

Distributional interpretation of Casimir energy

Linearized Einstein eq. (related to a particular path)

$$-\frac{d^2 h}{dx^2} = \frac{8}{\pi} \frac{4x^2 - 3t^2}{(t^2 + 4x^2)^3} \theta(x) \quad (1) \quad \left(h \text{ is the } 00\text{-component of the metric tensor} \right)$$

integrating twice and imposing that h as well as its derivative vanish at $x=0$ [note the half-space $x < 0$ is assumed to be empty]

$$\Rightarrow h(x) = \frac{\theta(x)}{\pi} \left[\frac{4x}{t^3} \tan^{-1}\left(\frac{2x}{t}\right) - \frac{1}{t^2 + 4x^2} + \frac{1}{t^2} \right] \quad (2)$$

As $t \rightarrow 0$ which corresponds to removing the regularization parameter this expression is ill-defined in the usual sense of h being a function. Therefore we consider a distributional interpretation of (2) and we check if eq. (1) remains valid as $t \rightarrow 0$ (in distributional sense).

The basis for the analysis to come is the following theorem.

Moment expansion theorem: Let $f \in S'(\mathbb{R})$ with support bounded on the left. Suppose

$f(x) = b_1 x^{\beta_1} + \dots + b_n x^{\beta_n} + o(x^\beta)$ as $x \rightarrow \infty$, where $\beta_1 > \dots > \beta_n > \beta$, and $-(k+1) > \beta > -(k+2)$. Then as $\lambda \rightarrow \infty$

$$f(\lambda x) = \sum_{j=1}^n b_j \frac{g_j(\lambda x)}{\lambda^{\beta_j}} + \sum_{j=0}^k (-1)^j \frac{\mu_j(f) \delta^{(j)}(\lambda x)}{j!} + o(\lambda^\beta)$$

in the space $S'(\mathbb{R})$, where $g_j(x) = P_f(x^{\beta_j} \theta(x))$ if $\beta_j = -1, -2, -3, \dots$

Here the moments are

$$\mu_j(f) = \int_{-\infty}^{\infty} f(x) x^j dx$$

$P_f =$ Pseudo-fct.
 $F.P. =$ finite part

Yasser working on more general results; leave the proof to him.

Explanation of notation:

Let $F(\epsilon) = \int_{\epsilon}^{\infty} f(x) dx$ be given as $F(\epsilon) = F_1(\epsilon) + F_2(\epsilon)$, where $F_1(\epsilon)$ is simple as $\epsilon \rightarrow 0$ and $\lim_{\epsilon \rightarrow 0} \frac{F_2(\epsilon)}{\epsilon} = A$ exists. We then define

$$\text{F.P. } F(\epsilon) = A.$$

[Similarly for $F(M) = \int_0^M f(x) dx$ as $M \rightarrow \infty$.]

Definition of Pseudofunctions:

$$\langle \text{Pf}(g(x)), \phi(x) \rangle = \text{F.P.} \int_0^{\infty} g(x) \phi(x) dx.$$

Applications of the Moment expansion theorem to $h(x)$, eq. (2) ($\lambda = \frac{1}{t}$):

Write $h(x) = h_1(x) + h_2(x) + h_3(x)$ with

$$h_1(x) = \frac{\theta(x)}{\pi} \frac{4x}{t^3} \tan^{-1}\left(\frac{2x}{t}\right) = \frac{4}{\pi} \lambda^2 \theta(x) (2x) \tan^{-1}(2\lambda x)$$

$$h_2(x) = -\frac{\theta(x)}{\pi} \frac{1}{t^2 + 4x^2} = -\frac{\lambda^2}{\pi} \theta(x) \frac{1}{1 + 4\lambda^2 x^2}$$

$$h_3(x) = \frac{\theta(x)}{\pi} \cdot \frac{1}{t^2} = \lambda^2 \cdot \frac{\theta(x)}{\pi}.$$

↑ done

Asymptotic behavior of h_3 : done.

Asymptotic behavior of h_2 :

$$f_2(x) = \frac{\theta(x)}{1+4x^2} = \frac{1}{4x^2} + o\left(\frac{1}{x^4}\right)$$

numbers, easily evaluated

$$-3 < \beta < -4$$

so and

$$h_2(x) = -\frac{\lambda^2}{\pi} \left\{ \frac{1}{4} Pf \left[\frac{\theta(\lambda x)}{(\lambda x)^2} \right] + \sum_{j=0}^2 \frac{(-1)^j \mu_j(f_2) \delta^{(j)}(\lambda x)}{j!} + o(\lambda^\beta) \right\}$$

Asymptotic behavior of h_1 :

$$f_1(x) = \frac{\theta(x)}{x} \tan^{-1}(\lambda x) = \frac{\pi}{2} \cdot x - \frac{1}{2} + \frac{1}{24} \cdot \frac{1}{x^2} + o\left(\frac{1}{x^4}\right), \text{ so } -3 < \beta < -4$$

and

$$h_1(x) = \frac{4}{\pi} \lambda^2 \left\{ \frac{\pi}{2} \theta(\lambda x) (\lambda x) - \frac{1}{2} \theta(\lambda x) (\lambda x)^0 + \frac{1}{24} Pf \left[\frac{\theta(\lambda x)}{(\lambda x)^2} \right] \right.$$

$$\left. + \sum_{j=0}^2 \frac{(-1)^j \mu_j(f_1) \delta^{(j)}(\lambda x)}{j!} + o(\lambda^\beta) \right\}$$

numbers, easily evaluated

In order to identify the $\lambda \rightarrow \infty$ behavior more clearly we need to understand scaling properties of distributions in some detail.

Def.: The distribution $g(Ax)$ is defined as

$$\langle g(Ax), \phi(x) \rangle = \frac{1}{|\det A|} \langle f(x), \phi(A^{-1}x) \rangle$$

Pedantic point of view | terms in h_1

Ex.: $\langle \theta(\lambda x)(\lambda x), \phi(x) \rangle = \frac{1}{\lambda} \langle \theta(x) \cdot x, \phi\left(\frac{x}{\lambda}\right) \rangle = \frac{1}{\lambda} \int x \phi\left(\frac{x}{\lambda}\right) dx$

sub.: $y = \frac{x}{\lambda}$
 $= \frac{1}{\lambda} \int_0^\infty \lambda y \phi(y) \lambda dy = \lambda \langle \theta(x) \cdot x, \phi(x) \rangle$

similarly for the second term

Ex.: $\langle \delta^{(j)}(\lambda x), \phi(x) \rangle = \frac{1}{\lambda} \langle \delta^{(j)}(x), \phi\left(\frac{x}{\lambda}\right) \rangle$

$$= \frac{(-1)^j}{\lambda} \langle \delta(x), \frac{d^j}{dx^j} \phi\left(\frac{x}{\lambda}\right) \rangle = \frac{(-1)^j}{\lambda} \left[\frac{d^j}{dx^j} \phi\left(\frac{x}{\lambda}\right) \right]_{x=0}$$

$$= \frac{(-1)^j}{\lambda^{j+1}} \phi^{(j)}(0) = \frac{(-1)^j}{\lambda^{j+1}} \langle \delta(x), \phi^{(j)}(x) \rangle = \frac{1}{\lambda^{j+1}} \langle \delta^{(j)}(x), \phi(x) \rangle$$

Ex. pseudo-fct.: Before we consider the scaling we derive some properties of the pseudo-fct.

Lemma: $\langle Pf\left(\frac{\theta(x)}{x}\right), \phi(x) \rangle = \int_0^1 \frac{\phi(x) - \phi(0)}{x} dx + \int_1^\infty \frac{\phi(x)}{x} dx$

Proof: Note, that [by def. of Pf this is what we have to consider]

$$\int_\epsilon^\infty \frac{\phi(x)}{x} dx = \int_\epsilon^1 \frac{\phi(x)}{x} dx + \int_1^\infty \frac{\phi(x)}{x} dx = \int_\epsilon^1 \frac{\phi(x) - \phi(0)}{x} dx + \int_\epsilon^1 \frac{\phi(0)}{x} dx + \int_1^\infty \frac{\phi(x)}{x} dx$$

$$= -\ln \epsilon \phi(0) + \int_\epsilon^1 \frac{\phi(x) - \phi(0)}{x} dx + \int_1^\infty \frac{\phi(x)}{x} dx$$

↳ $x \rightarrow 0$ behavior improved

This implies the anetide.

Lemma: $Pf\left[\frac{\theta(\lambda x)}{\lambda x}\right] = \frac{1}{\lambda} Pf\left[\frac{\theta(x)}{x}\right] + \frac{\ln \lambda \cdot \delta(x)}{\lambda}$

Its basically about understanding this Rev.

Proof: By def.

$$\langle Pf\left[\frac{\theta(\lambda x)}{\lambda x}\right], \phi(x) \rangle = \left(\frac{1}{\lambda}\right) \langle Pf\left[\frac{\theta(x)}{x}\right], \phi\left(\frac{x}{\lambda}\right) \rangle = \frac{1}{\lambda} \int_0^\infty \frac{\phi\left(\frac{x}{\lambda}\right)}{x} dx$$

Clearly, [this is what needs to be considered]

$$\int_{\varepsilon}^{\infty} \frac{\phi\left(\frac{x}{\lambda}\right)}{x} dx = \left(\int_{\varepsilon}^{\lambda} + \int_{\lambda}^{\infty} \right) \frac{\phi\left(\frac{x}{\lambda}\right)}{x} dx$$

reasonable splitting
to scaling to come

Furthermore,

$$\int_{\lambda}^{\infty} \frac{\phi\left(\frac{x}{\lambda}\right)}{x} dx = \int_1^{\infty} \frac{\phi(y)}{y} dy \quad \left[\text{fits the second term in previous lemma; will be part of first term in this lemma} \right]$$

$$\int_{\varepsilon}^{\lambda} \frac{\phi\left(\frac{x}{\lambda}\right)}{x} dx = \int_{\varepsilon/\lambda}^1 \frac{\phi(y)}{y} dx = \int_{\varepsilon/\lambda}^1 \frac{\phi(y) - \phi(0) + \phi(0)}{y} dy$$

$$= -\ln(\varepsilon/\lambda) \phi(0) + \int_{\varepsilon/\lambda}^1 \frac{\phi(y) - \phi(0)}{y} dy.$$

Using the definition and the previous lemma implies the assertion. •

Similarly to $\text{Pf} \left[\frac{\theta(x)}{x} \right]$; proof eg by induction or by direct calculation.

Adding up:
$$h(x) = 2\lambda^3 \theta(x) \cdot x - \frac{\lambda^2}{\pi} \theta(x) - \frac{1}{12\pi} \text{Pf} \left[\frac{\theta(x)}{x^2} \right] \quad (3)$$

$$- \frac{1}{18\pi} \delta'(x) + \frac{1}{12\pi} \ln(2\lambda) \delta'(x) + \dots$$

leave on board

plus terms that vanish as $\lambda \rightarrow \infty$

Check eq. (1):

$$-\frac{d^2 h}{dx^2} = \frac{8}{\pi} \frac{4x^2 - 3t^2}{(t^2 + tx^2)^3} \theta(x) = \frac{8}{\pi} \frac{t^2}{t^6} \frac{4(\frac{x}{t})^2 - 3}{(1 + 4(\frac{x}{t})^2)^3} \theta(x)$$

Rewrite in form suitable for $\lambda \rightarrow \infty$

$$= \frac{8}{\pi} \lambda^4 \frac{4\lambda^2 x^2 - 3}{(1 + 4\lambda^2 x^2)^3} \theta(x)$$

exactly as before; as $x \rightarrow \infty$ this behaves like $1/x^4$ and the following is apparent

I do not want to be as explicit as before, but the relevant $f(x)$ is

$$[f(x) = \theta(x) \frac{4x^2 - 3}{(1 + 4x^2)^3} \sim \frac{1}{16x^4} + \mathcal{O}\left(\frac{1}{x^6}\right); -5 > \beta > -6]$$

$j=4$ neglected as vanishes as $\lambda \rightarrow \infty$

$$= \frac{8}{\pi} \lambda^4 \left\{ \frac{1}{16} Pf\left(\frac{\theta(\lambda x)}{(\lambda x)^4}\right) + \sum_{j=0}^3 (-1)^j \mu_j(f) \frac{\delta^{(j)}(\lambda x)}{j!} + \mathcal{O}\left(\frac{1}{\lambda^5}\right) \right\}$$

apply derived scaling behavior

$$= \frac{1}{2\pi} Pf\left(\frac{\theta(x)}{x^4}\right) - 2\lambda^3 \delta(x) + \frac{1}{\pi} \lambda^2 \delta'(x) + \frac{1}{8\pi} \delta'''(x)$$

$$- \frac{1}{12\pi} \ln(2\lambda) \delta'''(x) + \mathcal{O}\left(\frac{1}{\lambda}\right)$$

This is the rhs! Evaluate lhs.

Evaluate $\frac{d^2 h}{dx^2}$: (go back to eq. (3) \rightarrow 2nd, 4th and 5th terms are trivial; 2nd and 5th terms already match; 4th not yet and more contributions need to generate δ''' terms)

Let's see how to evaluate derivatives of distributions!

$$\text{Ex.: } \left\langle \frac{d^2}{dx^2} \theta(x) \cdot x, \phi(x) \right\rangle = \int_0^{\infty} x \frac{d^2 \phi}{dx^2} dx = x \cdot \phi'(x) \Big|_0^{\infty} - \int_0^{\infty} \frac{d\phi}{dx} dx$$

(partial integration)

$$= \phi'(0) = \langle \delta(x), \phi(x) \rangle \quad (\text{first term in (3) matches 2nd above.})$$

Ex.: $\frac{d}{dx} \text{Pf} \left(\frac{\theta(x)}{x^k} \right) = -k \text{Pf} \left(\frac{\theta(x)}{x^{k+1}} \right) + \frac{(-1)^k \delta^{(k)}(x)}{k!}$

↑ expected
↑ needs explanation

Proof: Note, that

$$\left\langle \frac{d}{dx} \text{Pf} \left(\frac{\theta(x)}{x^k} \right), \phi(x) \right\rangle = - \left\langle \text{Pf} \left(\frac{\theta(x)}{x^k} \right), \phi'(x) \right\rangle.$$

consider: partial integration

$$\int_{\epsilon}^{\infty} \frac{\phi'(x)}{x^k} dx = \left. \frac{\phi(x)}{x^k} \right|_{\epsilon}^{\infty} - \int_{\epsilon}^{\infty} \left(\frac{d}{dx} \frac{1}{x^k} \right) \phi(x) dx$$

$$= - \frac{\phi(\epsilon)}{\epsilon^k} - k \int_{\epsilon}^{\infty} \frac{\phi(x)}{x^{k+1}} dx$$

Taylor series exp.

$$= - \frac{1}{\epsilon^k} \sum_{j=0}^{\infty} \frac{\phi^{(j)}(0)}{j!} \epsilon^j - k \int_{\epsilon}^{\infty} \frac{\phi(x)}{x^{k+1}} dx.$$

Taking the ~~limit~~ ^{F.P.} ~~part~~ implies the assertion. ■

So: $-\frac{1}{12\pi} \frac{d^2}{dx^2} \text{Pf} \left(\frac{\theta(x)}{x^2} \right) = -\frac{1}{2\pi} \text{P.f.} \left(\frac{\theta(x)}{x^4} \right) - \frac{5}{72\pi} \delta'''(x).$

Adding up all terms in $-\frac{dh^2}{dx^2}$, eq. (1) follows as $\eta \rightarrow \infty$.

In particular, the finite part as $\eta \rightarrow \infty$ satisfies the differential equation. So eventually defining the metric as the finite part of the distributional answer makes sense!

Proof of Moment exp. theorem:

It is possible to find a constant M such that we can write $f = f_0 + f_1$, where f_0 has compact support and where $f_1(x)$ has support in $[0, \infty)$ and with $|f_1(x) - \sum_{j=1}^n b_j x^j| \leq Mx^\beta, x \geq 0$.

This is clearly possible for $x \geq 1$ from the asymptotic behavior. By suitably choosing f_0 it can be enforced also for $0 \leq x < 1$.

Then $f_0(\lambda x) = \sum_{j=0}^k \frac{(-1)^j \mu_j(f_0) \delta^{(j)}(x)}{j! \lambda^{j+1}} + o(\lambda^{-(k+2)}), \text{ as } \lambda \rightarrow \infty$

Furthermore:

$\langle f_1(\lambda x) - \sum_{j=1}^n b_j g_j(\lambda x) - \sum_{j=0}^k \frac{(-1)^j \mu_j(f_1) \delta^{(j)}(x)}{j! \lambda^{j+1}}, \phi(x) \rangle$

give me a second to define this

$= \langle f_1(\lambda x) - \sum_{j=1}^n b_j g_j(\lambda x), \phi(x) \rangle -$

well def quantity by construction

$\sum_{j=0}^k \frac{\phi^{(j)}(0)}{j! \lambda^{j+1}} \int_0^\infty (f_1(x) - \sum_{j=1}^n b_j g_j(x)) x^j dx$

can be replaced by $x^\beta \lambda^\beta$ & integral is finite.

$= \int_0^\infty (f_1(\lambda x) - \sum_{j=1}^n b_j g_j(\lambda x)) (\phi(x) - \sum_{j=0}^k \frac{\phi^{(j)}(0)}{j!} x^j) dx$

$\leq \int_0^1 Mx^\beta \lambda^\beta kx^{k+1} dx + \int_1^\infty Mx^\beta \lambda^\beta kx^k dx = Mk \left[\frac{1}{k+\beta+2} - \frac{1}{k+\beta+1} \right] \lambda^\beta$

$= o(\lambda^\beta), \text{ implying the assertion.}$