PHYS 6433: Quantum Field Theory

Problem Set 3 – Due February 08, 2017

Problem (1): Finite Rotation for a Dirac Spinor

Let us consider a finite rotation of angle ϕ about the z-axis with $x' = \Lambda x$ and

 $\omega_{12} = \epsilon_{123}\phi_3 = \phi_3 = \phi$ where $\omega_{ij} = \epsilon_{ijk}\phi_k$

Then the Dirac spinor will change under this rotation as $\psi'(x') = S_R(\Lambda)\psi(x)$.

- (a) Find the rotation operator Λ as a 4×4 matrix.
- (b) Find σ_{ij} for an infinitesimal rotation

$$S_R(\Lambda) = I - \frac{i}{4} \sigma_{ij} \omega^{ij}$$
.

(c) Show that for a finite rotation of angle ϕ about the z-axis,

$$S_R(\Lambda) = e^{-\frac{i}{2}\Sigma_3\phi_3} = I\cos\left(\frac{\phi}{2}\right) - i\Sigma_3\sin\left(\frac{\phi}{2}\right) \quad \text{where} \quad \Sigma_3 = \sigma_{12} = \begin{pmatrix} \sigma_3 & 0\\ 0 & \sigma_3 \end{pmatrix}$$

(d) Find the rotation operator $S_R(\Lambda)$ as a 4×4 matrix for a Dirac spinor.

Problem (2): Maxwell equations

In classical electrodynamics with the Heaviside-Lorentz conventions and c = 1, the 4-vector potential is defined as

$$A^{\mu} = (\Phi, A);$$

and the 4-vector current density is defined to be

$$J^{\mu} \equiv (\rho, J)$$
.

Furthermore, let us defined a second-rank antisymmetric field-strength tensor as

$$F^{\mu\nu} \equiv \partial^{\mu}A^{\nu} - \partial^{\nu}A^{\mu},$$

then we can express the first two Maxwell equations with sources as

$$\partial_{\mu}F^{\mu\nu} = \partial_{\mu}\partial^{\mu}A^{\nu} = J^{\nu},$$

with the Lorenz condition $\partial_{\mu}A^{\mu} = 0$.

Show that

$$\Box A^{\nu} = J^{\nu},$$

that is

$$\partial_{\mu}F^{\mu 0} = \partial_{\mu}\partial^{\mu}A^{0} = J^{0} = \rho, \text{ and}$$
$$\partial_{\mu}F^{\mu i} = \partial_{\mu}\partial^{\mu}A^{i} = J^{i}, i = 1, 2, 3.$$

Problem (3): Lagrangian for Vector Particles

The Lagrangian density of electrodynamics is

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - J^{\mu}A_{\mu}$$

for the photon field A^{μ} . Find the equation of motion for A^{μ} .

Problem (4): Polarization Vectors for a Massive Vector Particle

For a vector particle of mass M, energy E, and momentum \vec{k} along the z axis, the polarization vectors can be expressed as

$$\epsilon^{\mu}(\lambda = \pm 1) = \mp \frac{1}{\sqrt{2}}(0, 1, \pm i, 0), \text{ and}$$

 $\epsilon^{\mu}(\lambda = 0) = (|\vec{k}|, 0, 0, E)/M$

that are called circular polarization vectors.

- Determine how $\vec{\epsilon}(\lambda = 1)$ and $\vec{\epsilon}(\lambda = -1)$ transform under a rotation ϕ about the z axis.
- Show that every polarization vector satisfies the following relations

$$\epsilon_{\mu}k^{\mu} = 0$$
 and
 $\epsilon^{*}_{\mu}\epsilon^{\mu} = -1$.

• Show that the completeness relation is

$$\sum_{\lambda} \epsilon^*_{\mu}(\lambda) \epsilon_{\nu}(\lambda) = -g_{\mu\nu} + \frac{k_{\mu}k_{\nu}}{M^2} \,.$$

N.B. There is a minus sign in the definition of $\epsilon(\lambda = \pm 1)$. This is a standard phase convention used in the construction of the spherical harmonics $Y_{\ell m}$ for $\ell = 1$ and $m = \pm 1$.