

5. Consider an infinitely long, solid, nonmagnetic conducting rod of radius  $a$  centered on the  $z$  axis. An infinitely long, hollow, conducting cylinder with inner radius  $b > a$  and outer radius  $d$  is coaxial with the rod. Let  $r$  be the radial distance perpendicular to the axis of the rod and the cylinder. The region between the conducting rod and the conducting cylinder (that is,  $a < r < b$ ) is filled with a nonconducting, linear, isotropic magnetic material with a constant relative permeability  $K = \mu/\mu_0$ , where  $\mu$  is the permeability of the material, and  $\mu_0$  is the permeability of free space ( $\mu_0 = 1$  in Gaussian units).

The rod carries a current  $I$  in the  $+z$  direction while the concentric cylinder carries a current  $I$  in the  $-z$  direction. We assume that the current density  $\mathbf{j}$  is uniform and of the same magnitude in both the rod and the cylinder,

$$j = \frac{I}{\pi a^2} = \frac{I}{\pi(d^2 - b^2)}.$$

- a) 3 pts. Calculate the magnetic field  $H(r)$  for the four regions

$$\text{I: } r \leq a, \quad \text{II: } a \leq r \leq b, \quad \text{III: } b \leq r \leq d, \quad \text{IV: } d \leq r.$$

- b) 3 pts. Calculate the magnetic flux (per unit length in the  $z$  direction) crossing a half-plane extending from the axis of the coaxial system and extending to infinity, that is, the surface defined by  $x > 0$ ,  $y = 0$ ,  $-\infty < z < \infty$ . Use this result to find the self-inductance  $L$  per unit length of the coaxial conductor.
- c) 2 pts. Compute the magnetic energy  $U$  per unit length along the  $z$  axis stored in the region filled with the linear magnetic material, that is for region II,  $a < r < b$ .
- d) 2 pts. Using the result from part c), show that the contribution to  $L$  coming from the region  $a \leq r \leq b$ ,  $L_{\text{II}}$ , is consistent with the contribution from the same region that you calculated in part b) above. That is, compute  $\frac{1}{2}L_{\text{II}}I^2$  and compare with the result of part c).

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E+M #5

Gaussian

a) We can compute the magnetic field via Ampere's Law (from  $\nabla \times \mathbf{B} = \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} = \frac{4\pi}{c} \mathbf{j}$ )

$$\int \nabla \times \mathbf{B} \cdot d\mathbf{a} = \int \frac{4\pi}{c} \mathbf{j} \cdot d\mathbf{a}$$

$$\mathbf{H} = \mathbf{B} - 4\pi \mathbf{M}$$

$$\oint \mathbf{B} \cdot d\mathbf{l} = \frac{4\pi}{c} \int \mathbf{j} \cdot d\mathbf{a}$$

$$\mathbf{B} = \mu \mathbf{H}, \quad \mu = 1 + 4\pi \chi_m$$

$$\mathbf{M} = \chi_m \mathbf{H} \quad \text{for linear, isotropic material}$$

- We have 4 regions to consider:

Region I:  $r \leq a$

$$\oint \mathbf{B} \cdot d\mathbf{l} = \frac{4\pi}{c} \int \left( \frac{I}{\pi a^2} \right) \cdot d\mathbf{a}$$

$$\mathbf{B} \cdot 2\pi r = \frac{4\pi}{c} \frac{I}{\pi a^2} \cdot \pi r^2$$

$$\mathbf{B} = \frac{2I}{ca^2} \hat{\phi} \quad (\text{direction determined via symmetry})$$

$$\mathbf{H} = \mathbf{B} - 4\pi \mathbf{M}$$

$$\mathbf{H} = \frac{2I}{ca^2} \hat{\phi} - 0 \quad (\text{non-magnetic material})$$

$$\mathbf{H} = \frac{2I}{ca^2} \hat{\phi}$$

Region II:  $a \leq r \leq b$

$$\oint \mathbf{B} \cdot d\mathbf{l} = \frac{4\pi}{c} \int \frac{I}{\pi a^2} \cdot d\mathbf{a}$$

$$\mathbf{B} \cdot 2\pi r = \frac{4\pi}{c} \frac{I}{\pi a^2} \cdot \pi a^2$$

$$\mathbf{B} = \frac{2I}{cr} \hat{\phi}$$

\* Because our magnetization is non-zero,  $\mathbf{M} = \chi_m \mathbf{H}$ ,  $\mathbf{H} = \frac{1}{\mu} \mathbf{B}$ ,  $\mu = 1 + 4\pi \chi_m$

$$\mathbf{H} = \mathbf{B} - 4\pi \mathbf{M}$$

$$\Rightarrow \mathbf{H} = \frac{2I}{\mu r} \hat{\phi}$$

$$\mathbf{H} = \mathbf{B} - 4\pi \chi_m \mathbf{H}$$

$$\mathbf{H}(1 + 4\pi \chi_m) = \mathbf{B}$$

$$\mathbf{H} \mu = \mathbf{B}$$

$$\mathbf{H} = \frac{1}{\mu} \mathbf{B}$$

## #5 (cont.)

Region III:  $b \leq r \leq d$

\* Must separate  $\vec{j}$ -integral to account for multiple currents

$$\oint \vec{B} \cdot d\vec{\ell} = \frac{4\pi}{c} \int \vec{j} \cdot d\vec{a}$$

$$B \cdot 2\pi r = \frac{4\pi}{c} \left[ \int \vec{j}_{rod} \cdot d\vec{a}_{rod} + \int \vec{j}_{shell} \cdot d\vec{a}_{shell} \right]$$

$$B \cdot 2\pi r = \frac{4\pi}{c} \left[ \frac{I}{\pi a^2} \cdot \pi a^2 + \frac{-I}{\pi(d^2 - b^2)} \cdot \pi(r^2 - b^2) \right]$$

$$B \cdot 2\pi r = \frac{4\pi}{c} \left[ I - \frac{I(r^2 - b^2)}{d^2 - b^2} \right]$$

$$\vec{B} = \frac{2I}{cr} \left[ 1 - \frac{r^2 - b^2}{d^2 - b^2} \right] \hat{\phi}$$

\* We again must consider magnetization in region II; following above work,

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

$$= \frac{2I}{c\mu r} \left[ 1 - \frac{r^2 - b^2}{d^2 - b^2} \right] \hat{\phi}$$

Region IV:  $r \geq d$

$$\oint \vec{B} \cdot d\vec{\ell} = \frac{4\pi}{c} \int \vec{j} \cdot d\vec{a}$$

$$B \cdot 2\pi r = \frac{4\pi}{c} \left[ \int \vec{j}_{rod} \cdot d\vec{a}_{rod} + \int \vec{j}_{shell} \cdot d\vec{a}_{shell} \right]$$

$$B \cdot 2\pi r = \frac{4\pi}{c} \left[ \frac{I}{\pi a^2} (\pi a^2) + \frac{-I}{\pi(d^2 - b^2)} (d^2 - b^2)\pi \right]$$

$$\vec{B} = \frac{4\pi}{c} \left[ I - I \right] \cdot \frac{1}{2\pi r}$$

$$\vec{B} = 0$$

$$\vec{H} = \frac{1}{\mu} \vec{B}$$

$$\vec{H} = 0$$

#5 (cont.)

b) \*We know that the self-inductance  $L$  is related to the flux via:  $\Phi = LI$

To get total flux of system, we sum the fluxes from the individual regions

$$\frac{\Phi}{l} = \frac{1}{l} \int_0^l dz \int dx B$$

Region I:  $r \leq a$

$$\begin{aligned} \frac{\Phi}{l} &= \frac{1}{l} \int_0^l dz \int_0^a \frac{2I \sqrt{x^2 + y^2}}{ca^2} dx \\ &= \int_0^l \frac{2I}{ca^2} x dx \\ &= \frac{I}{ca^2} x^2 \Big|_0^a \\ &= \frac{I}{c} \end{aligned}$$

Region II:  $a \leq r \leq b$

$$\begin{aligned} \frac{\Phi}{l} &= \frac{1}{l} \int_0^l dz \int_a^b \frac{2I}{c\sqrt{x^2 + y^2}} dx \\ &= \int_a^b \frac{2I}{cx} dx \\ &= \frac{2I}{c} \ln(x) \Big|_a^b \\ &= \frac{2I}{c} (\ln(b/a)) \end{aligned}$$

Region III:  $b \leq r \leq d$

$$\begin{aligned} \frac{\Phi}{l} &= \frac{1}{l} \int_0^l dz \int_b^d \frac{2I}{cr} \left[ 1 - \frac{r^2 - b^2}{a^2 - b^2} \right] dx \\ &= \int_b^d \frac{2I}{cr} \left[ 1 - \frac{r^2 - b^2}{a^2 - b^2} \right] dx \\ &= \frac{2I}{c} \left[ \frac{1}{x} - \frac{x}{a^2 - b^2} + \frac{b^2}{x(a^2 - b^2)} \right] dx \\ &= \frac{2I}{c} \left[ \ln(x) - \frac{1}{2} \frac{x^2}{a^2 - b^2} + \frac{b^2}{a^2 - b^2} \ln(x) \right] \Big|_b^d \end{aligned}$$

#5 (cont.)

$$\begin{aligned} \text{b) } \frac{\Phi}{\lambda} &= \frac{2I}{c} \left[ \ln\left(\frac{d}{b}\right) - \frac{1}{2} + \frac{b^2}{d^2+b^2} \ln\left(\frac{d}{b}\right) \right] \\ &= \frac{2I}{c} \left[ \left(1 + \frac{b^2}{d^2+b^2}\right) \ln\left(\frac{d}{b}\right) - \frac{1}{2} \right] \end{aligned}$$

Region IV:  $r \geq d$

$$B=0 \Rightarrow \frac{\Phi}{\lambda} = 0$$

\* Thus our overall flux per unit length is:

$$\begin{aligned} \frac{\Phi}{\lambda} &= \frac{I}{c} + \frac{2I}{c} \left( \ln(b/a) \right) + \frac{2I}{c} \left[ \left(1 + \frac{b^2}{d^2+b^2}\right) \ln\left(\frac{d}{b}\right) - \frac{1}{2} \right] \\ &= \frac{I}{c} + \frac{2I}{c} \ln(b/a) + \frac{2I}{c} \ln\left(\frac{d}{b}\right) + \frac{2I}{c} \frac{b^2}{d^2+b^2} \ln\left(\frac{d}{b}\right) - \frac{I}{c} \\ &= \frac{2I}{c} \left[ \ln(b/a) + \left(1 + \frac{b^2}{d^2+b^2}\right) \ln\left(\frac{d}{b}\right) \right] \end{aligned}$$

$$\frac{\Phi}{\lambda} = \frac{LI}{\lambda} \Rightarrow \frac{L}{\lambda} = \frac{\Phi}{LI}$$

$$\Rightarrow \frac{L}{\lambda} = \frac{2}{c} \left[ \ln(b/a) + \left(1 + \frac{b^2}{d^2+b^2}\right) \ln\left(\frac{d}{b}\right) \right]$$

$$\text{c) } U = \frac{1}{8\pi} B^2$$

$$\begin{aligned} U &= \frac{1}{8\pi} \left( \frac{2I}{cr} \right)^2 \\ &= \frac{1}{8\pi} \frac{4I^2}{c^2 r^2} \end{aligned}$$

$$\begin{aligned} U &= \int \frac{1}{4\pi c^2 r^2} r dr d\phi dz \\ &= \frac{2I^2}{c^2} \int_a^b \frac{1}{r} dr \\ &= \frac{2I^2}{c^2} \ln(b/a) \end{aligned}$$

$$\frac{U}{\lambda} = \frac{I^2}{2c^2} \ln(b/a)$$

#5 (cont.)

$$\begin{aligned} d) \quad \frac{1}{2} L_{II} I^2 &= \frac{1}{2} \left( \frac{2}{c} \ln(b/a) \right) I^2 \\ &= \frac{I^2}{c} \ln(b/a) \end{aligned}$$

\* off by factor  $\frac{1}{2c}$ , likely due to units  
being Gaussian v SI