0 \sin \delta),$$

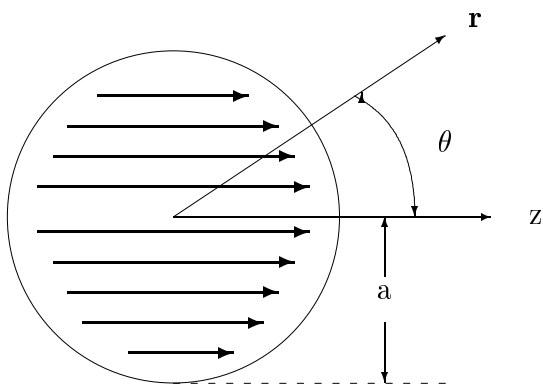
with

$$\delta \equiv \frac{\omega d}{2c} \delta n.$$

Hint: Write the wave at  $z=0$  as a combination of slow and fast plane polarized parts using  $\hat{\mathbf{y}} = \sin \phi_0 \hat{\mathbf{e}}_s + \cos \phi_0 \hat{\mathbf{e}}_f$ .

- (b) {3 pts} For what values of  $\delta$  and  $\theta_0$  will the wave emerge from the crystal as a circularly polarized wave? ( $E_x/E_y = \pm i$ ).
- (c) {3 pts} For what minimum crystal thicknesses  $d = d_{min}$  will the wave emerge as a plane polarized wave ( $E_x/E_y = \text{real}$ ) and what will its polarization direction be?





6. A permanent magnet in the shape of a solid sphere of radius  $a$  is oriented on the  $z$ -axis as shown in the figure. The magnetization of the magnet is given by  $\vec{M} = M_0 \hat{z}$ . [Recall that  $\nabla \times \mathbf{H} = 0$  implies the existence of a magnetic scalar potential  $\Phi_m(r, \theta)$  related to the magnetic field by  $\mathbf{H} = -\vec{\nabla} \Phi_m(r, \theta)$ .]
- (a) {4 pts} Compute the scalar magnetic potential  $\Phi_m(r, \theta)$  at all points  $r < a$  and  $r > a$ .
  - (b) {3 pts} Compute the magnetic Field  $\mathbf{H} = -\vec{\nabla} \Phi_m(r, \theta)$  at all points  $r < a$  and  $r > a$ .
  - (c) {3 pts} Compute the magnetic induction  $\mathbf{B}$ , where

$$\begin{aligned} \mathbf{B}/\mu_0 &= \mathbf{H} + \mathbf{M}, & (SI) \\ \mathbf{B} &= \mathbf{H} + 4\pi\mathbf{M}, & (Gaussian) \end{aligned}$$

at all points  $r < a$  and  $r > a$ .

Hints: The magnetic potential is axial symmetric about the  $z$ -axis and satisfies the Laplace equation at all points except  $r = a$ . Legendre polynomials are useful.