

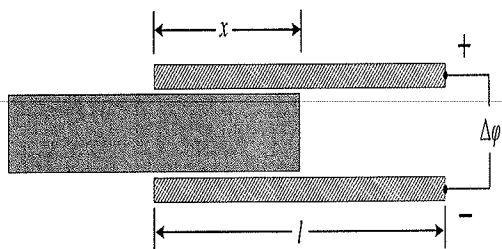
E & M Qualifier

January 13, 2011

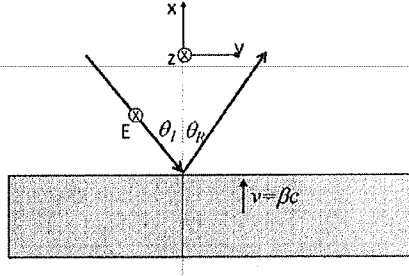
To insure that the your work is graded correctly you MUST:

1. use only the blank answer paper provided,
2. write only on one side of the page,
3. put your alias on every page,
4. put the problem # on every page,
5. start each problem by stating your units e.g., SI or Gaussian,
6. number every page starting with 1 for each problem,
7. put the total # of pages you use for that problem on every page,
8. staple your exam when done.

Use only the reference material supplied (Schaum's Guides).



1. A parallel plate capacitor has the region between its plates filled with a dielectric slab of dielectric constant $K = \epsilon/\epsilon_0$ and mass m . The plate dimensions are: width w , length ℓ , and plate separation d . The capacitor plates are connected to a battery of constant voltage V ($\Delta\phi = V$ in the figure). Neglect the fringe field and friction, and assume the slab is constrained to move in the plane parallel to the capacitor plates.
 - (a) {2 pts} Compute the capacitance $C \equiv q/V$ of this capacitor as a function of x .
 - (b) {2 pts} If the slab is withdrawn half way (to $x = \ell/2$) and held in place, what is the magnitude and direction of the force on the slab caused by the electric field?
 - (c) {2 pts} At $x = \ell/2$ the slab is released and given a velocity v_0 to the right. Find the current supplied by the battery at the instant it is released.
 - (d) {2 pts} At $x = \ell/2$ the slab is again released but with zero velocity. Describe the motion of the slab (in words). What is the maximum velocity achieved by the slab?
 - (e) {2 pts} Sketch the displacement of the slab versus time.



2. This problem investigates the shifting frequency of electromagnetic radiation that is reflected off a moving target. Incident and reflected frequencies and angles are not the same if the target is moving.

Assume that in the lab frame of reference, the target is a flat mirror traveling upward in the positive x -direction parallel to the mirror's normal with velocity $\mathbf{v} = \beta c \hat{\mathbf{x}}$ (see the figure). Also assume the wave is a linearly polarized plane wave traveling in vacuum towards the moving mirror at angle θ_I (relative to the mirror's normal). If the polarization is in the $\hat{\mathbf{z}}$ direction, the incident electric field is given by

$$\mathbf{E}_I = E_0 \hat{\mathbf{z}} e^{i(\mathbf{k}_I \cdot \mathbf{r} - \omega_I t)},$$

with

$$\mathbf{k}_I = \frac{\omega_I}{c} (-\cos \theta_I \hat{\mathbf{x}} + \sin \theta_I \hat{\mathbf{y}}).$$

- (a) {2 pts} Write the Lorentz boost A as a function of β and $\gamma \equiv \sqrt{1 - \beta^2}$ that transforms the Lab coordinates \mathbf{r} and ct to coordinates \mathbf{r}' and ct' co-moving with the mirror. Also give the inverse A^{-1} of the Lorentz boost A that transforms the moving coordinates \mathbf{r}' and ct' into Lab coordinates \mathbf{r} and ct .
- (b) {3 pts} By rewriting the above wave's phase in both reference frames, i.e.,

$$\mathbf{k}_I \cdot \mathbf{r} - \omega_I t = \mathbf{k}'_I \cdot \mathbf{r}' - \omega'_I t'$$

as a function of the co-moving mirror coordinates \mathbf{r}' and ct' (i.e., use A^{-1}) find \mathbf{k}'_I and ω'_I as observed in the co-moving frame. These will be functions of β, γ , and θ_I as well as ω_I .

- (c) {2 pts} By writing the incident wave vector just obtained in the moving frame in the form

$$\mathbf{k}'_I = \frac{\omega'_I}{c} (-\cos \theta'_I \hat{\mathbf{x}} + \sin \theta'_I \hat{\mathbf{y}}),$$

determine the incident angle θ'_I as seen by observers moving with the mirror (e.g., give $\cos \theta'_I$ as a function of θ_I, ω_I and the Lorentz parameters β, γ).

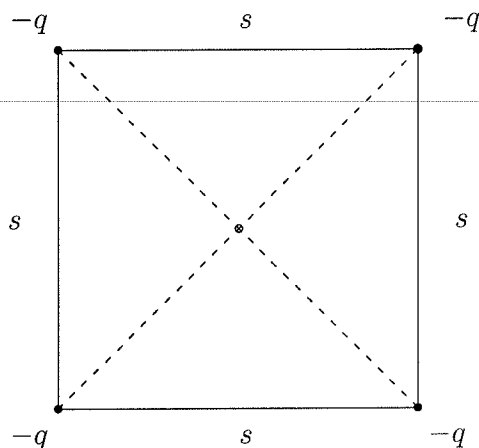
- (d) {3 pts} If, as seen by observers moving with the mirror, the reflected wave has the same frequency as the incident wave $\omega'_R = \omega'_I$ and a reflection angle that is the same as the incidence angle $\theta'_R = \theta'_I$, i.e.,

$$\mathbf{k}'_R = \frac{\omega'_I}{c}(\cos \theta'_I \hat{\mathbf{x}} + \sin \theta'_I \hat{\mathbf{y}}),$$

what is the frequency ω_R of the reflected light as measured in the laboratory frame? Hint: again use

$$\mathbf{k}_R \cdot \mathbf{r} - \omega_R t = \mathbf{k}'_R \cdot \mathbf{r}' - \omega'_R t',$$

and the Lorentz boosts A .



3. Consider a square with sides of length s and charges $-q$ at the corners as shown:
- (a) {2 pts} What is the potential at the center of the square if the potential is zero at ∞ ?
 - (b) {2 pts} How much work does it take to bring in another charge $-q$ from ∞ to the center of the square?
 - (c) {3 pts} How much work does it take to assemble the original configuration of 4 negative charges (no charge at center)?
 - (d) {3 pts} Now suppose that instead of the 4 charges being located at the corners of a square, a net charge of $-4q$ is distributed uniformly on the surface of a sphere of radius s . How much work does it take to bring in another charge q from ∞ to the center of the sphere?

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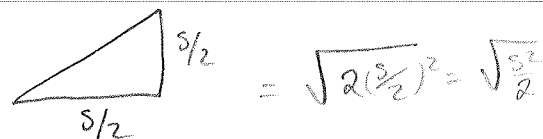
E+M #3

Gaussian

a) We will assume the origin lies at the center of the square

$$\Rightarrow V = \sum_i \frac{q_i}{r_i}$$

$$= 4 \frac{-q\sqrt{2}}{s}$$



$$\begin{aligned} \text{b) } W &= \int \vec{F} \cdot d\vec{\ell} \\ &= q (V(b) - V(a)) \\ &= -q \left(-\frac{\sqrt{2}q}{s} \right) \\ &= \frac{\sqrt{2}q^2}{s} \end{aligned}$$

c) The work necessary to assemble the charge distribution is:

$$\begin{aligned} W &= \frac{1}{2} \sum_i q_i V_i \quad \text{where } i \text{ is each individual charge} \\ &= \frac{1}{2} \left(0 + (-q) \left(-\frac{q}{s} \right) + (-q) \left(-\frac{q}{s} + \frac{-q}{\sqrt{2}s} \right) + -q \left(-\frac{q}{s} + \frac{-q}{s} + \frac{-q}{\sqrt{2}s} \right) \right) \\ &= \frac{1}{2} \left(\frac{q^2}{s} + \frac{q^2}{s} + \frac{q^2}{\sqrt{2}s} + \frac{q^2}{s} + \frac{q^2}{s} + \frac{q^2}{\sqrt{2}s} \right) \\ &= \frac{2q^2}{s} + \frac{q^2}{\sqrt{2}s} \end{aligned}$$

d) From symmetry, and the fact that the charge is distributed evenly over the surface of the sphere,

$$\vec{E} = \begin{cases} -\frac{4q}{r^2} & r > s \\ 0 & r < s \end{cases}$$

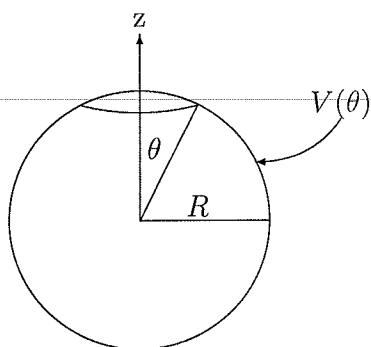
* We also know that

$$V = - \int_{\infty}^0 \vec{E} \cdot d\vec{\ell} = - \left[\int_{\infty}^s \vec{E}_{\text{out}} \cdot d\vec{\ell} - \int_s^0 \vec{E}_{\text{in}} \cdot d\vec{\ell} \right]$$

#3 (cont.)

$$\begin{aligned} \text{d) } V &= - \int_{\infty}^s \frac{-4q}{r^2} dr \\ &= - \left[-4q \left(\frac{-1}{r} \Big|_{\infty}^s \right) \right] \\ &= 4q \left(\frac{-1}{s} - \frac{-1}{\infty} \right) \\ &= \frac{-4q}{s} \end{aligned}$$

$$\begin{aligned} W &= -q V(0) \\ &= -q \left(\frac{-4q}{s} \right) \\ &= \frac{4q^2}{s} \end{aligned}$$



4. Consider an isolated spherical surface of radius R centered on the origin, that is kept at a known potential $V(\theta)$, i.e.,

$$\Phi(r = R, \theta) = V(\theta)$$

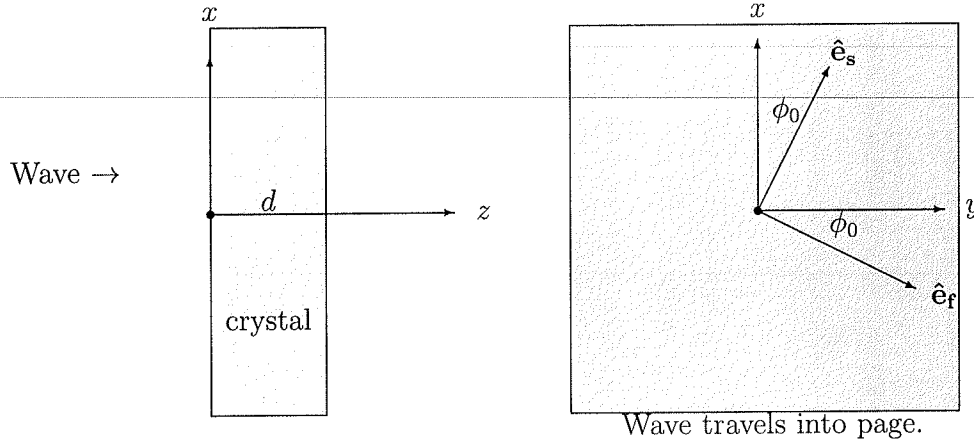
where (r, θ, ϕ) are the usual spherical polar coordinates, i.e., θ is measured with respect to a z -axis passing through the center of the sphere and ϕ is the azimuthal angle about the z -axis measured from the x axis.

- (a) {2 pts} Write down expressions for the general solution to $\nabla^2\Phi(r, \theta) = 0$ for the electrostatic potential as a linear combination of Legendre polynomials in the respective regions $0 \leq r < R$ and $r > R$. Assume that the potential vanishes at $r \rightarrow \infty$ and has azimuthal symmetry i.e., no dependence on the angle ϕ . Do not include terms that must vanish. Do not attempt to evaluate the constants that appear in the linear combination but do give the correct r dependence of each term.
- (b) {2 pts} What boundary conditions must your two expressions satisfy at the junction $r = R$ to have a unique solution to Maxwell's equations?
- (c) {2 pts} If the particular surface potential imposed is

$$\Phi(r = R, \theta) = V_0 \cos \theta$$

where V_0 is a constant, what is the explicit form of your potential for both regions $r \leq R$ and $r > R$?

- (d) {2 pts} Determine the resulting electric field on both sides of the $r=R$ surface.
- (e) {2 pts} What is the surface charge density $\sigma(\theta)$ on the spherical shell at $r=R$.



5. A plane polarized monochromatic light wave traveling in the $+z$ direction enters a large flat slab of transparent crystal of thickness d , located between $z = 0$ and $z = d$. This crystal has the property that the index of refraction depends on the direction of polarization as follows: Plane waves traveling in the z direction but polarized in the direction

$$\hat{\mathbf{e}}_s = \cos \phi_0 \hat{\mathbf{x}} + \sin \phi_0 \hat{\mathbf{y}},$$

travel with speed $v_s = c/n_s < c$ but those polarized in the orthogonal direction

$$\hat{\mathbf{e}}_f = -\sin \phi_0 \hat{\mathbf{x}} + \cos \phi_0 \hat{\mathbf{y}},$$

travel with the faster speed $v_f = c/n_f < c$ where $n_s = n_f + \delta n$.

Assume the wave, just after entering the crystal (i.e., for very small $z \ll \lambda < d$), is polarized in the y direction and hence has the form

$$\mathbf{E}(z \approx 0, t) = E_0 \hat{\mathbf{y}} e^{-i\omega t}.$$

- (a) {4 pts} Prove that in general the initial plane wave becomes elliptically polarized when it reaches $z = d$ by deriving the following expression

$$\mathbf{E}(z = d, t) = [E_x \hat{\mathbf{x}} + E_y \hat{\mathbf{y}}] e^{i(\bar{k}d - \omega t)},$$

where

$$\bar{k} \equiv \frac{\omega}{c} \left(\frac{n_s + n_f}{2} \right),$$

and

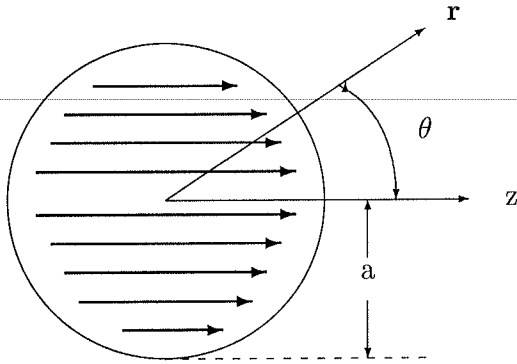
$$\begin{aligned} E_x &= iE_0 \sin 2\phi_0 \sin \delta, \\ E_y &= E_0 (\cos \delta - i \cos 2\phi_0 \sin \delta), \end{aligned}$$

with

$$\delta \equiv \frac{\omega d}{2c} \delta n.$$

Hint: Write the wave at $z=0$ as a combination of slow and fast plane polarized parts using $\hat{\mathbf{y}} = \sin \phi_0 \hat{\mathbf{e}}_s + \cos \phi_0 \hat{\mathbf{e}}_f$.

- (b) {3 pts} For what values of δ and θ_0 will the wave emerge from the crystal as a circularly polarized wave? ($E_x/E_y = \pm i$).
-
- (c) {3 pts} For what minimum crystal thicknesses $d = d_{min}$ will the wave emerge as a plane polarized wave ($E_x/E_y = \text{real}$) and what will its polarization direction be?



6. A permanent magnet in the shape of a solid sphere of radius a is oriented on the z -axis as shown in the figure. The magnetization of the magnet is given by $\vec{M} = M_0 \hat{z}$. [Recall that $\nabla \times \mathbf{H} = 0$ implies the existence of a magnetic scalar potential $\Phi_m(r, \theta)$ related to the magnetic field by $\mathbf{H} = -\vec{\nabla} \Phi_m(r, \theta)$.]

- (a) {4 pts} Compute the scalar magnetic potential $\Phi_m(r, \theta)$ at all points $r < a$ and $r > a$.
- (b) {3 pts} Compute the magnetic Field $\mathbf{H} = -\vec{\nabla} \Phi_m(r, \theta)$ at all points $r < a$ and $r > a$.
- (c) {3 pts} Compute the magnetic induction \mathbf{B} , where

$$\begin{aligned} \mathbf{B}/\mu_0 &= \mathbf{H} + \mathbf{M}, & (SI) \\ \mathbf{B} &= \mathbf{H} + 4\pi\mathbf{M}, & (Gaussian) \end{aligned}$$

at all points $r < a$ and $r > a$.

Hints: The magnetic potential is axial symmetric about the z -axis and satisfies the Laplace equation at all points except $r = a$. Legendre polynomials are useful.