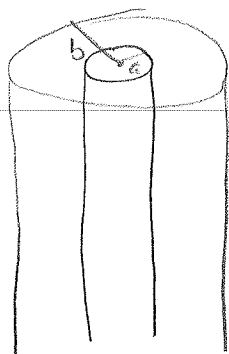


1. A very long conducting wire of radius a , carrying free positive charge per unit length λ , is surrounded by a dielectric coating of outside radius b and relative dielectric constant $\kappa = \epsilon/\epsilon_0$.
 - (a) {2 pts} Find the displacement vector \mathbf{D} everywhere.
 - (b) {2 pts} Find the electric field \mathbf{E} everywhere.
 - (c) {2 pts} Find the polarization density \mathbf{P} everywhere.
 - (d) {2 pts} Find the bound volume charge density ρ_b and the bound surface charge density σ_b everywhere.
 - (e) {2 pts} Show that the total charge densities, bound and free, produce the same \mathbf{E} found in (a).

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E+M #1



⇒ Long conducting wire of radius a w/ free charge per unit length λ

⇒ dielectric coating of radius b , w/ $\kappa = \epsilon/\epsilon_0$

a) Gauss Law states:

$$\oint_S \vec{D} \cdot d\vec{a} = Q_{enc}, \quad Q_{enc} = \lambda L$$

* Long wire means D has azimuthal symmetry \rightarrow points along \hat{r}

$$\Rightarrow \oint_S \vec{D} \cdot d\vec{a} = \lambda L$$

$$D \int_S d\vec{a} = \lambda L$$

$$D 2\pi r L = \lambda L$$

$$D = \frac{\lambda}{2\pi r} \hat{r} \quad (\text{outside wire})$$

b/c our wire is conducting, all the charge flows on the surface

$$\Rightarrow D = \begin{cases} 0 & r < a \\ \frac{\lambda}{2\pi r} \hat{r} & r \geq a \end{cases}$$

assumption of evenly distributed charge means this is an electrostatics problem and any field inside the conductor causes moving charges

b) We know that $D = \epsilon E$

$$\Rightarrow E = \begin{cases} \frac{\lambda}{2\pi\epsilon_0 r} & r > b \\ \frac{\lambda}{2\pi\kappa\epsilon_0 r} & a < r < b \\ 0 & r < a \end{cases}$$

#1 (cont.)

c) We know that $\vec{P} = \vec{D} - \epsilon_0 \vec{E}$ (from $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$)

$$\Rightarrow P = \begin{cases} 0 & r > b \\ \frac{\lambda(1-\frac{1}{k})}{2\pi r} \hat{r} & a < r < b \\ 0 & r < a \end{cases}$$

d) The bound charge is given by: $\rho_b = -\nabla \cdot \vec{P}$

$$\begin{aligned} \rho_b &= -\nabla \cdot \vec{P} \\ &= -\frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\lambda(1-\frac{1}{k})}{2\pi s} \right) \\ &= -\frac{1}{s} \frac{\partial}{\partial s} \left(\frac{\lambda(1-\frac{1}{k})}{2\pi} \right) \\ &= 0 \end{aligned}$$

$\rho_b = 0$ everywhere

The surface bound charge $\sigma_b = \vec{P} \cdot \hat{n}$

$$\hookrightarrow \text{at } s=b, P = \frac{\lambda(1-\frac{1}{k})}{2\pi b} \hat{s} \cdot \hat{n}, \hat{n} = \hat{s}$$

$$\Rightarrow \sigma_b|_{s=b} = \frac{\lambda(1-\frac{1}{k})}{2\pi b}$$

$$\sigma_b|_{s=a} = -\frac{\lambda(1-\frac{1}{k})}{2\pi a} \quad (\text{by similar logic as above, } \hat{n} = -\hat{s})$$

e) According to Gauss Law: $\int \vec{E} \cdot d\vec{a} = Q_{enc}/\epsilon_0$

$$\text{* if } s > b: Q_{enc} = \lambda L - \frac{\lambda(1-\frac{1}{k})}{2\pi a} (2\pi a L) + \frac{\lambda(1-\frac{1}{k})}{2\pi b} (2\pi b L)$$

$$= \lambda L - \lambda L(1-\frac{1}{k}) + \lambda L(1-\frac{1}{k})$$

$$= \lambda L$$

$$\hookrightarrow \vec{E} = \frac{\lambda}{2\pi r \epsilon_0} \hat{r}$$

$$\text{* if } a < s < b: Q_{enc} = \lambda L - \frac{\lambda(1-\frac{1}{k})}{2\pi a} (2\pi a L)$$

$$= \lambda L - (1-\frac{1}{k})\lambda L$$

$$= \frac{\lambda L}{k}$$

$$\hookrightarrow = \frac{\lambda}{2\pi r k \epsilon_0} \hat{r}$$

$$\text{* if } s < a: Q_{enc} = 0$$

$$\hookrightarrow \vec{E} = 0$$

All fields match!