

Physics 5970
Homework Assignment I

1. Let $\beta = \tau_3 \times \sigma_2$, $\alpha_1 = 1 \times \sigma_1$, $\alpha_2 = \tau_2 \times \sigma_2$, $\alpha_3 = 1 \times \sigma_3$
 $\beta^2 = \tau_3^2 \times \sigma_2^2 = 1 \times 1 = 1$, likewise $\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = 1$

$$\{\beta, \alpha_1\} = (\tau_3 \times \sigma_2)(1 \times \sigma_1) + (1 \times \sigma_1)(\tau_3 \times \sigma_2) \\ = \tau_3 \times (\sigma_1 \sigma_2 + \sigma_2 \sigma_1) = 0$$

$$\{\beta, \alpha_2\} = (\tau_3 \times \sigma_2)(\tau_2 \times \sigma_2) + (\tau_2 \times \sigma_2)(\tau_3 \times \sigma_2) \\ = (\tau_3 \tau_2 + \tau_2 \tau_3) \times \sigma_2^2 = 0$$

$$\{\beta, \alpha_3\} = (\tau_3 \times \sigma_2)(1 \times \sigma_3) + (1 \times \sigma_3)(\tau_3 \times \sigma_2) \\ = \tau_3 \times (\sigma_2 \sigma_3 + \sigma_3 \sigma_2) = 0$$

from properties of Pauli matrices

Similarly $\{\alpha_1, \alpha_2\} = (1 \times \sigma_1)(\tau_2 \times \sigma_2) + (\tau_2 \times \sigma_2)(1 \times \sigma_1) \\ = \tau_2 \times (\sigma_1 \sigma_2 + \sigma_2 \sigma_1) = 0$

$$\{\alpha_1, \alpha_3\} = (1 \times \sigma_1)(1 \times \sigma_3) + (1 \times \sigma_3)(1 \times \sigma_1) \\ = 1 \times (\sigma_1 \sigma_3 + \sigma_3 \sigma_1) = 0$$

$$\{\alpha_2, \alpha_3\} = (\tau_2 \times \sigma_2)(1 \times \sigma_3) + (1 \times \sigma_3)(\tau_2 \times \sigma_2) \\ = \tau_2 \times (\sigma_2 \sigma_3 + \sigma_3 \sigma_2) = 0$$

This verifies that the Dirac algebra is satisfied. Obvious $\beta, \vec{\alpha}$ are Hermitian, being the direct product of Hermitian matrices. Finally

$$\beta^* = \tau_3^* \times \sigma_2^* = \tau_3 \times (-\sigma_2) = -\beta \quad \checkmark$$

$$\alpha_1^* = 1^* \times \sigma_1^* = 1 \times \sigma_1 = \alpha_1$$

$$\alpha_2^* = \tau_2^* \times \sigma_2^* = (-\tau_2) \times (-\sigma_2) = \alpha_2$$

$$\alpha_3^* = 1^* \times \sigma_3^* = 1 \times \sigma_3 = \alpha_3 \quad \checkmark$$

2. Let $\frac{1}{2} [\alpha_i, \alpha_j] = i \varepsilon_{ijk} \Sigma_k$

or since $\varepsilon_{ijk} \varepsilon_{ijl} = \delta_{jj} \delta_{kl} - \delta_{jl} \delta_{jk} = 2\delta_{kl}$

$$\begin{aligned} \Sigma_k &= \frac{1}{4i} \varepsilon_{ijk} [\alpha_i, \alpha_j] \\ &= \frac{1}{2i} \varepsilon_{ijk} \alpha_i \alpha_j \quad \text{since } \varepsilon_{ijk} = -\varepsilon_{jik} \end{aligned}$$

Then $\Sigma_k \Sigma_l = -\frac{1}{4} \varepsilon_{ijk} \varepsilon_{rsl} \alpha_i \alpha_j \alpha_r \alpha_s$

There are only three α 's so at least two of these must be the same, say $i=r$ (equivalently $j=s, i=s, j=s$; other possibilities give 0)

$$\sum_k \sum_l = -\frac{1}{4} \underbrace{\epsilon_{ijk} \epsilon_{isl}}_{\delta_{js} \delta_{kl} - \delta_{jl} \delta_{ks}} \underbrace{\alpha_i \alpha_j \alpha_i \alpha_s}_{-\alpha_i^2 \alpha_j} \times \begin{cases} 4 & \text{if } j \neq s \\ 2 & \text{if } j = s \end{cases}$$

since $i \neq j$

$$= \left(\frac{2}{3} \delta_{js} \delta_{kl} - \delta_{jl} \delta_{ks} \right) \alpha_j \alpha_s$$

$$= \frac{2}{3} \delta_{kl} \underbrace{\vec{\alpha} \cdot \vec{\alpha}}_3 - \alpha_l \alpha_k$$

$$\alpha_l \alpha_k = \frac{1}{i} \{ \alpha_l, \alpha_k \} + \frac{1}{2} [\alpha_l, \alpha_k]$$

$$= \delta_{kl} + i \epsilon_{lkm} \sum_m$$

$$\therefore \sum_k \sum_l = \delta_{kl} + i \epsilon_{klm} \sum_m \quad \checkmark$$

$$\text{Explicitly: } \Sigma_1 = \frac{1}{2i} \epsilon_{1jk} \alpha_j \alpha_k = \frac{1}{i} \alpha_2 \alpha_3$$

$$= \frac{1}{i} (\tau_1 \times \sigma_2) (1 \times \sigma_3) = \frac{1}{i} \tau_2 \times i \sigma_1$$

$$= \tau_2 \times \sigma_1$$

$$\Sigma_2 = \frac{1}{2i} \epsilon_{2jk} \alpha_j \alpha_k = \frac{1}{i} \alpha_3 \alpha_1$$

$$= \frac{1}{i} (1 \times \sigma_3) (1 \times \sigma_1) = \frac{1}{i} (1 \times \sigma_3 \sigma_1) = 1 \times \sigma_2$$

$$\text{and } \Sigma_3 = \frac{1}{i} \alpha_1 \alpha_2 = \frac{1}{i} (1 \times \sigma_1) (\tau_2 \times \sigma_2) \\ = \tau_2 \times \sigma_3$$

$$\text{check } \Sigma_1 \Sigma_2 = (\tau_2 \times \sigma_1) (1 \times \sigma_2) \\ = \tau_2 \times \sigma_1 \sigma_2 = i \tau_2 \times \sigma_3 = i \Sigma_3$$

3. Verify rotational invariance of the Dirac equation: If in the primed frame,

$$i \frac{\partial}{\partial t'} \bar{\Psi}(\vec{x}', t') = (\frac{1}{i} \vec{\alpha} \cdot \vec{\nabla}' + \beta m) \bar{\Psi}(\vec{x}', t')$$

$$\text{where } t' = t, \vec{x}' = \vec{x} - \delta \vec{\omega} \times \vec{x}$$

$$\vec{\nabla}' = \vec{\nabla} - \delta \vec{\omega} \times \vec{\nabla}$$

$$\bar{\Psi}(\vec{x}', t) = U \Psi(\vec{x}, t)$$

$$= (1 + i \delta \vec{\omega} \cdot \frac{1}{2} \vec{\Sigma}) \Psi$$

$$\text{then } (1 + i \delta \vec{\omega} \cdot \frac{1}{2} \vec{\Sigma}) i \frac{\partial}{\partial t} \Psi(\vec{x}, t)$$

$$= (1 + i \delta \vec{\omega} \cdot \frac{1}{2} \vec{\Sigma}) (\vec{\alpha} \frac{1}{i} \cdot \vec{\nabla} + \beta m) \Psi(\vec{x}, t)$$

$$- \frac{1}{i} \vec{\alpha} \cdot \delta \vec{\omega} \times \vec{\nabla} \Psi + \left[\frac{1}{i} \vec{\alpha} \cdot + i \delta \vec{\omega} \cdot \frac{1}{2} \vec{\Sigma} \right] \vec{\nabla} \Psi$$

But $[\vec{\alpha}, \delta \vec{\omega} \cdot \vec{\Sigma}]_i$

$$= [\alpha_i, \frac{1}{2i} \delta \omega_j \varepsilon_{jkl} \alpha_k \alpha_l]$$

$$= \frac{1}{2i} \delta \omega_j \varepsilon_{jkl} \left(\underbrace{\{\alpha_i, \alpha_k\}}_{2\delta_{ik}} \alpha_l - \alpha_k \underbrace{\{\alpha_i, \alpha_l\}}_{2\delta_{il}} \right)$$

$$= \frac{2}{i} \delta \omega_j \varepsilon_{jil} \alpha_l = 2i (\delta \vec{\omega} \times \vec{\alpha})_i$$

The last two terms on p. 4 are

$$i \alpha \cdot \delta \vec{\omega} \times \vec{\nabla} \psi + i \delta \vec{\omega} \times \vec{\alpha} \cdot \vec{\nabla} \psi = 0$$

$$\text{since } \vec{\alpha} \cdot \delta \vec{\omega} \times \vec{\nabla} = \vec{\alpha} \times \delta \vec{\omega} \cdot \vec{\nabla}$$

4. Since $G = \frac{1}{2} J_{\mu\nu} \delta \omega^{\mu\nu} = \delta \vec{\omega} \cdot \vec{J} + \delta \vec{v} \cdot \vec{N}$

$$\delta \omega^{ij} = \varepsilon^{ijk} \delta \omega_k, \quad \delta \omega^{0i} = \delta v^i$$

$$\frac{1}{2} J_{ij} \varepsilon^{ijk} = J_k, \quad J_{ij} = \varepsilon_{ijk} J_k$$

$$J_{0i} = -J^0_i = N_i$$

Then, since $[J_i, J_j] = i \varepsilon_{ijk} J_k$

$$[J_i, N_j] = i \varepsilon_{ijk} N_k, \quad [N_i, N_j] = -i \varepsilon_{ijk} J_k$$

$$\begin{aligned} \delta x^\mu &= \delta \omega^{\nu\mu} x_\nu \\ \Rightarrow \delta \vec{x} &= \delta \vec{\omega} \times \vec{x} \quad R \\ \delta t &= -\delta \vec{v} \cdot \vec{x} \quad B \\ \delta \vec{x} &= -\delta \vec{v} t \end{aligned}$$

consider $[J_{12}, J_{13}] = [J_3, -J_2]$

$$= i J_1 = i J_{23}$$

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$$[J_{12}, J_{30}] = [J_3, -N_3] = 0$$

$$[J_{12}, J_{10}] = [J_3, -N_1] = -i N_2 = -i J_{02}$$

$$[J_{10}, J_{20}] = [-N_1, -N_2] = -i J_3 = -i J_{12}$$

By cyclic permutation we can deduce all the others. Since exactly two indices must be equal (otherwise we get zero), we therefore have

$$[J_{\mu\nu}, J_{\kappa\lambda}] = i [g_{\mu\kappa} J_{\nu\lambda} - g_{\nu\kappa} J_{\mu\lambda} - g_{\mu\lambda} J_{\nu\kappa} + g_{\nu\lambda} J_{\mu\kappa}]$$

since it must be antisymmetric under $\mu \leftrightarrow \nu, \kappa \leftrightarrow \lambda$.

Alternative argument

$$[J_{ij}, J_{kl}] = \epsilon_{ijm} \epsilon_{kln} [J_m, J_n]$$

$$= \epsilon_{ijm} \epsilon_{kln} i \epsilon_{mnp} J_p$$

Lemma: $\epsilon_{ijm} \delta_{nq} - \epsilon_{jmn} \delta_{iq} + \epsilon_{mni} \delta_{jq} - \epsilon_{nij} \delta_{mq} = 0$

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Proof: evidently the LHS = 0 if $i=j$, $i=m$, or $j=m$
● If $i \neq j \neq m$ say $i=1, j=2, m=3$ LHS =

$$\begin{aligned} & \epsilon_{123} \delta_{nq} - \epsilon_{23n} \delta_{1q} + \epsilon_{3n1} \delta_{2q} - \epsilon_{n12} \delta_{3q} \\ &= \delta_{nq} - \delta_{1n} \delta_{1q} - \delta_{2n} \delta_{2q} - \delta_{3n} \delta_{3q} = 0 \checkmark \end{aligned}$$

$$\text{Then } [J_{ij}, J_{kl}] = i [\epsilon_{jmn} \delta_{iq} - \epsilon_{mni} \delta_{jq} + \epsilon_{niz} \delta_{mq}] \epsilon_{mqp} \epsilon_{kln} J_p$$

$$= i [\epsilon_{jmn} \epsilon_{kln} J_{mi} - \epsilon_{mni} \epsilon_{kln} J_{mj} + \epsilon_{niz} \epsilon_{kln} \cancel{J_{mm}}^0]$$

$$= i [(\delta_{jk} \delta_{ml} - \delta_{jl} \delta_{mk}) J_{mi} - (\delta_{ml} \delta_{ik} - \delta_{mk} \delta_{il}) J_{mj}]$$

$$= i [\delta_{jk} J_{li} - \delta_{jl} J_{ki} - \delta_{ik} J_{lj} + \delta_{il} J_{kj}]$$

$$= i [\delta_{ik} J_{jl} - \delta_{il} J_{jk} - \delta_{jk} J_{il} + \delta_{jl} J_{ik}]$$

which is the space version of desired result.

$$[J_{ij}, J_{k0}] = \varepsilon_{ijl} [J_l, -N_k]$$

$$= -\varepsilon_{ijl} i \varepsilon_{lkp} N_p$$

$$= -i(\delta_{ik} \delta_{jp} - \delta_{ip} \delta_{jk}) J_{0p}$$

$$= -i(\delta_{ik} J_{0j} - \delta_{jk} J_{0i})$$

$$= i[\delta_{ik} J_{j0} - \delta_{jk} J_{i0}]$$

which is the desired time-space result. Finally

$$[J_{i0}, J_{j0}] = [-N_i, -N_j] = -i\varepsilon_{ijk} J_k$$

$$= -i J_{ij} = i g_{00} J_{ij}$$

which completes the proof.

$$\text{For spin-}\frac{1}{2}: \vec{J} = \frac{1}{2} \vec{\Sigma}, \quad \vec{\Sigma} = \frac{1}{2i} \vec{\alpha} \times \vec{\alpha}$$

$$\vec{N} = +\frac{i}{2} \vec{\alpha}$$

$$\text{so } J_{ij} = \varepsilon_{ijk} \frac{1}{2} \Sigma_k = \frac{i}{4} [\alpha_i, \alpha_j] = \frac{i}{4} [\gamma_i, \gamma_j]$$

$$\text{since } \gamma^0 \alpha_i \gamma^0 \alpha_j - \gamma^0 \alpha_j \gamma^0 \alpha_i = -\alpha_i \alpha_j + \alpha_j \alpha_i$$

$$\text{and } J^0_i = -N_i = \frac{i}{2} \alpha_i = \frac{i}{4} [\gamma^0, \gamma_i]$$

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since $\frac{1}{2}[\gamma^0, \gamma^0 \alpha_i] = \frac{1}{2}\gamma^{02}\alpha_i - \frac{1}{2}\gamma^0\alpha_i\gamma^0 = \alpha_i$

• $\therefore J^{\mu\nu} = \frac{1}{2}\sigma^{\mu\nu}$ where $\sigma^{\mu\nu} = \frac{i}{2}[\gamma^\mu, \gamma^\nu]$
 and $\sigma^{ij} = \varepsilon_{ijk}\Sigma_k$, $\frac{1}{2}\sigma^{0i} = \frac{i}{2}\alpha_i = -N_i$
