

Stellar Nucleosynthesis

I The nature of stars & the Sun

1. ~ 1 million miles in diameter
2. $\sim 91\%$ H $\sim 9\%$ He $\sim 1\%$ other "metals"
(% of total number of atoms)

3. The core

Size $\sim \frac{1}{4} R_0$

density $160 \times 10^3 \text{ kg/m}^3$

Temp. 14 million K

II Energy Released by a Nuclear Reaction

found from mass defect

= mass of particles going into reaction - mass of particles created

Proton-Proton chain - more later



$$4.0312 - 4.0026 = .0286 \text{ u} = 26.6 \text{ MeV}$$

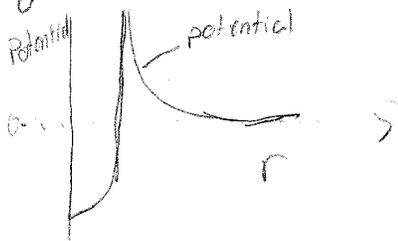
III Reactions - Rates + Probabilities

1 Most reactions in the sun require tunneling

p + nucleus means that there is a coulomb barrier with potential energy $\frac{e^2}{4\pi\epsilon_0 r}$. For fusion to take place, the $\frac{1}{2}mv^2$ kinetic energy must exceed the potential at small r

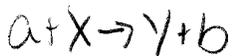
$$r \leq \frac{2e^2}{4\pi\epsilon_0 m_p v^2} \quad \frac{1}{2}mv^2 = \frac{3}{2}kT \quad v = \left(\frac{3kT}{m}\right)^{1/2}$$

v must be very large so T must be huge



2. Reaction Rates

Particle a collides with nucleus X to produce nucleus Y and particle b



$dn_a(v_a)$ number of particles a in 1cm^3 in volume element d^3v_a
in the velocity space

$dn_x(v_x)$

$v = v_a - v_x$ relative velocity for 2 particles

$\sigma(v)$ reaction crosssection (Probability of interaction between the incident particle and the target nucleus)

r_{ax} number of nuclear reactions in 1cm^3 per second

$$r_{ax} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \sigma(v) v dn_a(v) dn_x(v) \quad (1)$$

The speed of the particles follow the Maxwell distribution

$$dn_i(v_i) = n_i \left(\frac{m_i}{2\pi kT} \right)^{3/2} e^{-\frac{m_i v_i^2}{2kT}} d^3v_i \quad (2) \quad \text{put (2) into (1)}$$

$$r_{ax} = n_a n_x \left(\frac{m_a}{2\pi kT} \right)^{3/2} \left(\frac{m_x}{2\pi kT} \right)^{3/2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\left(\frac{m_a v_a^2}{2kT} + \frac{m_x v_x^2}{2kT} \right)} \sigma(v) v d^3v_a d^3v_x \quad (3)$$

if U is the velocity of the COM of the two particles

$$(m_a + m_x)U = m_a v_a + m_x v_x \quad v = v_a - v_x$$

$$\Rightarrow v_a = U + \frac{m_x}{m_a + m_x} v \quad v_x = U - \frac{m_a}{m_a + m_x} v$$

$$\mu = \frac{m_x m_a}{m_a + m_x} \quad \frac{1}{2} m_a v_a^2 + \frac{1}{2} m_x v_x^2 = \frac{1}{2} (m_a + m_x) U^2 + \frac{1}{2} \mu v^2 \quad (4)$$

$$\text{Put (4) into (3)} \quad \text{and } d^3v_a \cdot d^3v_b = d^3U \cdot d^3v = 4\pi U^2 dU \cdot 4\pi v^2 dv$$

$$r_{ax} = n_a n_x \left(\frac{m_a}{2\pi kT} \right)^{3/2} \left(\frac{m_x}{2\pi kT} \right)^{3/2} \int_0^{\infty} \int_0^{\infty} e^{-\left(\frac{(m_a + m_x) U^2}{2kT} + \frac{\mu v^2}{2kT} \right)} \sigma(v) v 4\pi U^2 dU \cdot 4\pi v^2 dv \quad (5)$$

$$\int_0^{\infty} e^{-\frac{(m_a + m_x) U^2}{2kT}} 4\pi U^2 dU = \left(\frac{2\pi kT}{m_a + m_x} \right)^{3/2}$$

$$r_{ax} = 4\pi n_a n_x \left(\frac{\mu}{2\pi kT} \right)^{3/2} \int_0^{\infty} e^{-\frac{\mu v^2}{2kT}} \sigma(v) v^3 dv \quad (6)$$

$E = \frac{1}{2} \mu v^2$ (6) becomes

$$r_{ax} = n_a n_x \int_0^\infty f(E) \sigma v dE = n_a n_x \langle \sigma v \rangle \quad (7)$$

$$\langle \sigma v \rangle = \int_0^\infty f(E) \sigma v dE \quad f(E) = \frac{2}{\pi^{1/2}} \frac{1}{(kT)^{3/2}} e^{-\frac{E}{kT}} E^{1/2}$$

if a and x are the same then the probability of having a collision $\propto \frac{n_a(n_a-1)}{2} \approx \frac{n_a^2}{2} = \langle \sigma v \rangle$

$$\text{so } r_{ax} = \frac{1}{1 + \delta_{ax}} n_a n_x \langle \sigma v \rangle$$

3. Reaction Probability $\langle \sigma v \rangle$

$$\sigma v = \frac{r_{ax}}{n_a n_x} \quad \text{dn}_i \text{ number of particles in } 1 \text{ cm}^3$$

r_{ax} is nuclear reaction rate
= product of two probabilities $P_{\text{coll}}, P_{\text{nu}}$

P_{coll} = probability of tunneling = ()

P_{nu} = probability of the particle having a reaction in the nuclear force range

$$P_{\text{coll}} = \left(\frac{2\mu}{m}\right)^{1/2} E^{1/2} e^{-B/E^{1/2}} \quad B = \frac{4\pi^2 Z_1 Z_2 e^2}{h} \left(\frac{m}{2}\right)^{1/2}$$

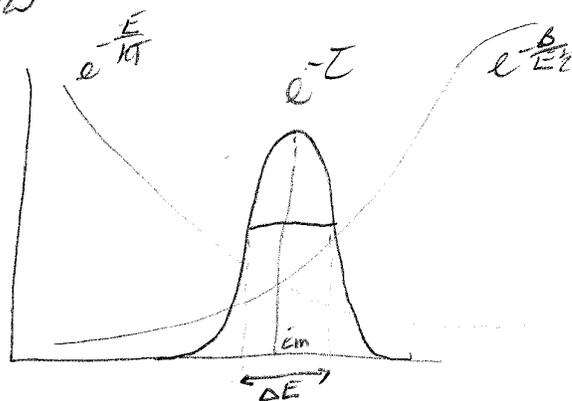
$\sigma v \propto P_{\text{coll}}$ for non-resonance

$$\langle \sigma v \rangle \propto \left(\frac{2}{\pi m}\right)^{1/2} \frac{2}{(kT)^{3/2}} \int_0^\infty e^{-\frac{E}{kT} - \frac{B}{E^{1/2}}} dE$$

$$\tau = \frac{3E_m}{kT} \quad \Delta E = \frac{\sqrt{8} E_m}{\tau^2}$$

$$\int e^{-\frac{E}{kT} - \frac{B}{E^{1/2}}} dE \approx \sqrt{8} \frac{E_m}{\tau^2} e^{-\tau}$$

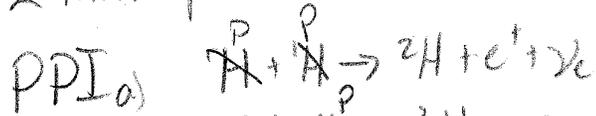
$$\langle \sigma v \rangle \propto \frac{1}{T^{3/2}} e^{-\tau}$$



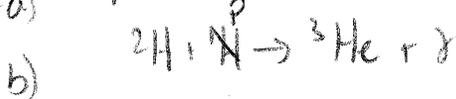
The integrand is known as the Gamow Peak
 $E = E_m$ at the peak
 ΔE is the half width

H-burning - can start at 10 million K

2 main processes



(releases 1.442 MeV) $\sim 10^6$ billion yrs



5.49 MeV

~ 6 s



12 σ_0 MeV

10^8 yrs

+ 26.72 MeV



1.02 MeV

instant

(26.72 MeV / 2)

a) is weak $\sigma < 10^{-51} \text{ m}^2$ (10^{-24} fm^2) at $T \sim 1.5 \times 10^7 \text{ K}$
 too small to be measured, calculated from theory

b)+c) involve Coulomb barrier penetration or much much larger

b) EM c) strong

At these densities (10^5 kg/m^3) d has a mean life of 1s

${}^3_2\text{He}$ " " " 2×10^5 yrs

reaction A controls the production of He
 this chain accounts for 91% of reactions in the sun

PPII

a)

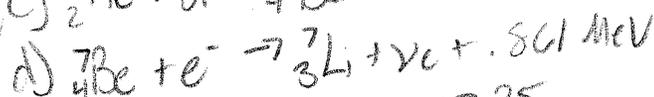
b)



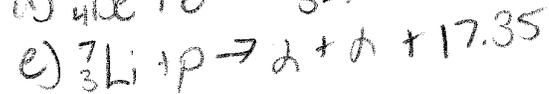
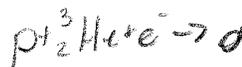
Beryllium

EM

.8 MeV / 2 lost



weak

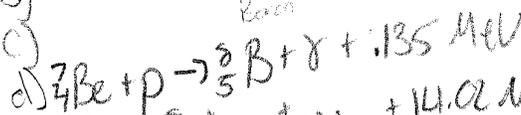


PPIII

a)

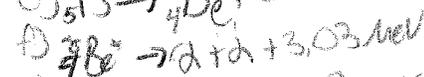
b)

c)

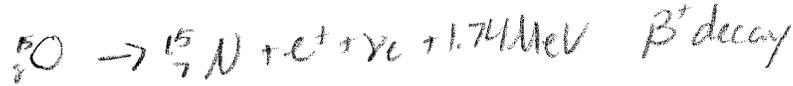
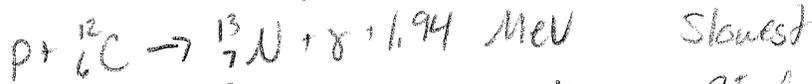


Boron

7.2 MeV / 2

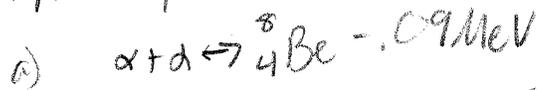


CNO important for high mass stars $> 8 M_{\odot}$



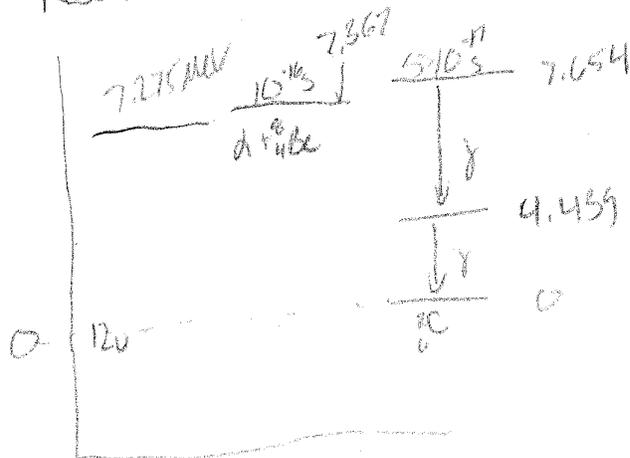
He burning 100 Million K 10^8 kg/m^3

triple α process



${}^8_4\text{Be}$ has a mean life of 10^{-16} s so a) leads to an equilibrium concentration of ${}^8_4\text{Be} \rightarrow N_d \sim 10^{-9} {}^8_4\text{Be}$

Resonance in reaction b) at COM Energy @ 207 KeV $\rightarrow {}^{12}_6\text{C}$ excited \rightarrow photoemission



This is the last thing the sun can produce. it will end its life as a C+O white Dwarf

For more massive stars ...

