

Chapter 7

Wavefunctions

The transformation function $\langle a'|b'\rangle$ tells how to go from one description (a states) to another (b states)—from one class of states to another. But we don't have to think of all states in a class; we can talk of two states only:

$$b' \rightarrow 2, \quad a' \rightarrow 1. \quad (7.1)$$

Then the probability that if we have selected 2 of subsequently finding 1 (or *vice versa*) is

$$p(1, 2) = |\langle 1|2\rangle|^2, \quad (7.2)$$

where $\langle 1|2\rangle$ is the (inner) product of the vector representing state 1 with the vector representing state 2. It tells how alike or different the two states are.

Now suppose we describe the system in terms of A measurements. We then use the algebraic construction of unity,

$$1 = \sum_{a'} |a'\rangle\langle a'|. \quad (7.3)$$

This carries us from vectors to components of vectors in some coordinate system:

$$\langle 1|2\rangle = \langle 1| \sum_{a'} |a'\rangle\langle a'|2\rangle = \sum_{a'} \langle 1|a'\rangle\langle a'|2\rangle. \quad (7.4)$$

To reiterate, since

$$|2\rangle = \sum_{a'} |a'\rangle\langle a'|2\rangle, \quad (7.5)$$

we can think of $|a'\rangle$ as a unit vector in some coordinate system, and $\langle a'|2\rangle$ as the component of the vector $|2\rangle$ in that coordinate system, the projection of $|2\rangle$ on the basis vector $|a'\rangle$. Since

$$\langle 1|2\rangle = \sum_{a'} \langle a'|1\rangle^* \langle a'|2\rangle, \quad (7.6)$$

we have here a complex scalar product, the sum of the products of components with complex conjugate components. The set of components are also denoted by

$$\langle a' | \rangle = \psi(a'), \quad (7.7)$$

where $| \rangle$ represents any state, which is the *wavefunction* of that state in the A description. Thus the probability of finding the state $|1\rangle$ given that the system was prepared in the state $|2\rangle$ is

$$p(1, 2) = \left| \sum_{a'} \psi_1(a')^* \psi_2(a') \right|^2. \quad (7.8)$$

This is the form we saw much earlier when we discussed spin-1/2 system. See Sec. 4.2. Note that this probability makes no reference to A measurements, so is independent of the “coordinate system.”

The probability that if the system is known to be in state 2, and measurement of A will yield the value a' is

$$p(a', 2) = |\langle a' | 2 \rangle|^2 = |\psi_2(a')|^2. \quad (7.9)$$

In general, if ψ is the wavefunction of the state in the A description, $|\psi(a')|^2$ is the probability of finding $A = a'$ in that state. Note the factorization referred to earlier:

- ψ refers to how the system is prepared,
- a' refers to what particular measurement is performed on that state.

The two probability statements (7.8) and (7.9) are not independent. For if

$$\langle 1 | = \langle a' |, \quad |1\rangle = |a'\rangle, \quad (7.10)$$

then

$$\psi_1(a'') = \langle a'' | a' \rangle = \delta(a'', a'), \quad (7.11)$$

and so Eq. (7.8) implies Eq. (7.9).

Finally, we note that the vector representing a state is a unit vector,

$$\langle 1 | 1 \rangle = 1, \quad (7.12)$$

which physically expresses the fact that if we initially measure the system to be in state 1, it will be found with certainty in state 1 in a subsequent measurement. In terms of the a -wavefunctions,

$$1 = \langle 1 | 1 \rangle = \sum_{a'} |\psi_1(a')|^2, \quad (7.13)$$

which says that the square of the length of a unit vector is the sum of the absolute squares of the components.

In summary, in mathematical language, we have gone just one step beyond Euclidean geometry to a *unitary geometry*. The space, the totality of vectors is not Euclidean, but what is called *Hilbert space*. In physical terms, we are describing a geometry of physical measurements—a geometry of states, and a state space.

7.1 Spin-1/2 example

We return to spin 1/2. We want to compute $\psi_{\sigma_z''}(\sigma_z')$, which is the component, with respect to spin measurements along the z direction, of a state with a certain spin value (σ_z'') along the z' direction. The measurement symbols for these states are

$$|\sigma_{z'}'' = +1\rangle = \frac{1 + \sigma_{z'}}{2}, \quad |\sigma_{z'}'' = -1\rangle = \frac{1 - \sigma_{z'}}{2}. \quad (7.14)$$

We describe the system in terms of σ_z :

$$\sigma_{z'} = U^{-1} \sigma_z U, \quad (7.15)$$

by means of the unitary transformation

$$U = e^{i\frac{\theta}{2}\sigma_y} e^{i\frac{\phi}{2}\sigma_z}, \quad U^{-1} = e^{-i\frac{\phi}{2}\sigma_z} e^{-i\frac{\theta}{2}\sigma_y}. \quad (7.16)$$

Thus

$$\frac{1 \pm \sigma_{z'}}{2} = U^{-1} \frac{1 \pm \sigma_z}{2} U, \quad (7.17)$$

or

$$|\sigma_{z'}'' = \pm 1\rangle = U^{-1} |\sigma_z'' = \pm 1\rangle U. \quad (7.18)$$

Since the measurement symbols may be factored,

$$|\sigma_{z'}'' = \pm 1\rangle = |\sigma_{z'}''\rangle \langle \sigma_{z'}''|, \quad |\sigma_z'' = \pm 1\rangle = |\sigma_z''\rangle \langle \sigma_z''|, \quad (7.19)$$

we can factor statement (7.18) into

$$\langle \sigma_{z'}''| = \langle \sigma_z''| U, \quad |\sigma_{z'}''\rangle = U^{-1} |\sigma_z''\rangle. \quad (7.20)$$

Note that $\sigma_{z'}''$ and σ_z'' represent the same outcome, the same number, but referring to measurements along different axes.

- All information about the direction of the axis of measurement is contained in U , relative to z , the standard direction.
- All information about the outcome of the measurement is contained in $|\sigma_z''\rangle$.

To work out the wavefunctions, the components of $|\sigma_{z'}''\rangle$ with respect to $|\sigma_z'\rangle$,

$$\psi_{\sigma_{z'}''}(\sigma_z') = \langle \sigma_z' | \sigma_{z'}'' \rangle = \langle \sigma_z' | U^{-1} | \sigma_z'' \rangle, \quad (7.21)$$

we have to know the action of σ_x , σ_y , σ_z on $\langle \sigma_z'|$, $|\sigma_z'\rangle$. We recall

$$\sigma_x = | - + \rangle \langle + - | + | + - \rangle \langle - + | = | - \rangle \langle + | + | + \rangle \langle - |, \quad (7.22a)$$

$$\sigma_y = i | - + \rangle \langle - i | + | - \rangle \langle + i | = i | - \rangle \langle + | - i | + | + \rangle \langle - i |, \quad (7.22b)$$

$$\sigma_z = | + + \rangle \langle - - | - | + - \rangle \langle - + | = | + \rangle \langle + | - | - \rangle \langle - |, \quad (7.22c)$$

$$1 = | + + \rangle \langle + + | + | - - \rangle \langle - - | = | + \rangle \langle + | + | - \rangle \langle - |. \quad (7.22d)$$

Therefore,

$$\sigma_z|+\rangle = |+\rangle, \quad \sigma_z|-\rangle = -|-\rangle, \quad (7.23a)$$

$$\sigma_x|+\rangle = |-\rangle, \quad \sigma_x|-\rangle = |+\rangle, \quad (7.23b)$$

$$\sigma_y|+\rangle = i|-\rangle, \quad \sigma_y|-\rangle = -i|+\rangle, \quad (7.23c)$$

$$1|+\rangle = |+\rangle, \quad 1|-\rangle = |-\rangle. \quad (7.23d)$$

Similarly,

$$\langle +|\sigma_z = \langle +|, \quad \langle -|\sigma_z = -\langle -|, \quad (7.24a)$$

$$\langle +|\sigma_x = \langle -|, \quad \langle -|\sigma_x = \langle +|, \quad (7.24b)$$

$$\langle +|\sigma_y = -i\langle -|, \quad \langle -|\sigma_y = i\langle +|, \quad (7.24c)$$

$$\langle +|1 = \langle +|, \quad \langle -|1 = \langle -|. \quad (7.24d)$$

Therefore,

$$\begin{aligned} |\sigma''_{z'} = +1\rangle &\equiv |+, z'\rangle = U^{-1}|+, z\rangle \\ &= e^{-i\frac{\phi}{2}\sigma_z} \left(\cos \frac{\theta}{2} - i\sigma_y \sin \frac{\theta}{2} \right) |+\rangle \\ &= e^{-i\frac{\phi}{2}\sigma_z} \left(\cos \frac{\theta}{2} |+\rangle + \sin \frac{\theta}{2} |-\rangle \right) \\ &= e^{-i\frac{\phi}{2}} \cos \frac{\theta}{2} |+\rangle + e^{i\frac{\phi}{2}} \sin \frac{\theta}{2} |-\rangle \\ &= \sum_{\sigma'_z} |\sigma'_z\rangle \langle \sigma'_z | \sigma''_{z'} = +1\rangle = \psi_{+z'}(+)|+\rangle + \psi_{+z'}(-)|-\rangle, \end{aligned} \quad (7.25)$$

where $+z'$ means $\sigma''_{z'} = +1$, and $|+\rangle$ means $|+, z\rangle$. Thus we read off

$$\psi_{+z'}(+)=e^{-i\frac{\phi}{2}}\cos\frac{\theta}{2}, \quad \psi_{+z'}(-)=e^{i\frac{\phi}{2}}\sin\frac{\theta}{2}. \quad (7.26)$$

Similarly,

$$\begin{aligned} |-, z'\rangle &= e^{-i\frac{\phi}{2}\sigma_z} \left(\cos \frac{\theta}{2} - i\sigma_y \sin \frac{\theta}{2} \right) |-\rangle = e^{-i\frac{\phi}{2}\sigma_z} \left(\cos \frac{\theta}{2} |-\rangle - \sin \frac{\theta}{2} |+\rangle \right) \\ &= -e^{-i\frac{\phi}{2}} \sin \frac{\theta}{2} |+\rangle + e^{i\frac{\phi}{2}} \cos \frac{\theta}{2} |-\rangle, \end{aligned} \quad (7.27)$$

from which we read off

$$\psi_{-z'}(+)= -e^{-i\frac{\phi}{2}}\sin\frac{\theta}{2}, \quad \psi_{-z'}(-)=e^{i\frac{\phi}{2}}\cos\frac{\theta}{2}. \quad (7.28)$$

To get further exercise, work out the left vectors,

$$\begin{aligned} \langle +, z'| &= \langle +|U = \langle +| \left(\cos \frac{\theta}{2} + i\sigma_y \sin \frac{\theta}{2} \right) e^{i\frac{\phi}{2}\sigma_z} \\ &= \left(\cos \frac{\theta}{2} \langle +| + \sin \frac{\theta}{2} \langle -| \right) e^{i\frac{\phi}{2}\sigma_z} \\ &= \langle +| \cos \frac{\theta}{2} e^{i\frac{\phi}{2}} + \langle -| \sin \frac{\theta}{2} e^{-i\frac{\phi}{2}}, \end{aligned} \quad (7.29)$$

which implies that

$$\psi_{+z'}(+)^* = \cos \frac{\theta}{2} e^{i\frac{\phi}{2}}, \quad \psi_{+z'}(-)^* = \sin \frac{\theta}{2} e^{-i\frac{\phi}{2}}. \quad (7.30)$$

which are indeed true, and also note

$$|\psi_{+z'}(+)|^2 + |\psi_{+z'}(-)|^2 = 1. \quad (7.31)$$

Similarly,

$$\begin{aligned} \langle -, z' | &= \langle - | U = \langle - | \left(\cos \frac{\theta}{2} + i \sigma_y \sin \frac{\theta}{2} \right) e^{i\frac{\phi}{2} \sigma_z} \\ &= \left(\cos \frac{\theta}{2} \langle - | - \sin \frac{\theta}{2} \langle + | \right) e^{i\frac{\phi}{2} \sigma_z} \\ &= \langle + | \left(-\sin \frac{\theta}{2} e^{i\frac{\phi}{2}} \right) + \langle - | \cos \frac{\theta}{2} e^{-i\frac{\phi}{2}}, \end{aligned} \quad (7.32)$$

which correctly implies

$$\psi_{-z'}(+)^* = -\sin \frac{\theta}{2} e^{i\frac{\phi}{2}}, \quad \psi_{-z'}(-)^* = \cos \frac{\theta}{2} e^{-i\frac{\phi}{2}}. \quad (7.33)$$

Finally, check that

$$\begin{aligned} 0 &= \langle +z' | -z' \rangle = \psi_{+z'}(+)^* \psi_{-z'}(+) + \psi_{+z'}(-)^* \psi_{-z'}(-) \\ &= \left(e^{i\frac{\phi}{2}} \cos \frac{\theta}{2} \right) \left(-e^{-i\frac{\phi}{2}} \sin \frac{\theta}{2} \right) + \left(e^{-i\frac{\phi}{2}} \sin \frac{\theta}{2} \right) \left(e^{i\frac{\phi}{2}} \cos \frac{\theta}{2} \right) \\ &= -\cos \frac{\theta}{2} \sin \frac{\theta}{2} + \cos \frac{\theta}{2} \sin \frac{\theta}{2} = 0. \end{aligned} \quad (7.34)$$

The \pm signs, and the phases, which don't show up in the probabilities

$$|\psi_{+z'}(+)|^2 = \cos^2 \frac{\theta}{2}, \quad |\psi_{-z'}(+)|^2 = \sin^2 \frac{\theta}{2}, \quad (7.35)$$

etc., are crucial for the above cancellation.

We can write the above wavefunction as column vectors,

$$\psi_{+z'} = \begin{pmatrix} e^{-i\frac{\phi}{2}} \cos \frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \sin \frac{\theta}{2} \end{pmatrix}, \quad \psi_{-z'} = \begin{pmatrix} -e^{-i\frac{\phi}{2}} \sin \frac{\theta}{2} \\ e^{i\frac{\phi}{2}} \cos \frac{\theta}{2} \end{pmatrix}, \quad (7.36)$$

where the first row refers to the $\sigma_z = +1$ element, and the second row to $\sigma_z = -1$. These are nearly the same as the wavefunctions found earlier, in Eq. (4.35); they are equally as good as those (see homework).

Let's represent not just the states (vectors) but the algebraic symbols by arrays. In Eqs. (7.22a)–(7.22d), we express the symbols σ , 1, in terms of the coefficients of the four measurement symbols $|\sigma'_z\rangle\langle z''|$:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad (7.37a)$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad 1 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad (7.37b)$$

where the rows represent the values of σ'_z , the columns the values of σ''_z . This is an example of a more general procedure, which we will describe in the next chapter.