# PHYS 3803: Quantum Mechanics I, Spring 2021 Lecture 9, February 25, 2021 (Thursday)

- Reading: Mathematical Tools (Chapter 3 in Griffiths)
- Assignment: Problem Set 4 due February 26 (Friday). Submit your homework assignments to Canvas.

Topics for Today: Mathematical Tools
2.9 Hilbert Space and Fourier Transform
2.10 Introduction to the Uncertainty Principle
Topics for Next Lecture: Mathematical Tools
2.11 Bonus: The Momentum Space or the *p*-basis
3.1 Postulates of Quantum Mechanics

### 2.9 Hilbert Space and Fourier Transform

The Hilbert space is a linear vector space where every vector can be normalized either to unity or the Dirac delta function.

Any state vector  $|f\rangle$  can be expanded in terms of a complete set of orthonormal basis.

• In the coordinate space with basis vectors  $|x\rangle$ ,

$$|f\rangle = \left(\int |x\rangle \langle x|dx\right) |f\rangle = \int \langle x|f\rangle |x\rangle dx = \int f(x)|x\rangle dx,$$

where  $f(x) \equiv \langle x | f \rangle$  is the wave function in the x-basis.

• In the momentum space with k-basis,

$$|f\rangle = \left(\int |k\rangle \langle k|dk\right) |f\rangle = \int \langle k|f\rangle |k\rangle \, dk = \int g(k)|k\rangle \, dk \,,$$

where  $g(k) \equiv \langle k | f \rangle$  is the wave function in the k-basis.

Applying completeness relations

$$\int |x\rangle \langle x| \, dx = \mathbf{I}$$

we can find the relations between wave functions g(k) and f(x)

$$g(k) \equiv \langle k|f \rangle$$
  
=  $\int \langle k|x \rangle \langle x|f \rangle dx$   
=  $\int \psi_k^*(x) f(x) dx$   
=  $\frac{1}{\sqrt{2\pi}} \int e^{-ikx} f(x) dx$ .

where  $g(k) = \langle k | f \rangle$  and  $f(x) = \langle x | f \rangle$ .

This is just the Fourier transform.

Similarly we can apply the completeness relations

$$\int |k\rangle \langle k| \, dk = \mathbf{I}$$

and show that

$$f(x) \equiv \langle x|f \rangle$$
  
=  $\int \langle x|k \rangle \langle k|f \rangle dk$   
=  $\int \psi_k(x)g(k) dk$   
=  $\frac{1}{\sqrt{2\pi}} \int e^{ikx}g(k) dk$ 

where  $f(x) = \langle x | f \rangle$  and  $g(k) = \langle k | f \rangle$ .

This is the inverse Fourier transforms.

Fourier transforms take the state vector from one basis to another.

In the eigenbasis of K (the k-basis), the eigenvalue equation is

 $K|k\rangle = k|k\rangle$ .

The normalized eigenvectors form a complete set of orthonormal basis

$$\langle k|q \rangle = \delta(k-q)$$
 (orthonormal relation)  
 $\int |k\rangle \langle k| dk = I$  (completeness relation).

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In the k-basis,

- a state vector  $|\psi\rangle$  is represented by a wave function  $\phi(k) \equiv \langle k | \psi \rangle$ ,
- an operator is represented with matrix elements

$$\Omega_{kq} = \langle k | \Omega | q \rangle \,.$$

For example,

$$\langle k|K|q 
angle = q \langle k|q 
angle = q \delta(k-q)$$
  
=  $k \langle k|q 
angle = k \delta(k-q)$ .

In the eigenbasis of X (the x-basis), the eigenvalue equation is

 $X|x\rangle = x|x\rangle.$ 

The normalized eigenvectors form a complete set of orthonormal basis

$$\langle x|y \rangle = \delta(x-y)$$
 (orthonormal relation)  
 $\int |x \rangle \langle x| dx = I$  (completeness relation).

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In the x-basis,

- a state vector  $|\psi\rangle$  is represented by a wave function  $\psi(x) \equiv \langle x | \psi \rangle$ ,
- an operator is represented with matrix elements

$$\Omega_{xy} = \langle x | \Omega | y \rangle \,.$$

For example,

$$\langle x|X|y\rangle = y\langle x|y\rangle = y\delta(x-y)$$
  
=  $x\langle x|y\rangle = x\delta(x-y)$ .

The operator X usually takes a state vector  $|f\rangle$  to a different state vector  $|h\rangle$  (transformation)

$$X|f\rangle = |h\rangle$$
.

Multiplying both sides with a dual basis vector  $\langle x |$ , we obtain

$$\begin{aligned} x|X|f\rangle &= \int \langle x|X|y\rangle \langle y|f\rangle \, dy \\ &= \int y\delta(x-y)f(y) \, dy \\ &= xf(x) = h(x) \, . \end{aligned}$$

Thus the effect of X on a vector  $|f\rangle$  is to multiply it by x.

The matrix element of the operator X in the  $|k\rangle$  basis is

$$\begin{split} k|X|q\rangle &= \int dx \int dy \, \langle k|x \rangle \langle x|X|y \rangle \langle y|q \rangle \\ &= \int dx \int dy \, \psi_k^*(x) y \delta(x-y) \psi_q(y) \\ &= \frac{1}{2\pi} \int x e^{-i(k-q)x} \, dx \\ &= \frac{1}{2\pi} i \frac{d}{dk} \int e^{-i(k-q)x} \, dx \\ &= i \frac{d}{dk} \delta(k-q) \, . \end{split}$$

- In the x-basis, X acts as x and K acts as -id/dx on functions.
- In the k-basis, K acts as k and X acts as id/dk on functions.
- Operators with such reciprocity are called conjugate operators.

Clearly, conjugate operators do not commute.

In the x-basis, we have

$$\begin{array}{lll} X|f(x)\rangle & \to & xf(x) \quad \text{and} \quad K|f(x)\rangle \to -i\frac{df(x)}{dx} \\ XK|f(x)\rangle & \to & -ix\frac{df(x)}{dx} \\ KX|f(x)\rangle & \to & -i\frac{d}{dx}\left[xf(x)\right] \,. \end{array}$$

Thus

$$(XK - KX)|f(x)\rangle \to -ix\frac{df(x)}{dx} + if(x) + ix\frac{df}{dx} = if(x) \to i|f\rangle.$$

Thus

$$[X,K]|f\rangle = i|f\rangle$$
 or  $[X,K] = iI$ ,

and

$$[X, P]|f\rangle = i\hbar|f\rangle$$
 or  $[X, P] = i\hbar I$ .

X is the position operator and  $P = \hbar K$  is the momentum operator.

### 2.10 Introduction to the Uncertainty Principle

Let A and B be two non-commuting operators with

$$[A,B] = i\hbar$$

As we have seen before, these are conjugate operators. Let  $\Delta A$  be the root mean square deviation of the operator A. Thus

$$(\Delta A)^2 \equiv \langle A^2 \rangle - \langle A \rangle^2$$

Similarly let  $\Delta B$  be the root mean square deviation of the operator B. Thus

$$(\Delta B)^2 = \langle B^2 \rangle - \langle B \rangle^2$$

then

$$\Delta A \Delta B \ge \frac{\hbar}{2}$$
 (Heisenberg's uncertainty principle)

First of all notice that

$$(\Delta \Omega)^{2} = \langle \Omega^{2} \rangle - \langle \Omega \rangle^{2}$$
  
=  $\langle \Omega^{2} - 2\Omega \langle \Omega \rangle + \langle \Omega \rangle^{2} \rangle$   
=  $\langle (\Omega - \langle \Omega \rangle)^{2} \rangle$   
 $\Delta \Omega \equiv \langle (\Omega - \langle \Omega \rangle)^{2} \rangle^{1/2} = [\Omega^{2} \rangle - \langle \Omega \rangle^{2}]^{1/2}$ 

where (i)  $\Delta\Omega$  is the standard deviation which measures the average fluctuation around the mean.

N.B. (i) ΔΩ is often called the root mean squared deviation or the uncertainty in Ω.
(ii) (ΔΩ)<sup>2</sup> is called the mean square deviation or the variance.

## **Expectation Value in the x-Basis**

In the x-basis with  $|x\rangle$ , the state vector can be expanded as

$$|\psi\rangle = \left(\int |x\rangle\langle x\right) |\psi\rangle \, dx = \int \psi(x) |x\rangle dx \,, \quad \text{where} \quad \psi(x) \equiv \langle x|\psi\rangle \,.$$

And the expectation value of  $\Omega$  becomes

$$\begin{split} \langle \psi | \Omega | \psi \rangle &= \langle \psi | \left( \int |x\rangle \langle x| \, dx \right) (\Omega) \left( \int |y\rangle | \langle y| \, dy \right) | \psi \rangle \\ &= \int dx \int dy \, \langle \psi | x \rangle \langle x| \Omega | y \rangle \langle y| \psi \rangle \,. \end{split}$$

For example, the matrix elements of X and P are

$$\langle x|X|y\rangle = y\delta(x-y) = x\delta(x-y)$$
, and  
 $\langle x|P|y\rangle = -i\hbar \frac{d}{dx}\delta(x-y)$ .

**Examples:** (a)  $\Omega = X$  and (b)  $\Omega = P$ In the coordinate space we have a complete set of orthonormal basis  $\langle x|y \rangle = \delta(x-y)$  (orthonormal),  $\int |x \rangle \langle x| \, dx = I$  (completeness). The state vector  $|\psi \rangle$  and operators ( $\Omega$ ) are represented with  $\psi(x) \equiv \langle x|\psi \rangle$  (wave function),  $\Omega_{xy} \equiv \langle x|\Omega|y \rangle$  (matrix element). (a) The expectation value of X becomes

$$egin{aligned} &\langle\psi|\left(\int|x
angle\langle x|\,dx
ight)(X)\left(\int|y
angle|\langle y|\,dy
ight)|\psi
ight)\ &=&\int dx\int dy\,\langle\psi|x
angle\langle x|X|y
angle\langle y|\psi
angle\ &=&\int dx\int dy\,\psi^*(x)\left[y\delta(x-y)
ight]\psi(y)\ &=&\int\psi^*(x)\left[x\psi(x)
ight]\,dx\,. \end{aligned}$$

(b) The expectation value of P becomes

$$\begin{split} \langle \psi | P | \psi \rangle &= \langle \psi | \left( \int |x\rangle \langle x| \, dx \right) (P) \left( \int |y\rangle | \langle y| \, dy \right) | \psi \rangle \\ &= \int dx \int dy \, \langle \psi | x \rangle \langle x | P | y \rangle \langle y | \psi \rangle \\ &= \int dx \int dy \, \psi^*(x) \left[ -i\hbar \frac{d}{dx} \delta(x-y) \right] \psi(y) \\ &= -i\hbar \int \psi^*(x) \left[ \frac{d}{dx} \psi(x) \right] dx \,, \end{split}$$

where

$$\langle x|P|y\rangle = -i\hbar \frac{d}{dx}\delta(x-y)$$
 where  $P = \hbar K = -i\hbar D$ 

in the x-basis.

Useful formulas:

(a) 
$$\int_0^\infty \frac{1}{(x^2 + \alpha^2)^2} dx = \left(\frac{1}{2\alpha^2}\right) \left[\frac{x}{x^2 + \alpha^2} + \frac{1}{\alpha} \tan^{-1}\left(\frac{x}{\alpha}\right)\right]_0^\infty$$
  
(b) 
$$\int_0^\infty \frac{x^{m-1}}{(1+x)^{m+n}} dx = \frac{\Gamma(m)\Gamma(n)}{\Gamma(m+n)}.$$

(c) To apply Gaussian Integrals, you might need to complete the square for a quadratic function

$$ax^{2} + bx + c = a\left(x^{2} + \frac{b}{a}x\right) + c$$

$$= a\left[x^{2} + \frac{b}{a}x + \left(\frac{b}{2a}\right)^{2}\right] - a \cdot \left(\frac{b}{2a}\right)^{2} + c$$

$$= a\left[x + \left(\frac{b}{2a}\right)\right]^{2} - \frac{b^{2}}{4a} + c$$

$$= a\left[x + \left(\frac{b}{2a}\right)\right]^{2} + c - \frac{b^{2}}{4a}.$$

#### **A Brief Dictionary for Notations**

Here are some helpful relations between two sets of notations when you read the textbook:

- State vector:  $|\Psi(t)\rangle \rightarrow |S(t)\rangle$
- Operators:  $\Omega \to \hat{Q}$
- Basis vectors:  $|e_1\rangle \rightarrow |1\rangle$  and  $|e_2\rangle \rightarrow |2\rangle$
- Orthonormal relations:  $\langle x|y \rangle = \delta(x-y) \rightarrow \langle g_x|g_y \rangle = \delta(x-y)$  and  $\langle p|q \rangle = \delta(p-q) \rightarrow \langle f_p|f_q \rangle = \delta(p-q).$
- Special wave functions:  $\psi_k(x) = \langle x|k \rangle = \frac{1}{\sqrt{2\pi}} e^{ikx} \to f_k(x)$  and  $\psi_p(x) = \langle x|p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{(i/\hbar)px} \to f_p(x).$
- Uncertainty or root mean square deviation  $\Delta \Omega \rightarrow \sigma_{\Omega}$ .