

5. A plane-polarized harmonic ( $e^{-i\omega t}$ ) plane electromagnetic wave traveling to the right in a homogeneous dielectric medium described by an dielectric constant  $\epsilon_1$ , strikes a second homogeneous dielectric material described by dielectric constant  $\epsilon_2 > \epsilon_1$  (see the figure). Assume that both materials have the same magnetic permeability  $\mu_0$  and that the incidence angle is  $0^\circ$  (i.e., the wave is traveling perpendicular to the junction). Assume the incoming wave is polarized in the  $\hat{x}$  direction and that its electric field amplitude is  $E_0$ , i.e., assume the incoming electric field is the real part of

$$\mathbf{E} = E_0 e^{i(kz - \omega t)} \hat{x}.$$

- (a) {3 pts} Give the magnetic induction  $\mathbf{B}$  associated with the above incoming wave. Make sure your wave satisfies Maxwell's equations, e.g., give  $k$  as a function of  $\omega$ , the direction of  $\mathbf{B}$ , and the amplitude of  $\mathbf{B}$  as a function of  $E_0$ .
- (b) {1 pts} Give similar expressions for the  $\mathbf{E}$  and  $\mathbf{B}$  components of the reflected and transmitted waves. Use  $E_0''$  and  $E_0'$  for the respective amplitudes of reflected and transmitted waves.
- (c) {2 pts} In general, what conditions must be satisfied at the junction between two materials by the electromagnetic fields  $\mathbf{E}$ ,  $\mathbf{B}$ ,  $\mathbf{D}$ , and  $\mathbf{H}$ , if Maxwell's equations are to be satisfied?
- (d) {2 pts} Apply these junction conditions to the combined incoming, reflected, and transmitted wave to compute  $E_0''$  and  $E_0'$  as functions of  $E_0$  and the two dielectric constants  $\epsilon_1$  and  $\epsilon_2$ .
- (e) {2 pts} Evaluate the time averages of the Poynting vectors of the incident, reflected, and transmitted waves. Recall that

$$\mathbf{S} \equiv \mathbf{E} \times \mathbf{H}, \quad (SI)$$

$$\mathbf{S} \equiv \frac{1}{4\pi} \mathbf{E} \times \mathbf{H}. \quad (Gaussian)$$

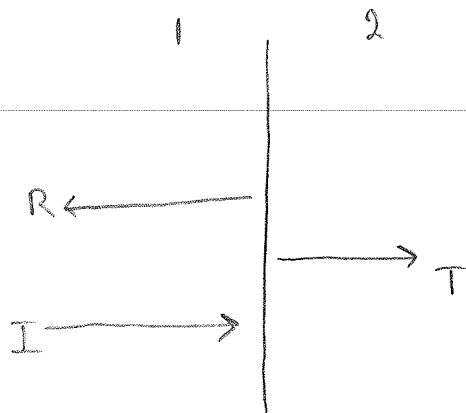
The sum of the magnitudes of the reflected and transmitted time averaged Poynting vectors should equal the magnitude of the incident wave's time averaged Poynting vector.

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# E+M #5

Gaussian

a)



$$\vec{E} = E_0 \exp[i(k_z z - \omega t)] \hat{x}$$

$$\vec{B} = n \hat{k} \times \vec{E}$$

$$= \sqrt{\epsilon_1} E_0 \exp[i(k_z z - \omega t)] \hat{y}$$

\* The relevant Maxwell equations in Gaussian units are:

$$\textcircled{1} \nabla \cdot \vec{D} = 4\pi \rho_f$$

$$\textcircled{3} \nabla \cdot \vec{B} = 0$$

$$\textcircled{2} \nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$$

$$\textcircled{4} \nabla \times \vec{H} - \frac{1}{c} \frac{\partial \vec{D}}{\partial t} = \frac{4\pi}{c} \vec{J}_f$$

\* Starting w/ equation  $\textcircled{3}$ :

$$\nabla \cdot \vec{B} = \frac{\partial}{\partial x} B_x + \frac{\partial}{\partial y} B_y + \frac{\partial}{\partial z} B_z$$

$$= \frac{\partial}{\partial y} (-\sqrt{\epsilon_1} E_0 \exp[i(k_z z - \omega t)])$$

$$= 0 \checkmark$$

\* Now equation  $\textcircled{2}$ :

$$\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$$

$$\hookrightarrow \nabla \times \vec{E} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & 0 & 0 \end{vmatrix} = \langle 0, +\frac{\partial}{\partial z} E_x, -\frac{\partial}{\partial y} E_x \rangle$$

$$= \langle 0, ik_z E_0 \exp[i(k_z z - \omega t)], 0 \rangle$$

# #5 (cont.)

$$\begin{aligned} \text{a) } \frac{1}{c} \frac{\partial}{\partial t} B &= \frac{1}{c} \frac{\partial}{\partial t} (-\sqrt{\epsilon_1} E_0 \exp[i(kz - \omega t)]) \hat{y} \\ &= -\frac{1}{c} \sqrt{\epsilon_1} i\omega E_0 \exp[i(kz - \omega t)] \hat{y} \end{aligned}$$

$$\Rightarrow \nabla \times \vec{E} + \frac{1}{c} \frac{\partial B}{\partial t} \stackrel{?}{=} 0$$

$$= -ik E_0 \exp[i(kz - \omega t)] + -i \frac{\omega}{c} \sqrt{\epsilon_1} E_0 \exp[i(kz - \omega t)]$$

$$= -i E_0 \exp[i(kz - \omega t)] \left( -k + \frac{\omega}{c} \sqrt{\epsilon_1} \right)$$

$$\hookrightarrow \text{Only true if } k = \frac{\omega}{c} \sqrt{\epsilon_1}$$

\* Finally equation ①

$$\nabla \times H - \frac{1}{c} \frac{\partial D}{\partial t} = \frac{4\pi}{c} \vec{J}_f$$

$$\vec{J}_f = 0$$

$$\nabla \times H = \nabla \times \frac{1}{\mu_0} \vec{B}$$

$$= \frac{1}{\mu_0} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \partial_x & \partial_y & \partial_z \\ 0 & B_y & 0 \end{vmatrix} = \langle -\partial_z B_y, 0, \partial_x B_y \rangle$$

$$= \frac{1}{\mu_0} \frac{\partial}{\partial z} (\sqrt{\epsilon_1} E_0 \exp[i(kz - \omega t)]) \hat{x}$$

$$= \frac{\sqrt{\epsilon_1}}{\mu_0} E_0 (ik) \exp[i(kz - \omega t)] \hat{x}$$

$$-\frac{1}{c} \frac{\partial D}{\partial t} = -\frac{1}{c} \frac{\partial}{\partial t} \epsilon_1 \vec{E}$$

$$= -\frac{\epsilon_1}{c} (-i\omega) E_0 \exp[i(kz - \omega t)] \hat{x}$$

$$\hookrightarrow \nabla \times H - \frac{1}{c} \frac{\partial D}{\partial t} = \frac{i\omega \epsilon_1}{c} E_0 \exp[i(kz - \omega t)] \hat{x} + \frac{\sqrt{\epsilon_1} k}{\mu_0} E_0 \exp[i(kz - \omega t)] \hat{x}$$

$$= 0 \quad \text{if} \quad \frac{\omega \epsilon_1}{c} - \frac{\sqrt{\epsilon_1} k}{\mu_0} = 0 \quad \Rightarrow \quad k = \frac{\omega \sqrt{\epsilon_1}}{c} \quad \text{as before} \quad \checkmark$$

### #5 (cont.)

b) \* For the transmitted wave:

$$\vec{E}' = E_0' \exp[i(k_2 z - \omega t)] \hat{x}$$

$$\vec{B}' = \sqrt{\epsilon_2} E_0' \exp[i(k_2 z - \omega t)] \hat{y}$$

$$k_2 = \frac{\omega \sqrt{\epsilon_2}}{c}$$

\* For the reflected waves

$$\vec{E}'' = E_0'' \exp[i(k_1 z - \omega t)] \hat{x}$$

$$\vec{B}'' = -\sqrt{\epsilon_1} E_0'' \exp[i(k_1 z - \omega t)] \hat{y}$$

$$k_1 = \frac{\omega \sqrt{\epsilon_1}}{c}$$

c) The general boundary conditions are:

$$-D_1^\perp - D_2^\perp = 4\pi\sigma_f$$

$$B_1^\perp - B_2^\perp = 0$$

$$E_1^\parallel = E_2^\parallel$$

$$H_1^\parallel - H_2^\parallel = 4\pi\vec{K}_f$$

d) Applying these boundary conditions to our situation

$$E_1'' - E_2'' = 0$$

$$(\vec{E}_I + \vec{E}_R) - \vec{E}_T = 0$$

$$[(E_I + E_R) - E_T] \times \hat{z} = 0 \Rightarrow \text{picks out } \hat{y} \text{ components}$$

$$-E_0 \exp[i(k_1 z - \omega t)] + E'' \exp[i(k_1 z - \omega t)] - E' \exp[i(k_2 z - \omega t)] = 0$$

\* if we set our boundary to be  $z=0$

$$-E_0 + E'' - E' = 0$$

#5 (cont.)

$$d) \quad H_1'' - H_2'' = \cancel{4\pi \vec{K}_f} \rightarrow 0$$

$$\frac{1}{\mu_0} (B_1'' - B_2'') = 0 \Rightarrow B_1'' - B_2'' = 0$$

$$\hookrightarrow B_1'' - B_2'' = 0$$

$$B_I'' + B_R'' - B_T'' = 0$$

$$[(B_I + B_R) - B_T] \times \hat{z} = 0 \Rightarrow \text{picks off } x\text{-component}$$

$$-E_0 \sqrt{\epsilon_1} \exp[i(k_z z - \omega t)] + E_0'' \sqrt{\epsilon_1} \exp[i(-k_z z - \omega t)] - E_0' \sqrt{\epsilon_2} \exp[i(k_z z - \omega t)]$$

$$(E_0 + E_0'') \sqrt{\epsilon_1} - \sqrt{\epsilon_2} E_0' = 0$$

$$\hookrightarrow E_0' = \sqrt{\frac{\epsilon_1}{\epsilon_2}} (E_0 + E_0'')$$

\*Using our equations to solve for  $E_0'$ ,  $E_0''$  in terms of  $E_0$ :

$$E_0' = \sqrt{\frac{\epsilon_1}{\epsilon_2}} (-E_0 + E_0'')$$

$$= \sqrt{\frac{\epsilon_1}{\epsilon_2}} (-E_0 + (E_0' + E_0))$$